



# DEVELOPMENT OF A QUANTUM ELECTRON BEAM DIAGNOSTIC APPARATUS

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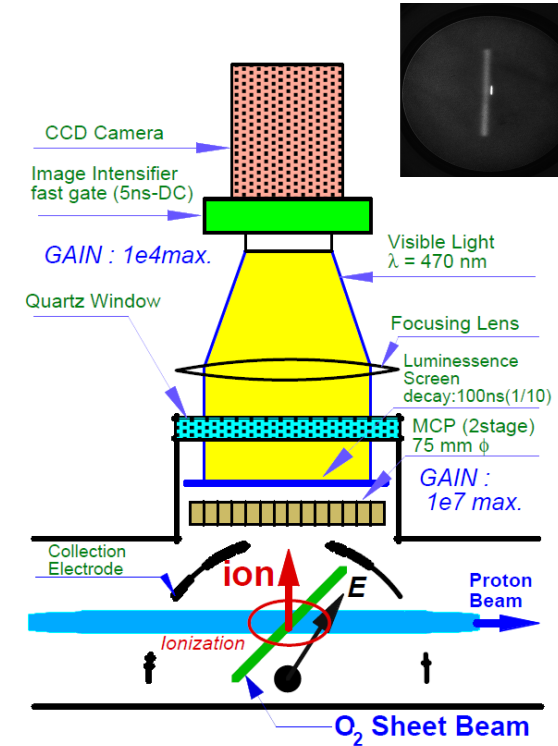
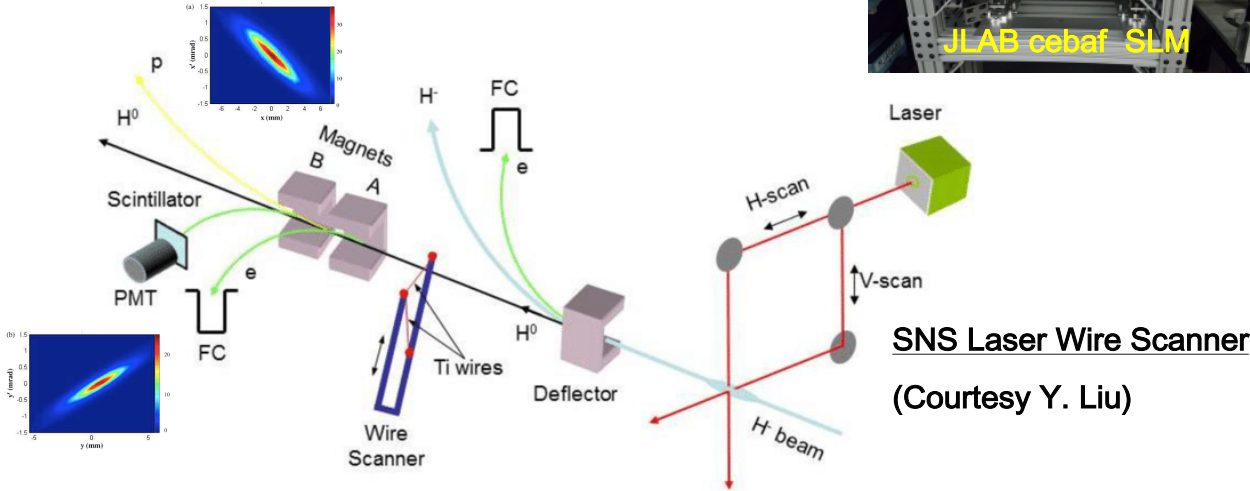
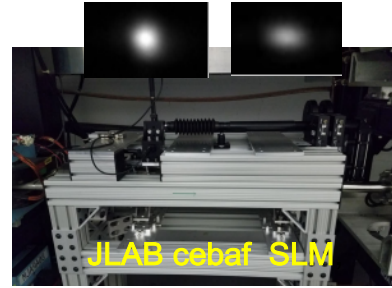
**S. Malinovskaya, A. Ramaswamy** – *Stevens Institute of Technology*

# Outlines

- Motivation
- Principles & Experimental Setup
- Results and Analysis
- Summary & Future Perspectives

# Various Electron Beam Profile Diagnostics

- Phosphor screen, OTR, Wire scanner, EO sampling, Laser scattering,
- Tomography, profile reconstruction,.....



- Is there a non-invasive beam apparatus measuring both longitudinal and transverse spatial profiles?

- Refs:
  - GJBHM (Courtesy A. Jeff)
  - H. Zhang, MONOPT045

# Potential Application - Charged Particles Tracking for NP

- Wire Chamber,
- Bubble Chamber,
- Hodoscope scintillator,
- Gem Chamber, ...

