

SPACE CHARGE EFFECTS IN THE KEK-BOOSTER SYNCHROTRON

Chihiro Ohmori

Institute for Nuclear Study, University of Tokyo, Tokyo, Japan

Toshikazu Adachi, Tadamichi Kawakubo, Motohiro Kihara and Isao Yamane
National Laboratory for High Energy Physics (KEK), Ibaraki, Japan

Space charge effects have been observed in the KEK-PS Booster synchrotron. Beam-profile measurements by the BEAMSCOPE method show that a beam blow-up in the vertical plane occurred during the RF capture process. The intensity dependence of beam blow-up and deformation of beam profile were observed.

the pulse width is $200\mu\text{s}$. M1-M8 are alternating-gradient magnets and have a focusing order of FDDF. The beam scraper is located in straight section S3, where the bump orbit has its extremum.

I. INTRODUCTION

The KEK-PS Booster is a 40- to 500-MeV proton synchrotron, which is the injector of the KEK 12 GeV Proton synchrotron, and also delivers a beam to the Booster Synchrotron Utilization Facility (BSF). It is a compact alternating-gradient synchrotron operated at fixed tunes. The repetition rate is 20Hz and the synchrotron magnets are excited by a resonance circuit. The linac beam is injected by charge exchange injection and is adiabatically captured in an RF bucket. A bunch is formed in $600\mu\text{s}$ and the bunching factor becomes ~ 0.3 . A transverse beam blow-up at the early stage of acceleration has been observed [1]. Since the Booster synchrotron was designed for H^+ multi-turn injection, it has apertures of 272π mm mrad and 49π mm mrad in the horizontal and vertical planes, respectively [2]. This vertical emittance growth becomes crucial as the beam intensity increases. Since the acceleration period is rather short, we chose the BEAMSCOPE method [3], which can measure a beam profile in $100\mu\text{s}$. It is difficult to use a wire scanner and other methods for a quick measurement. In order to measure the beam profile exactly, a precise calibration was performed [1].

II. MEASUREMENT

The beam-profile measurement system comprises a beam scraper and two pulsed bump magnets which kick the beam in the vertical direction. Figure 1 shows a plan view of the Booster synchrotron. Vertical bump magnets are installed in straight sections S2 and S4. Since the phase advance between the magnets is π radians and the beta functions at both magnets are same, a bump orbit is formed from S2 to S4 in the vertical plane. The magnets are connected in series and excited by a pulsed current having the same polarity. The pulse shape is a half-sine and

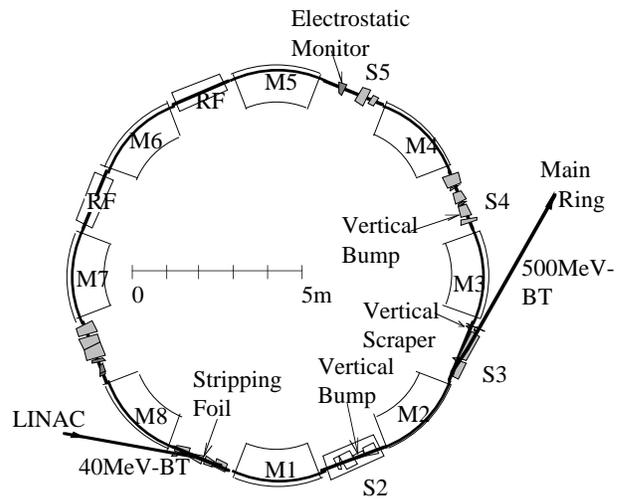


Fig. 1. Plan of the Booster synchrotron.

The ratio of the current of the bump magnets and the height of bump orbit at the scraper was precisely calibrated by means of a pencil beam, which was much smaller than the ring acceptance. In order to measure the beam profile exactly, it is essential to calibrate the bump current to the beam displacement with good accuracy. The ratio is reproducible and independent of the beam parameters, except for the energy. The beam scraper was movable and located at the optimum position where the aperture was minimum in the ring, but did not disturb the circulating beam. The current of the bump magnets was optimized to be able to scrape the whole beam as slowly as possible. Figure 2 shows a schematic view of the beam-profile measurement system. The current of the bump magnets and the beam-bunch signal were recorded simultaneously using a digital sampling oscilloscope, which records the all data during $200\mu\text{s}$ with a sampling rate of 200MHz. The current of the bump magnets and bunch signal were measured by a current transformer and a fast electrostatic monitor (ESM), respectively.

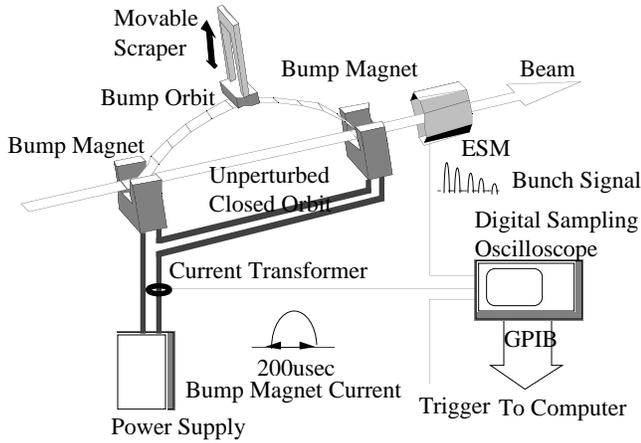


Fig. 2. Schematic of the profile measuring system.

