

PRE-INJECTOR OF THE KEK 2.5-GEV LINAC AND HIGH-CURRENT SINGLE-BUNCH ACCELERATION

Satoshi Ohsawa, Jae-Young Choi, Yujiro Ogawa, Atsushi Enomoto, Takuya Kamitani
 Hitoshi Kobayashi, and Isamu Sato
 National Laboratory for High Energy Physics (KEK)
 1-1 Oho, Tsukuba-shi, Ibaraki-ken, 305 Japan

Abstract

The timing stability required for the electron source of the pre-injector is determined by the energy acceptance of the KEKB rings. The jitter requirements for the gun pulse itself are somewhat mitigated by the subsequent bunching process. A system by which a beam trigger stability of 3–20 ps is achieved is described, and the method to estimate the phase jitter of the accelerated beam from the energy jitter is discussed.

Introduction

The KEK 2.5-GeV Linac which is presently being upgraded will provide beams of 8 GeV electrons and 3.5 GeV positrons for direct injection into the KEKB rings [1]. The energy acceptance of the rings is 0.5%, which places constraints on both the drift and the jitter in the energy of the injected beams. This paper discusses the manner in which the energy of the injected beams will be controlled to meet the requirements of the KEKB rings, with emphasis on the energy jitter.

The PF Linac pre-injector has been described elsewhere [2]. The half of the pre-injector is shown schematically in Fig. 1. A 119-MHz subharmonic buncher (SHB1) has been installed last January between the electron gun and a 476-MHz subharmonic buncher (SHB2) to produce S-band single-bunch beams with better bunch purity. A fast grid pulser located near the cathode of the triode gun produces an electron pulse of up to 15 A in a 2 ns pulse. The electron bunch length is first reduced by two pulsed RF cavities operating at respectively the 24th and the 6th subharmonics of the primary 2856 MHz of accelerating RF, then further compressed by two S-band prebunchers and a buncher. The fully bunched beam emerges at the end of the pre-injector with an energy of about 40 MeV, where the beam is momentum analyzed with a 90°-bending system.

The energy of a particle varies as $\cos\theta$, where θ is the angle between the particle and the negative crest of the RF. The buncher of the PF Linac consists of two parts: a bunching section of 6 cells with gradually different phase velocities and a

accelerating section of 29 cells with a phase velocity of the speed of light, where bunched electrons are accelerated at a large phase, typically $\theta = 40^\circ$. Furthermore a short bunch of high intensity must be positioned several degrees ahead of the crest of the accelerating RF in order to minimize the resultant energy spread.

If a particle is accelerated at an angle θ_0 in the accelerating section of the buncher and then at an angle θ_1 in the following accelerating sections, the resultant energy of the particle becomes $E_j = E_0 \cos(\theta_0 + \delta\theta_j) + E_1 \cos(\theta_1 + \delta\theta_j)$ at any position along the beam line of the linac. Then the energy jitter, $\delta E_j = \delta E_j / E_{center}$ is given by

$$\delta E_j = -\frac{E_0 \sin\theta_0 + E_1 \sin\theta_1}{E_0 \cos\theta_0 + E_1 \cos\theta_1} \delta\theta_j \quad (1)$$

where $\delta\theta_j$ is the jitter in the RF phase relative to the beam position. Consequently the deleterious effects of $\delta\theta_j$ grow as θ_0 and θ_1 increase. At the end of the linac where $E_0 \ll E_1$, the energy jitter is given approximately

$$\delta E_j = \delta\theta_j \tan\theta_1. \quad (2)$$

Typically θ_1 is set to about 10° for the 10-nC beam. If the phase jitter $\delta\theta_j$ is 1 ps, the energy jitter becomes as large as 0.3% at the end of the linac. Inevitably the phase jitter should be no more than 1 ps to meet the energy acceptance of 0.5% of the KEKB rings even if the energy compressing system is introduced.

Then it becomes important how to measure accurately the phase jitter as small as 1 ps. Because of large θ_0 and θ_1 , in the pre-injector the phase jitter $\delta\theta_j$ can easily be estimated by measuring the energy jitter δE_j at the exit of the pre-injector.