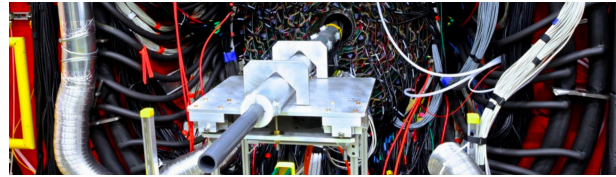
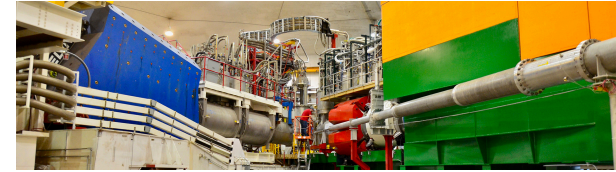
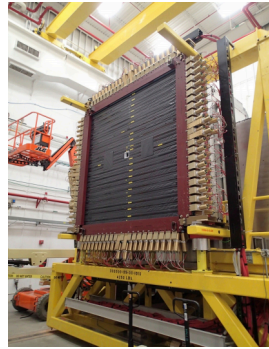
- Each one has limitations

- Common Features:

Sophisticated, Gigantic

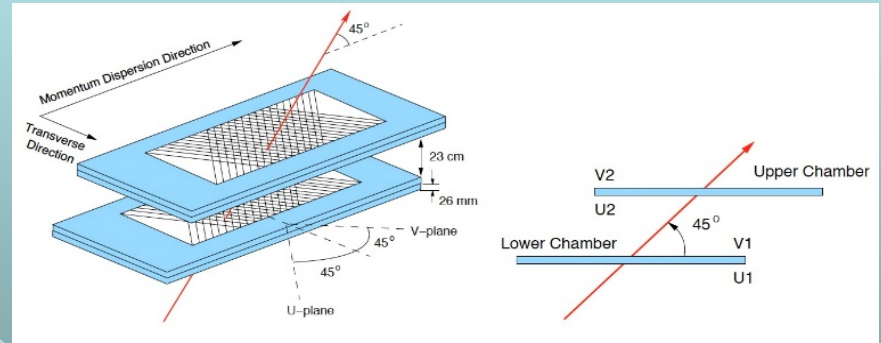
Expensive, Time consuming

- Any improvement in *rate capability*, *cost* or *size* of the tracking device may greatly benefit current and future physics experiments.
- Quantum Tracker maybe a glimpse of hope?



## Wire Chamber

- Ionization of gas collected on nearby wires, High-rate capability
- Multiple planes for 3-D track, High density electronic readout





# Quantum Idea for Charged Particle Monitoring/Tracking

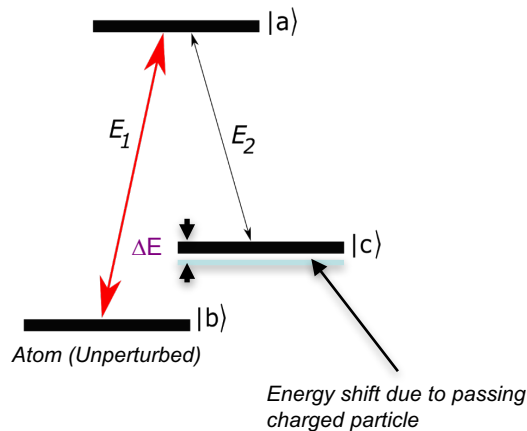
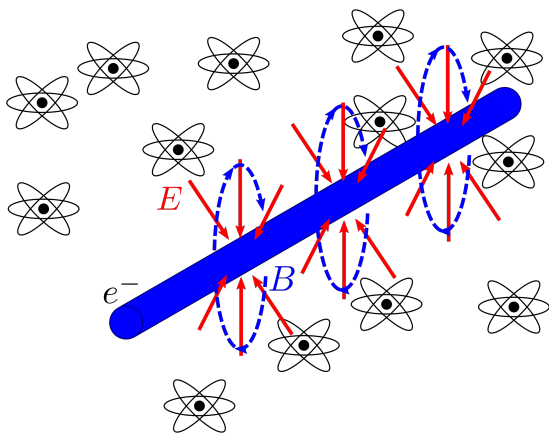
Charged particles produce electric/magnetic fields

The energy levels of Rb atoms are shifted, changing the atomic quantum state

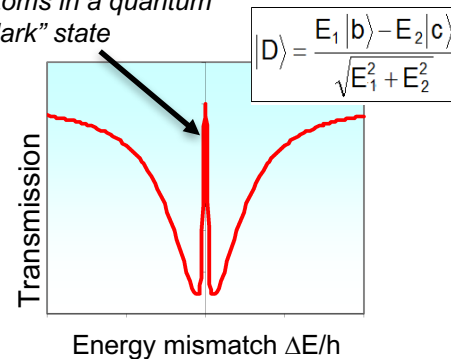
Particle track is recreated by detecting the altered optical properties of atoms via optical device

## Quantum Enhanced Tracker (QET)

Use phase sensitive superposition of atomic levels to greatly increase optical sensitivity to energy shifts.