We measured the intensity dependence of the beam profiles in the vertical plane for intensities from $1.8 \cdot 10^{11}$ to $2.2 \cdot 10^{12}$ particle per bunch (ppb). Since we wanted to use the same initial beam distributions for the different beam intensities, a beam slit of mesh shape was installed at the entrance of the 40-MeV Linac. Figures 3 and 4 show the normalized emittance containing 90% of the beam and normalized $1-\sigma$ emittance, respectively, against a delay time of up to 5 ms from injection for various intensities. In these figures, the emittance at injection, i.e. Delay Time = 0 ms, was obtained from the beam profile taken by moving the beam scraper without exciting the bump magnets, since the ESM cannot be available during injection where a bunch structure has not been sufficiently formed. The initial beam emittances are constant for different intensities for both 90% and $1-\sigma$ emittance (Figs. 3 and 4). The initial distribution was independent of the beam intensity. Figures 3 and 4 also show the emittance growth which occurred before ~ 1 ms.

The emittance growth is clearly observed at intensities larger than $1.66 \cdot 10^{12}$ ppb in Fig. 3. The emittance growth between intensities of $1.4 \cdot 10^{12}$ ppb and $1.66 \cdot 10^{12}$ ppb is much larger than that between intensities of $1.66 \cdot 10^{12}$ ppb and $2.2 \cdot 10^{12}$ ppb. However, we should note that there might be a small amount of beam loss for $2.2 \cdot 10^{12}$ ppb. The emittance growth is observed at intensities larger than $1 \cdot 10^{12}$ ppb. The normalized $1-\sigma$ emittance becomes larger according to the increment of the beam intensity. The threshold intensity for $1-\sigma$ emittance is about 40% less than that for 90% emittance. Figure 5 shows the FWHM and the width which contains 90% of the beam against the beam intensities. It clearly shows that the intensities which cause the emittance growths are different for the beam core and edge. The difference of the threshold intensities for

emittance growths suggests that it is caused by space charge effects.

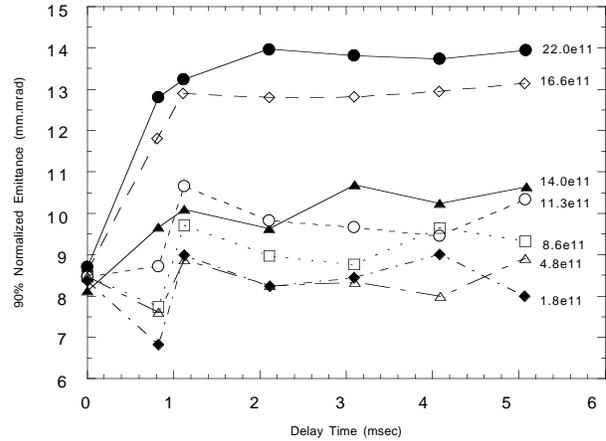


Fig. 3. Intensity dependence of the normalized emittance which contains 90% of the beam.

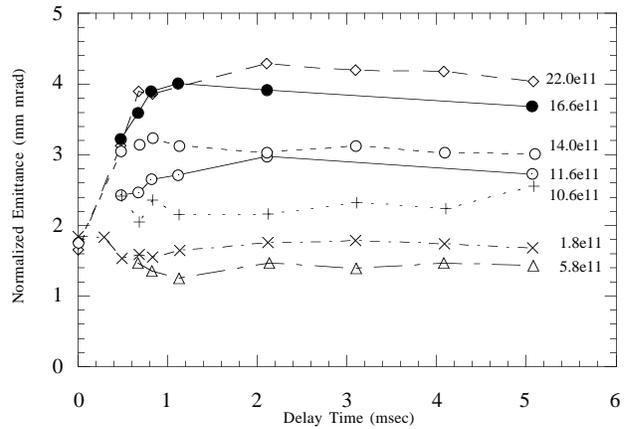


Fig. 4. Intensity dependence of the normalized $1-\sigma$ emittance.

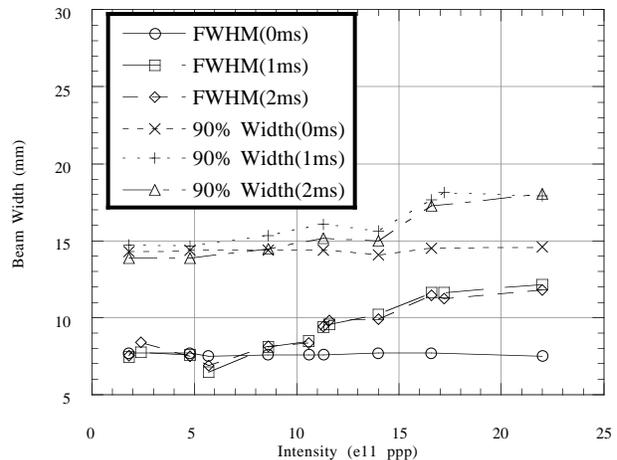


Fig. 5. Intensity dependence of the beam widths.