Source Timing Stability

Any timing change of the beam relative to the accelerating S-band phase results in a beam energy change. Thus, beam timing stability is required to avoid excessive energy jitter or drifts.

Within the first few centimeters of the buncher section the electrons become fully relativistic and subsequently do not

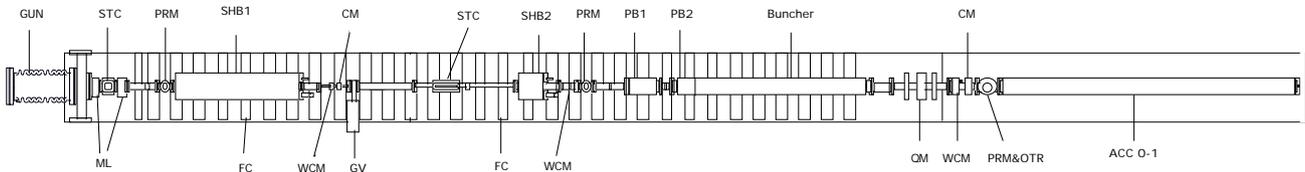


Fig. 1 New layout of the PF linac pre-injector. Abbreviations used are GUN for electron source, SHB1/SHB2 for 119-MHz/476-MHz SHB, PB1, PB2 for S-band prebunchers, and ACC for 2m-accelerating section. Others are beam monitors (screen monitor: PRM, and current monitor: WCM) and components of the beam transport system (magnetic lens: ML, focusing coil: FC, and quadrupole magnet: QM).

change their timing relative to a speed-of-light signal. However, before becoming relativistic the beam is susceptible to timing changes of several origins.

From the bunching process, three potential sources of jitter in the timing of the accelerated electron pulse can be identified: trigger timing jitter; grid pulser instability; and RF phase jitter in the source bunchers. The first two items give the same effect to the source pulse. As far as in the pre-buncher the effect of the last item is thought to be small, since the accelerated beam position is determined by the bunchers relative to the RF, which is supplied from a klystron except for SHBs. Therefore changes in phase relative to the beam can be the result of a shift in the timing of the electron bunch itself.

The basic timing trigger for the grid pulser is a trigger derived from the PF Linac timing system. The problem from the outset has been to place this timing signal on the gun high voltage deck without loss of stability. The fiber optical link which has been used was introducing a jitter of several ten picoseconds ($\sigma = 60$ ps). While it is expected to be possible to transmit the trigger with this improved link which have the required stability, this time as a trial, we placed a CW RF signal (476 MHz) on the high voltage deck where the trigger was resynchronized.

The fiber optical link is shown schematically in Fig. 2. The link utilizes about 50 m of a 100-m graded-index fiber optical cable. A high speed ECL circuit (TD2) was used to resynchronize [3]. With this new link, no timing jitter can be detected in the output of the TD2 above 20 ps threshold of the measuring instrumentation. But of course at $\theta = 10^\circ$, an energy change of 0.3% is produced at the end of the linac when the timing jitter of the accelerated gun pulse shifts by only 1 ps. Fortunately the timing jitter at the gun output is mediated by the bunching process itself.

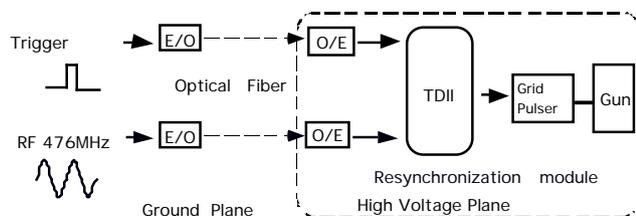


Fig. 2 Schematic diagram of the resynchronization system.

Beam Energy Stability

The overall bunch compression ratio is typically between 50 and 100 as is shown in Fig. 3, which varies depending on both the tuning conditions of the bunching system and the beam charge. The fully bunched beam of 10 nC emerges at the end of the pre-injector with an energy of about 40 MeV, compressed into a bunch length of about 13 ps. Since the bunch compression ratio decreases as the beam charge increases, it becomes difficult to stabilize the energy of the high current beams.