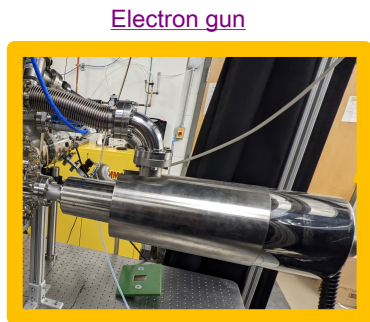
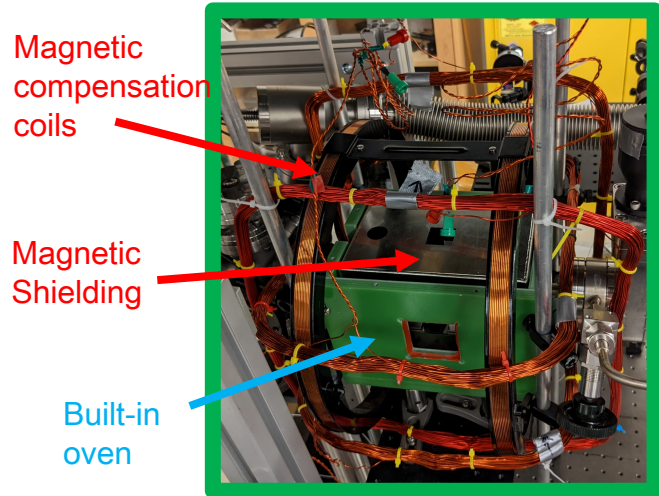
Atoms in a quantum "dark" state



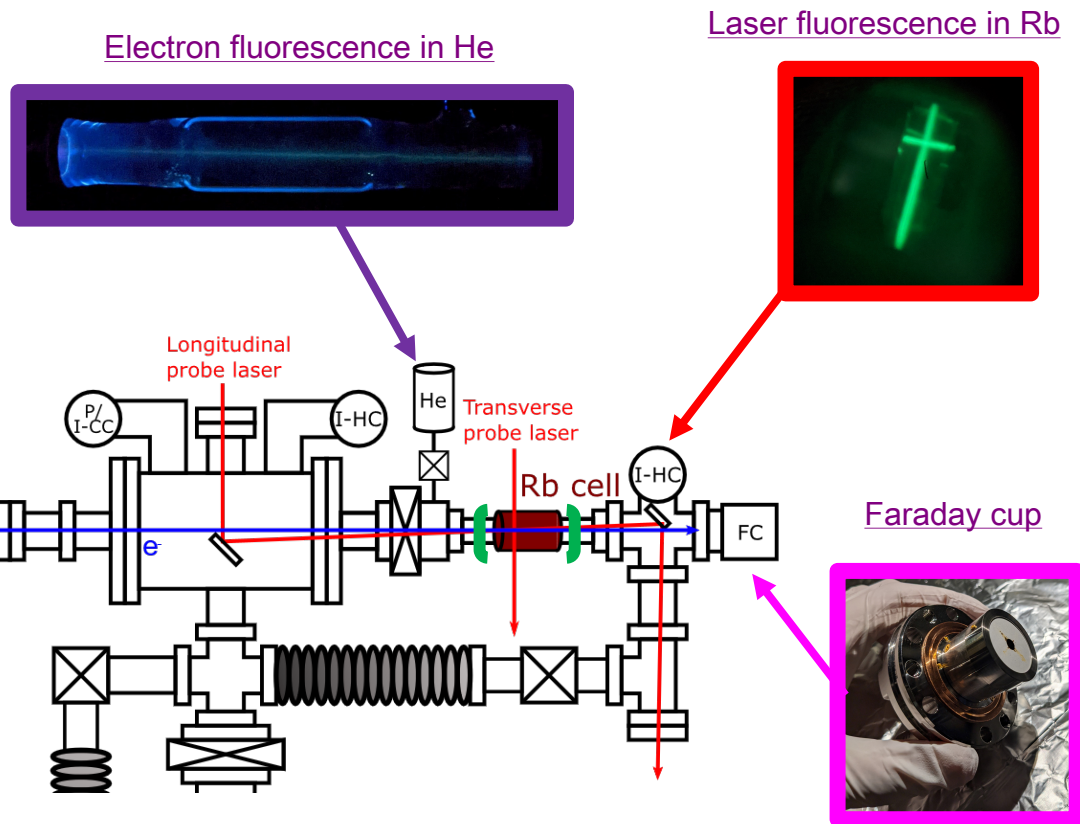
- Sensitive only to charged particles
- High Resolution, High Speed
- 3D tracking in single volume
- Small number of channels, high rate capability
- Modern 3D "bubble chamber"

An approach yet to be demonstrated!

