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BREMSSTRAHLUNG BY THE BUNCH OF ULTRARELATIVISTIC CHARGED PARTICLES INTO A THICK TARGET.

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1. INTRODUCTION.

Bremsstrahlung of a system of classically fast charged particles which do not interact with each other but which do undergo multiple elastic scattering by randomly positioned atoms of the medium is studied. We derived the spectrum of the bremsstrahlung of particles through a such analysis οf systimatic kinetic radiation process in the medium. is shown that the spectral Τt emission distribution of the the bremsstrahlung energy of depends significantly on both the characteristics of the scattering of the particles in the medium and the parameters characterizing the minitial set of the particles.

2.STATEMENT OF THE PROBLEM.

consider the system of We charged ultrarelativistic (E \gg m) clasically fast $(E \gg 0)$ is a frequency) particles radiation which do not interact with each Ε, and e are the other m, (the mass and the charge energy, These particle, h=c=1). of the homogeneous, particles enter а

semi-infinite, amorphous scattering medium. In the initial period,

t=0, particles are located in the points \vec{r}_{o1} , \vec{r}_{o2} , ..., \vec{r}_{oN} , and are the velocity \vec{v}_{o1} , \vec{v}_{o2} , ..., \vec{v}_{oN} , equal to $v_o = [1 - (m/E)^2]^{1/2}$ and they are directed under the $|\Delta_{\mu}| \ll \ll 1$ ($\mu = 1, ..., N$ - is the number of the particles) angle to the \vec{e}_z vector (vector of the inward normal to the boundary of the medium). Let the characteristic

longitudinal size of the beam L_b be such that $L_b v_o^{-1} \ll T$ (the time when the particles move in the medium).

3. SOLUTION OF THE PROBLEM.

The specrtal distribution of the energy emitted by these particles is

$$\begin{split} \frac{\mathrm{d}\varepsilon_{\omega}}{\mathrm{d}\omega} &= \frac{\mathrm{e}^{2}\omega^{2}}{2\pi^{2}} \frac{\mathrm{N}}{\mathrm{Re}} \sum_{\mu,\nu=1}^{\pi} \mathrm{fd}\Omega_{\pi} \int \mathrm{d}t \int \mathrm{d}\tau \left[\tilde{\pi}_{x} \tilde{\nabla}_{\mu} \right] \\ & \mu,\nu=1 \qquad \alpha \qquad (1) \\ & \cdot \left[\tilde{\pi}_{x} \tilde{\nabla}_{\nu} \right] \cdot \exp\{-\mathrm{i}\omega\tau + \mathrm{i}\tilde{R} \cdot \left(-\tilde{r}_{\mu} - \tilde{r}_{\nu} + \right. \\ & \left. + \left. \tilde{r}_{o\mu} - \tilde{r}_{o\nu} \right) \right\}, \end{split}$$

It is the wave vector of the radiation field, $d\Omega_{\pm}$ is an element of solid angle in the direction

 $\vec{n} = \vec{k} / k = \vec{k} / \omega$, $\vec{r}_{\mu} = \vec{r}_{\mu} (t + \tau)$, $\vec{\nabla}_{\mu} = \vec{\nabla}_{\mu} (t + \tau)$, condition [2]: $\vec{r}_{y} = \vec{r}_{y}(t)$, $\vec{v}_{y} = \vec{v}_{y}(t)$ are coordinates and velocities of the particles at the time $t+\tau$ and t respectively, τ is the time scale of the radiation formation (the coherence time), is the time at which and the t emitted. the radiation is

the observed То calculate spectral distribution of the particles emission energy of the in the medium, dE_ω /d ω , we must average expression (1) over all possible trajectories of the the scattering particles in matter [1]. To solve this problem necessary to find [2]it. is two-time distribution function of particles scattering the in а medium. the of Ιn case ultrarelativistic classically fast solution of particles the the problem is determined by the Fourier component of this function

$$F_{\mathbf{k}} \left(\vec{\nabla}_{\mu}, \vec{\nabla}_{\nu}, \mathbf{t}, \tau \right):$$

$$\frac{dE_{\omega}}{d\omega} = \frac{\mathbf{e}^{\mathbf{z}}\omega^{\mathbf{z}}}{2\pi^{\mathbf{z}}} \operatorname{Re} \sum_{\mu, \nu=1}^{N} \int d\Omega_{\mathbf{z}} d\vec{\nabla}_{\mu} d\vec{\nabla}_{\nu} \cdot \int d\mathbf{t} \int d\tau$$

$$[\vec{n}_{\mathsf{X}}\vec{\nabla}_{\mu}] \cdot [\vec{n}_{\mathsf{X}}\vec{\nabla}_{\nu}] \cdot \exp\{-\mathrm{i}\omega\tau + \mathrm{i}\vec{R} \cdot (\vec{r}_{\mu} - \vec{r}_{\nu} + \mathbf{i}\omega\tau + \mathrm{i}\omega\tau + \mathrm{i}\omega\tau$$

