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Abstract

Octupoles may be used to correct the third order spherical aberration of quadrupole transport systems. Crosstalk in the coupling of an octupole placed at a given point causes it to add a term with the wrong sign in the y-channel if it has the right sign in the x-channel, thus severely reducing efficiency. It is often convenient to utilize a special correcting section insertion which is seen as a +I transfer matrix by the first order focussing. Within point-to-point thin lens optics we give two-parameter systems with 16 magnets having locations with large S_x where $S_y = 0$ and vice versa for octupole placement.

The problem is to set up a -I quadrupole system which contains two special locations. At the first, S_x is large and S_y is zero; while at the second the reverse is true. Octupoles are placed at these locations and make corrections free of channel crosstalk. The layout is:

-I System

Section 1

subsection 1A	Note: Part B is the same as A but with the order of the elements reversed. Subsection 2 is the same as 1 but with the polarities of the quadrupoles reversed. Section 1 is X = -I, Y = +I. Section 2 is X = +I, Y = -I.
part 1Aa	
part 1Ab	
subsection 1B	
part 1Ba part 1Bb	

Section 2

subsection 2A
part 2Aa
part 2Ab
subsection 2B
part 2Ba
part 2Bb

The basic module for the 1Aa part is a five element OFODO sequence of drift spaces and thin lenses and part 1Aa consists of two such modules notated in English and Greek respectively as

1Aa:	ξ P m Q n	λ π μ τ ν
	one module	one module
	1 Aa	1Ab

For example, the second element (P) after the source has transfer matrix

$$\begin{pmatrix} 1 & 0 \\ -P & 1 \end{pmatrix}$$

in the X-channel. Each module is one of the modes:

PP = point-to-point	= $\begin{pmatrix} x & 0 \\ x & x \end{pmatrix} R_{12} = 0$
LP = parallel-to-point	= $\begin{pmatrix} 0 & x \\ x & x \end{pmatrix} R_{11} = 0$
PL = point-to-parallel	= $\begin{pmatrix} x & x \\ x & 0 \end{pmatrix} R_{22} = 0$
LL = parallel-to-parallel	= $\begin{pmatrix} x & x \\ 0 & x \end{pmatrix} R_{21} = 0$

For the present model we have chosen

X_{1A}		PL	and	LP	
Y_{1A}		PP	and	LL	
X_{1Aa}	LL	Y_{1Aa}	LP	and	PL
X_{1Ab}	LP	Y_{1Ab}	LP	and	PL

This decomposition makes the problem tractable because the one LL zero of X_{1Aa} and the two (PL and LP) zeroes of Y_{1Aa} are easy to solve. By matrix multiplication

$$(LP \text{ and } PL) \cdot (LP \text{ and } PL) = PP \text{ and } LL$$

$$Y_{1Aa} \cdot Y_{1Aa} = Y_{1A}$$

$$LP \cdot LL = LP$$

$$X_{1Ab} \cdot X_{1Aa} = X_{1A}$$

Thus Y_{1A} is automatically PP and LL, while X_{1A} needs to have one condition satisfied ($X_{1A}^{(22)} = 0$) to become LP and PL. We should also like to specify, at the end of subsection 1A, the value of S_x ($= [X_{1A}]_{12}$). A count shows that $3+3+1+1 = 8$ restrictions have been put on the 10 parameters in the two modules of 1Aa leaving two (m and μ) free. One finds, with notation

$$a = m Q$$

$$b = \pi \mu$$

the expressions

$$\xi = \frac{b(1-a^2)(2-b^2)S_x}{(1+b)(1-b^2)(1+a+2a^2)}$$

$$P = \frac{(1+a+2a^2)(1-b^2)(1+b)}{2ab(1+a)(2-b^2)S_x}$$

$$m = \frac{2a^2b(2-b^2)S_x}{(1+a+2a^2)(1-b^2)(1+b)}$$

$$Q = \frac{(1+a+2a^2)(1-b^2)(1+b)}{2ab(2-b^2)S_x}$$

$$n = \frac{b(2-b^2)(1+2a)S_x}{(1+a+2a^2)(1-b^2)(1+b)}$$

$$\lambda = \frac{(1-b-b^2)S_x}{(1+b)^2(1+a)}$$

$$\pi = \frac{(1+b)(1+a)}{bS_x}$$

$$\mu = \frac{b^2S_x}{(1+a)(1+b)}$$

$$\tau = \frac{(1+a)}{b(1-b)S_x}$$

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$$v = \frac{(1-b) S_x}{1+a}$$