# Experimental Setup

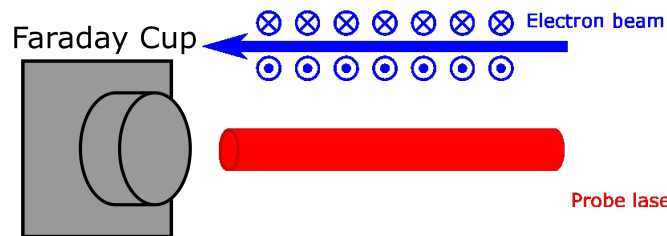


Electron gun  
30kV 200uA



# Longitudinal Detection

Longitudinal Polarization Rotation

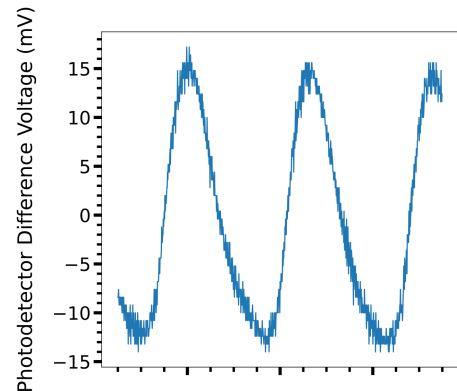


Probe laser

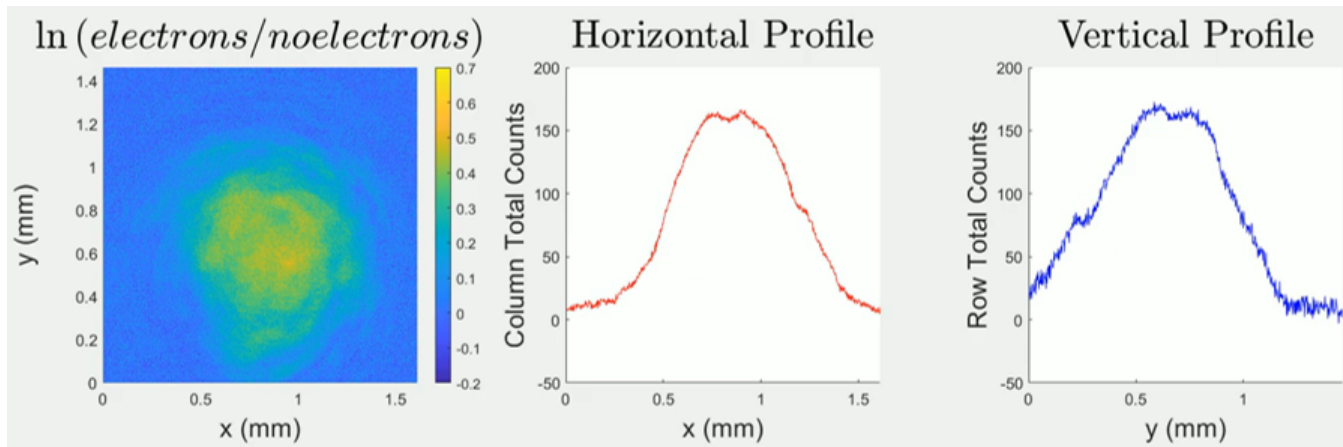
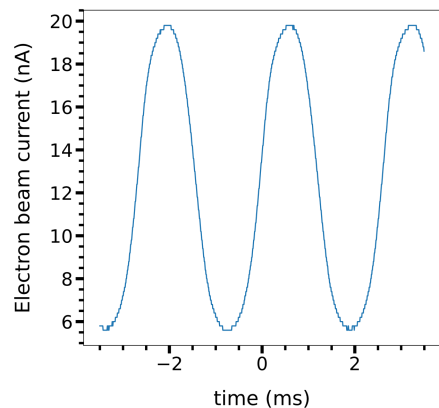
$$\sin(\phi) = \frac{1}{2} \frac{I_1 - I_2}{I_1 + I_2}$$