Energy shift of a beam of 10.5 nC was measured at the end of the pre-injector as a function of timing changes introduced in the electron gun trigger circuit. Substituting the measured

energy shift into Eq. (1), we obtain the phase shift of the beam in Fig. 4.

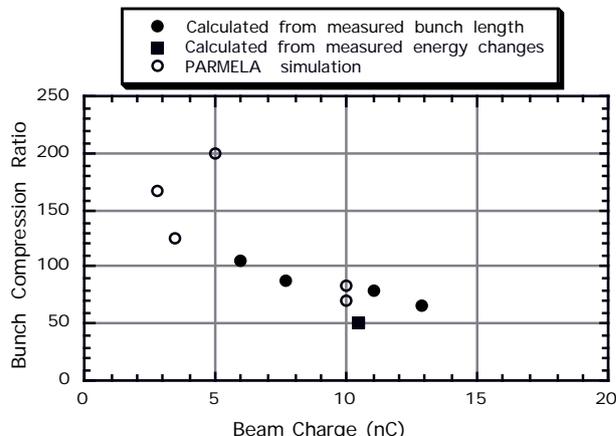


Fig. 3. Bunch compression ratios obtained experimentally and from PARMELA simulations.

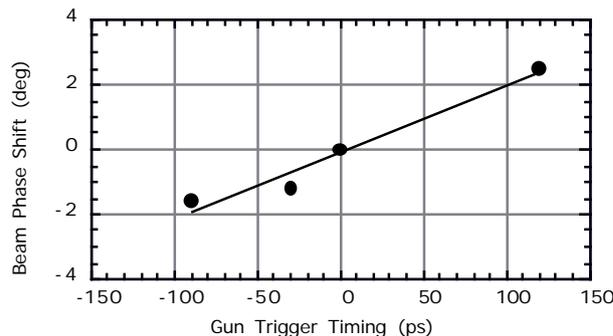


Fig. 4. Phase shift of the beam (10.5 nC) as a function of timing changes introduced in the electron gun trigger circuit. The slope of the linear fit is 0.02 (deg./ps).

The beam energy has been measured by a 90°-bending magnet system at the end of the pre-injector. A vidicon camera is used to view the beam hitting a fluorescent screen in vacuum at the exit of the analyzer where the dispersion is well known. The video is digitized in 2-dimensions to determine both the centered position and the width of the energy spectrum. When a beam charge was as low as 1 nC, the beam was stable and the energy jitter was not seen on the screen monitor. When the beam charge, however, was increased to 5 nC, the energy jitter of ± 200 keV ($\pm 0.5\%$) was observed before resynchronizing the beam trigger. By resynchronizing the timing jitter decreased to less than 20 ps, the energy jitter has been improved to $\pm 0.12\%$ as listed in Table 1.

Table 1. Phase jitter θ_j , θ'_j calculated from measured energy and timing jitter e_j , t_j respectively.

Beam trigger	e_j 3 (%)	θ_j 3 (deg)	t_j 3 (ps)	θ'_j 3 (deg)
not resynchronized	1.0	1.6°	60 ps	1.2°
resynchronized	0.25	0.41°	20 ps	0.4°

Phase jitter σ_j relative to the beam can be calculated from the measured energy jitter σ_{ej} using Eq. (1). On the other hand, from the timing jitter σ_{tj} , the phase jitter σ_j is estimated from the relationship of Fig. 5. The obtained values are well agree in each case, which indicates that energy jitter of the beam is due to the timing jitter of the electron bunch itself as far as pre-injector is concerned.