+
$$\mathbf{F}_{o\mu} - \mathbf{F}_{o\nu}^{\mathbf{Z}}$$
) } \cdot \mathbf{F}_{k}(\mathbf{V}_{\mu}, \mathbf{V}_{\nu}, \mathbf{t}, \tau)

∛_η, t,τ) The function F_k(∛_u, following the satisfies to equations and the initial

$$\frac{\partial F_{k}(t,\tau)}{\partial \tau} - i \mathbf{k} \cdot \mathbf{v}_{\mu}(\mathbf{\eta}) \cdot F_{k}(t,\tau) =$$

$$= \frac{q}{4} \cdot \frac{\partial F_{k}(t,\tau)}{\partial \mathbf{\eta}}$$

$$\frac{\partial}{\partial t} F_{\nu}(t,0) = -iR(\vec{\nabla}_{\mu}(\vec{\eta}) - \vec{\nabla}_{\nu}(\vec{\varphi}))F_{k}(t,0) =$$

$$= \frac{q}{4} \cdot \left[\frac{\partial}{\partial \eta} + \frac{\partial}{\partial \varphi} \right]^{2} F_{k}(t,0)$$

$$\mathbf{F}_{\mathbf{k}}(\mathbf{0},\mathbf{0}) = \delta(\eta - \Delta_{\mu}) \cdot \delta(\phi - \Delta_{\nu}).$$

$$\cdot \exp[i\mathbf{R} \cdot (\mathbf{T}_{ou} - \mathbf{T}_{ov})]$$
,

where q is the average square of multiple scattering angle per unit of path [3], \uparrow and ϕ the angle vectors which are connected with $\vec{\nabla}_{\mu}(\vec{\eta})$ and $\vec{\nabla}_{\nu}(\vec{\Phi})$ by the ordinary formulae of the theory of multiple scattering in a amorphous medium [3].

4.RESULTS

We have constructed a consistent kinetic theory for the radiation emission by a system of

classically fast noninteracting charged particles which undergo multiple elastic collisions in a scattering matter. We have found the spectral distribution of the such emission energy from particles. The obtained spectrum strongly both on the depends parameters of the scattering medium, and on the characteristics of the initial system οf the particles.

We have studied in detail the emisson by the beam of idetical particles. It is shown in this that the spectral case distribution οf the emission energy has at least one extremum in contradiction to the situation of an individual particle [1] when the emission spectrum is a monotonic function from emission frequency. If the pulse beam of identical particles takes place then the extremum is a maximum. the width i f of the Moreover initial beam D is such that the condition $qD\xi^{-a} \ll \xi(qT)^{-1/2} \ll 1$ the maximum of holds. bremsstrahlung energy spectrum is a plateau with a width on the $D^{-1}(qT)^{-1/2}$. The ratio of order $(dE_{\omega} \neq d\omega) = max$ to the background level ($(dE_{ij}/d\omega)_{o} = 2e^2 qT/3\pi\xi^3$) is approximately equal to N , the number of emitting particles. As the parameter $qD\xi^{-a}$ increases

 $qD\xi^{-3} \leq \xi \cdot (qT)^{-1/2} \ll 1$ the plateau convert into "strict" maximum. As before, we have $(dE_{\omega} \not d\omega)_{max} \cdot (dE_{\omega} \not d\omega)_{o}^{-1} \cong N.$ If we have $qD\xi^{-3} \gg 1$, then the quantities $(dE_{\omega} \not d\omega)_{max}$ and $(dE_{\omega} \not d\omega)_{o}$ become approximately equal to 1.

The bremsstrahlung by a higly anisotropic point source of ultrarelativistic particles is investigated. We show that in this case the spectral disrtibution of the emission energy has a maximum. The last is unique. The shape of

the maximum as well as the magnitude of the emission energy depend strongly both on the value $X_{o} \ll 1$ of cone vertex angle which

the includes the initial velocities of the particles and on on the scattering properties of a scattering medium.

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