20 kV, 110  $\mu$ A,  $\sim 1$  mm $\phi$  electron beam parameters

Minimum detection current 10  $\mu$ A room-temperature cell



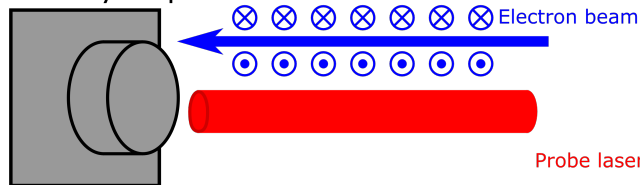
Electron Current



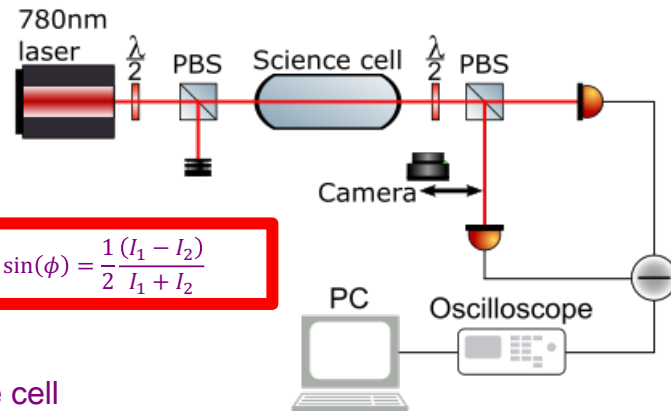
# Longitudinal Detection

Longitudinal Polarization Rotation

Faraday Cup



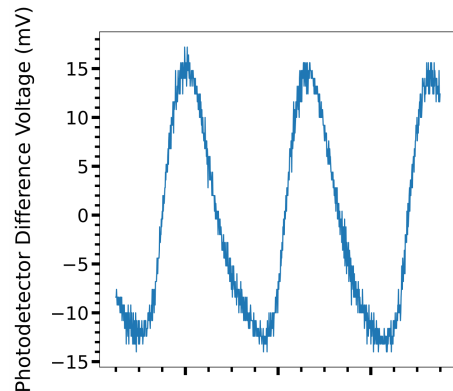
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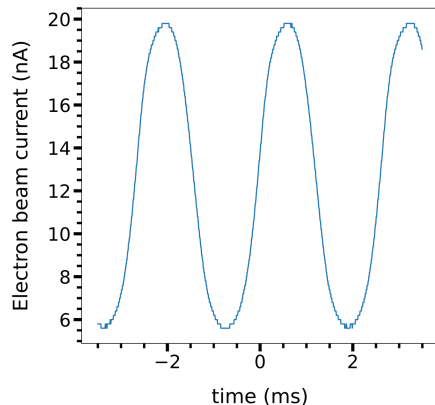
$$\sin(\phi) = \frac{1}{2} \frac{I_1 - I_2}{I_1 + I_2}$$

20 kV, 110  $\mu$ A,  $\sim 1$  mm $\phi$  electron beam parameters

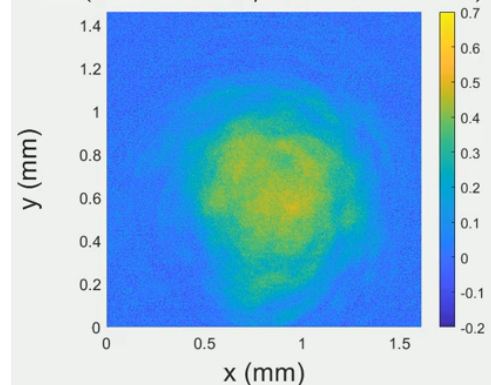
Minimum detection current 10  $\mu$ A room-temperature cell