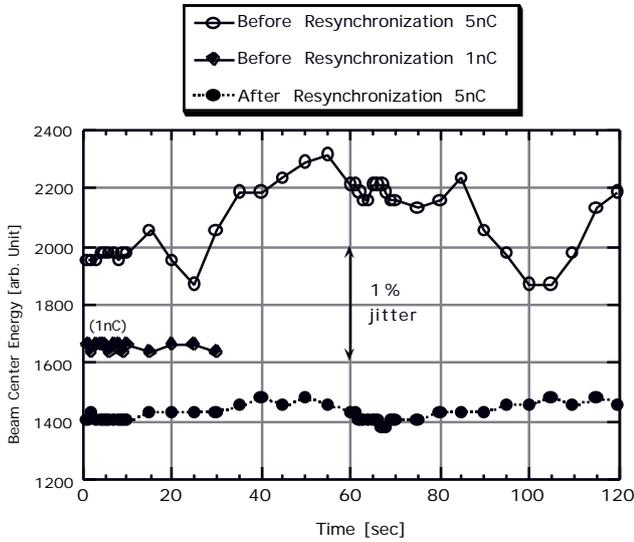


Fig. 5 Beam energy stability measured at the exit of the pre-injector. The beam became stable after the beam trigger was resynchronized.

The fact the the phase jitter was reduced to the required level has been comfermed by accelerating a beam to the end of the linac. Energy changes of a 0.5 nC beam were measured with two conditions. It was actually observed that before being resynchronized the beam energy was changing from every pulse to pulse. After being resynchronized, although residual drift exists in the beam energy, the fast energy jitter disappeared as being seen in Figs. 6a and 6b, which is consistent with the phase jitter ($\sigma_j = 0.12^\circ$) calculated from measured energy and timing jitter.

Conclusion

By resynchronizing the beam trigger, it was confirmed that the fast jitter of the beam energy from pulse to pulse is due to the timing jitter of the source, and reduced to the required level for the KEKB Linac.

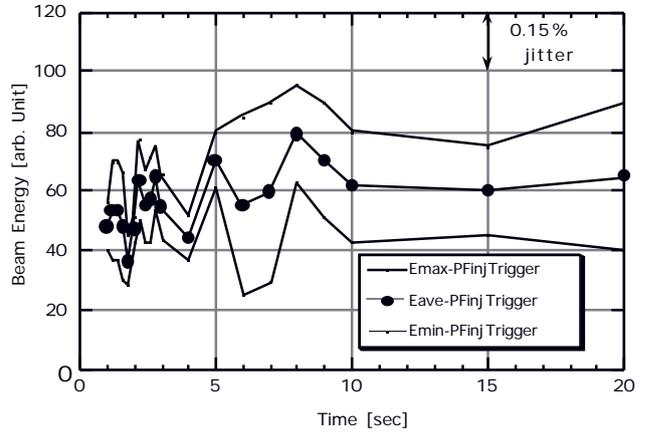


Fig. 6a. Beam energy jitter measured at the end of the PF 2.5-GeV linac. Fast jitter in pulse to pulse was observed before the beam trigger was resynchronized.

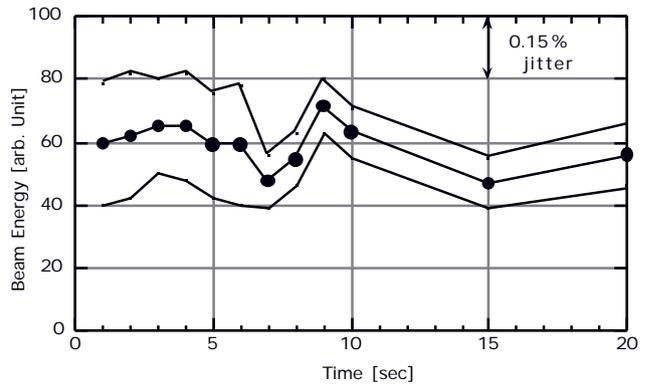


Fig. 6b. Beam energy jitter measured at the end of the PF 2.5-GeV linac. Fast jitter in pulse to pulse disappeared after the beam trigger was resynchronized.

References

- [1] A. Enomoto, "Upgrade to the 8-GeV Electron Linac for KEKB", this proceeding.
- [2] S. Ohsawa et al., "High-Intensity Single-Bunch Beam of the PF 2.5-GeV Linac", Proc. 1994 Int. Linac Conf., 193.
- [3] K. Ishii, "Digital Delay CAMAC Module with 550 Mhz Preset Counter (TD-2), KEK 83-14 , 1983, A/I.