III. SPACE CHARGE TUNE SHIFT

The Booster tune space is shown in Fig. 6. In order to estimate the space charge detuning in the smooth approximation for $\Delta p/p=0$, we use the expression [4] $\Delta Q_{H,V} = N r_p G [\pi B \beta^2 \gamma^3 (\epsilon_{H,V} + (\epsilon_H \epsilon_V Q_{H,V} / Q_{V,H})^{1/2})]^{-1}$, where N =number of protons, $r_p = 1.535 \cdot 10^{-18}$ m, and G is a form factor which depends on the transverse intensity distribution. For $1.0 \cdot 10^{12}$ ppp at the end of injection, where $\beta^2 \gamma^3 = 0.09$ and the bunching factor is ~ 1 , the 90% beam width was 14mm. Using $G=1.56$ for the Gaussian distribution, we obtain $\Delta Q_{ySC} = -0.10$ and $\Delta Q_{xSC} = -0.05$, as shown in Fig. 6(a), bounded by the bold line. At 1 ms after injection, where emittance growth has occurred for the beam core but not for the beam tail, the beam is captured in the RF bucket and the bunching factor becomes 0.3. Then, $\Delta Q_{ySC} = -0.34$ and $\Delta Q_{xSC} = -0.17$ are obtained for the same beam widths. This is the lowest point of the region in Fig. 6(a) bounded by the solid line, and touches both $Q_x=2$ and $Q_y=2$. For $1.66 \cdot 10^{12}$ ppp at the end of injection, we obtain $\Delta Q_{ySC} = -0.17$ and $\Delta Q_{xSC} = -0.08$, as shown in Fig. 6(b), bounded by the bold line. For bunched beam showing emittance growth for both the beam core and edge, the 90% beam width became 18mm 1 ms after injection. We obtain $\Delta Q_{ySC} = -0.40$ and $\Delta Q_{xSC} = -0.26$ which cross both $Q_x=2$ and $Q_y=2$ resonances largely, as shown in Fig. 6(b).

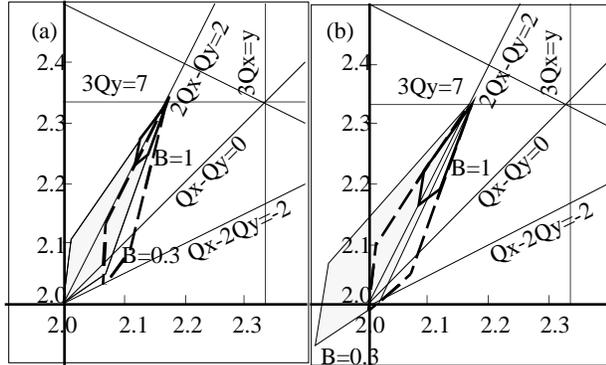


Fig. 6 KEK-PS Booster Tune Spaces at $1.0 \cdot 10^{12}$ ppp (a) and $1.66 \cdot 10^{12}$ ppp (b).

However, the real beam is not smooth and $\Delta p/p \neq 0$. For a more accurate estimation we used the computer code space SPACEX [5], which can solve envelope equations including the space charge force. The results are also shown in Fig. 6. We assume that the momentum spread was 0.30% and that the real beam was larger in the horizontal plane due to it. Using $G=1.56$ and $B=0.3$, we obtain $\Delta Q_{ySC} = -0.30$ and $\Delta Q_{xSC} = -0.11$ for $1 \cdot 10^{12}$ ppp. It is shown in Fig. 6(a), bounded by the dashed line. We also

obtained $\Delta Q_{ySC} = -0.34$ and $\Delta Q_{xSC} = -0.17$ for $1.66 \cdot 10^{12}$ ppp. This is shown in Fig. 6(b), bounded by the dashed line. Since the momentum spread was taken into account, the tune spreads obtained by SPACEX are smaller than those by Laslett formula. Since the 90% beam width during the emittance growth was smaller than 18 mm, the tune spread should have been much larger than the region in Fig. 6(b) bounded by the dashed line. It is reasonable that the large tune spread which crosses the integer and other resonances causes emittance growth, which stops when the tune spread becomes smaller.

IV. DISCUSSIONS

The emittance growth in the vertical plane can be explained based on the space charge effects. Operation at higher tunes and resonance correction techniques are used to avoid and overcome the effects. However, the Booster synchrotron is a small accelerator and has limited space. It is difficult to install the necessary instruments to change the tunes and for corrections. Reduction of the peak intensity by the second-harmonic buncher is another way, and might be possible if the instrument is designed to be compact.

V. REFERENCES

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