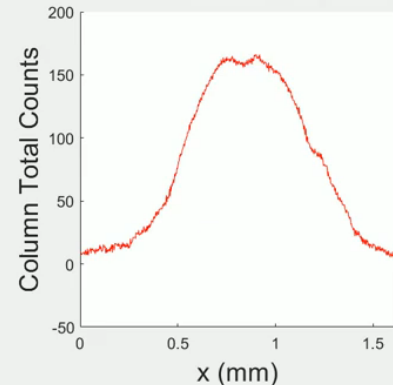
Electron Current



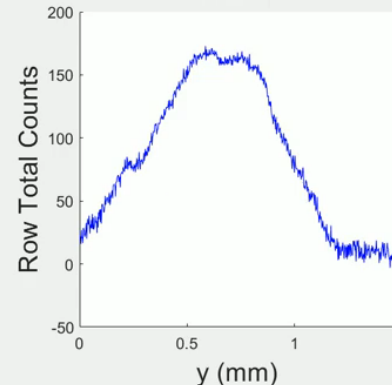
$\ln(\text{electrons/noelectrons})$



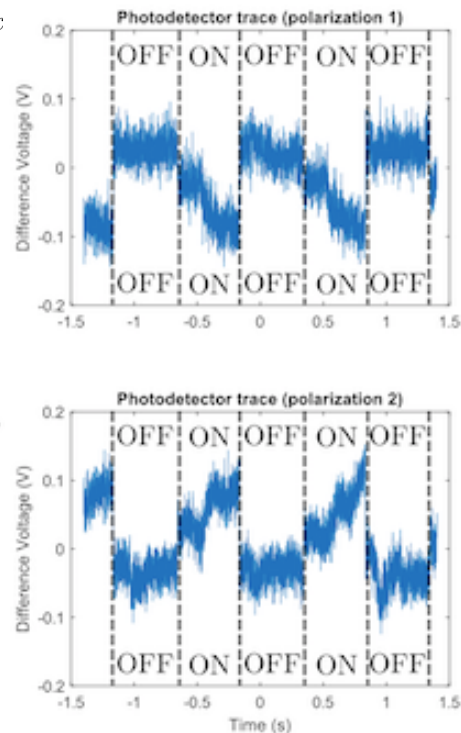
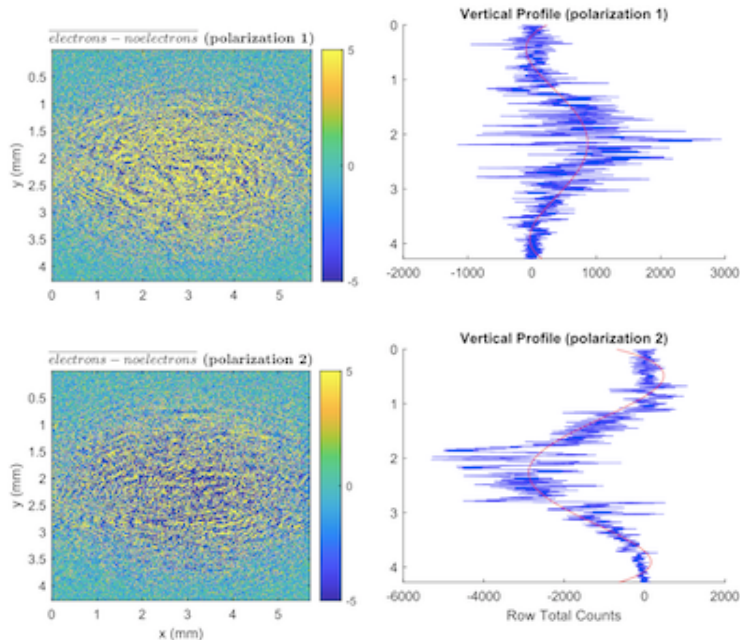
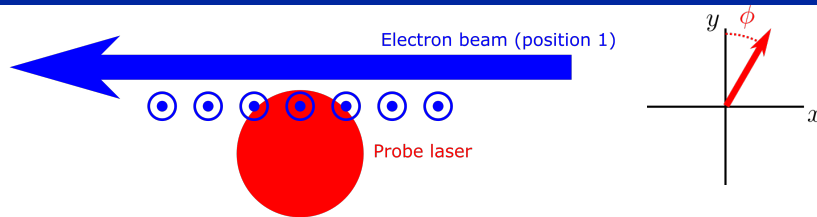
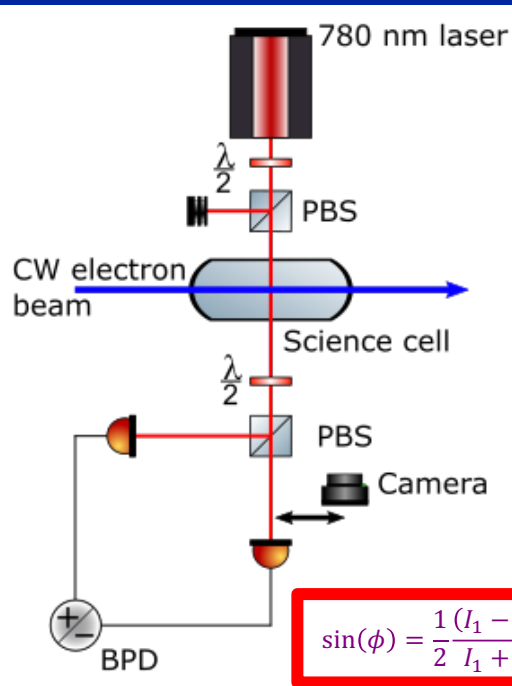
Horizontal Profile