$$S_y = (Y_{1A})_{12} = \frac{b(1+a) (2-b^2) S_x}{(1-b^2) (1+a+2a^2)}$$

where S_x is the peak of the X-channel sinelike function. The system allows all a, b choices within the limits

$$0 < a < 1$$

$$0 < b < \frac{\sqrt{5}-1}{2} = 0.618034$$

to keep all drift distances positive.

For a trial system it was found

$$(S_x)_{\max} = (S_y)_{\max} = 30 \text{ m}$$

was adequate to keep the octupole strength near 1 T/m. As an example we have chosen

$$a = 0.85 \quad b = 0.20$$

in an attempt to get a small total length

$$L = 4(\lambda+m+n+\lambda+\mu+v).$$

We find (MKS units)

$$\begin{aligned} \lambda &= 0.85973117 & \lambda &= 8.5585586 \\ P &= 0.10263132 & \pi &= 0.37000000 \\ m &= 4.4767956 & \mu &= 0.54054056 \\ Q &= 0.18986795 & \tau &= 0.38541667 \\ n &= 8.36494947 & v &= 12.972973 \end{aligned}$$

The resulting cosinelike and sinelike trajectories are listed below

s	C _x	S _x	C _y	S _y
0.0	1.0	0.0	1.0	0.0
0.85973117	1.0	0.8597311	1.0	0.85973117
5.3365268	0.54054056	4.9415152	1.459	5.7315382
13.701476	0.54054056	20.416672	0.0	5.7315382
22.260035	0.54054056	36.250005	-1.4932394	5.7315382
22.800575	0.43243243	30.000000	-1.8861971	6.8778458
35.773548	0.0	30.000000	-1.8861971	0.0

$$X_{1A} = \begin{pmatrix} 0 & 30 \\ -1 & 0 \\ 30 & \end{pmatrix}$$

$$Y_{1A} = \begin{pmatrix} -1.8861971 & 0 \\ 0 & -0.53016733 \end{pmatrix}$$

Octupoles are placed at 1A and 2A. A similar type of problem was solved for a beam rotator.

Alternate solutions are possible. One needs to consider only subsection 1A, which is divided into parts 1Aa, 1Ab, and repetitions as specified above. A part a or b may contain two or three magnets. There are sixteen cases, based on the following matrix multiplication table:

$$\begin{aligned} PL \cdot LP &= LL & LL \cdot PL &= PL \\ PL \cdot PP &= PL & LL \cdot LL &= LL \\ LP \cdot PL &= PP & PP \cdot LP &= LP \\ LP \cdot LL &= LP & PP \cdot PP &= PP \end{aligned}$$

In the list below, N indicates the number of magnets per module. The example₂ given above corresponds to the first line of this list.

X _{Ab}	X _{Aa}	Y _{Ab}	Y _{Aa}	N _{Ab}	N _{Aa}	X _A	Y _A
LP	LL	} (PL & LP)		2	2	LP	} (LL & PP)
LL	PL			2	2	PL	
PL	PP			3	3	PL	
PP	LP			3	3	LP	
PL	PP	} (LL & PP)		2	2	PL	
PP	LP			2	2	LP	
LP	LL			3	3	LP	
LL	PL			3	3	PL	
PP	LP	} (PL & LP)		3	2	LP	
LP	LL			3	2	LP	
LL	PL			2	3	PL	
PL	PP			2	3	PL	
LL	PL	} (LL & PP)		3	2	PL	
PL	PP			3	2	PL	
PP	LP			2	3	LP	
LP	LL			2	3	LP	

References

1. S. Kowalski and H. Enge, "Beam Rotator", Proc. Fourth Int'l. Conf. on Magnet Technology, CONF-720908, Brookhaven, 1972.
2. Karl Brown has emphasized +I sections as sextupole sacs.