Vertical Profile

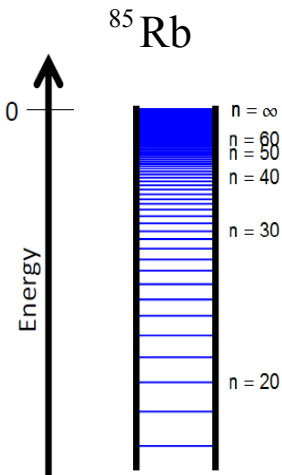


# Transverse Detection



- E-Source: 20 kV, 200  $\mu$ A,  $\sim 1$  mm $\phi$
- Optimal laser power in: 3 mW
- Minimum detection current 1  $\mu$ A 70°C cell

# Another Approach - Detection with Rydberg Atoms



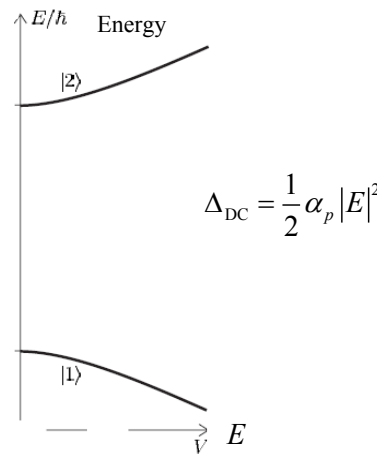
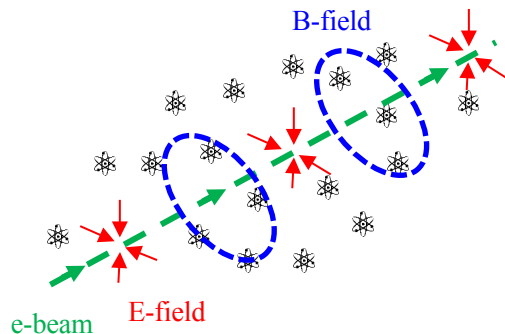
charged particles produce electric and magnetic

E-field shifts energy state of atoms (Stark effect)

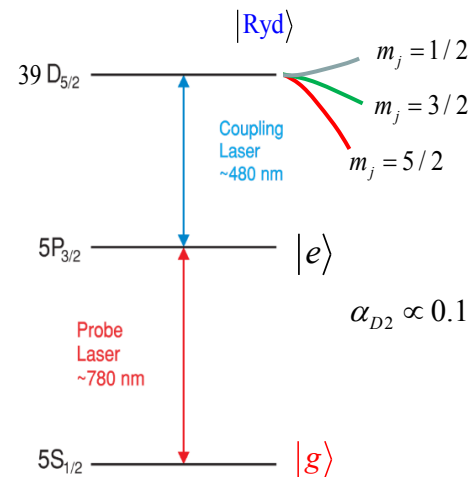
particles are detected by monitoring Rydberg atom

$$\alpha_p \propto n^7$$

$$\sim 40 \text{ (MHz/(V/cm)}^2\text{)}$$

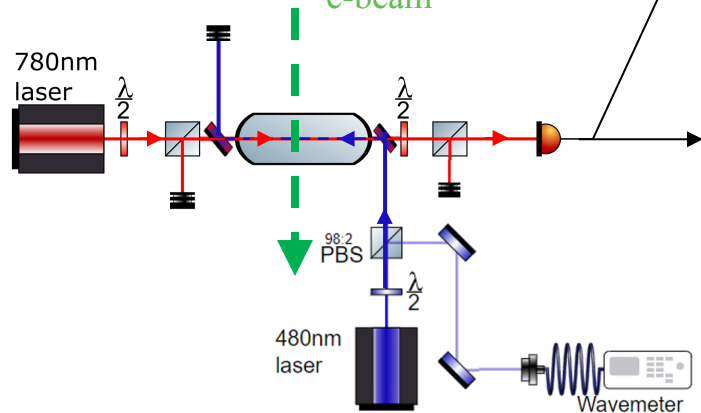
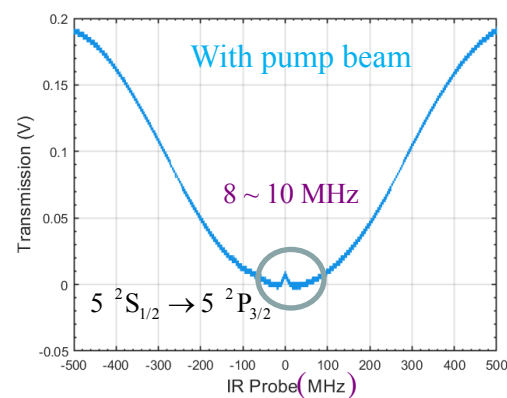
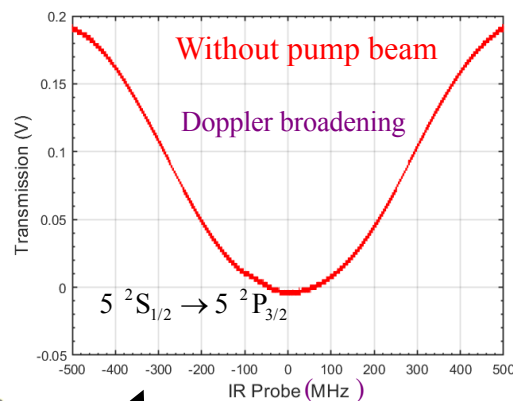
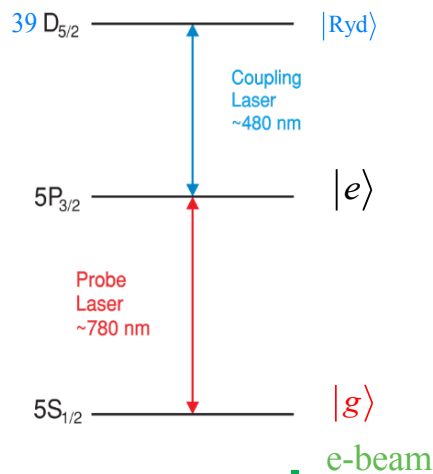


Presence of E-field

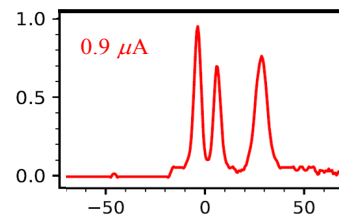
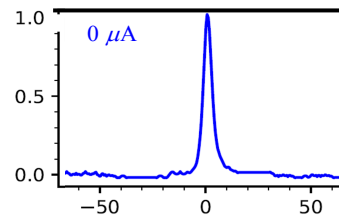


$$\alpha_{D2} \propto 0.13 \text{ Hz/(V/cm)}^2$$

# Rydberg EIT for e-Beam Sensor



IR Probe beam (A/U)



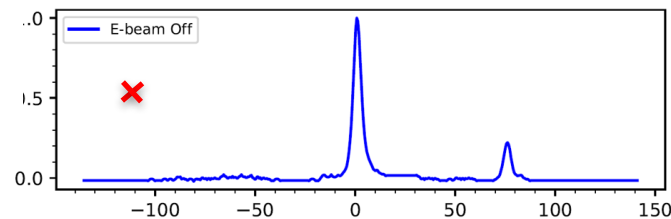
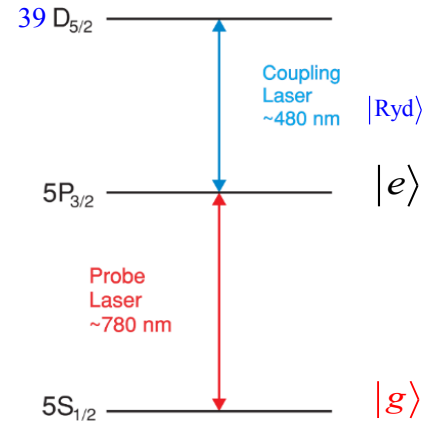
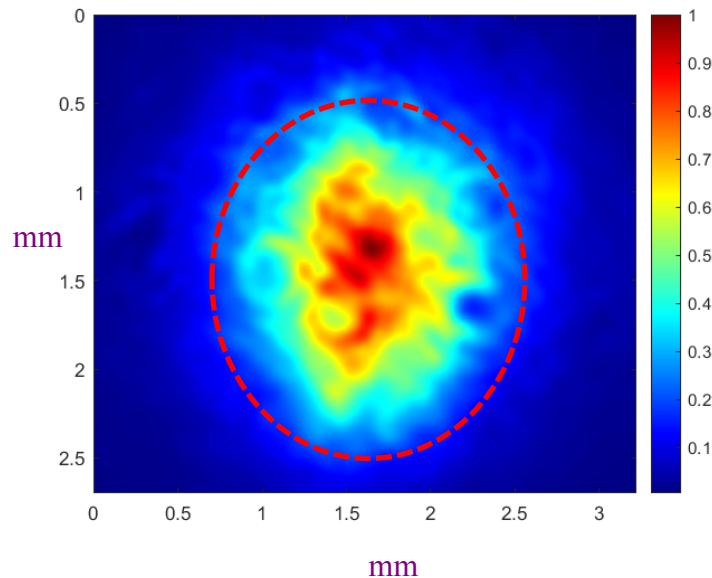
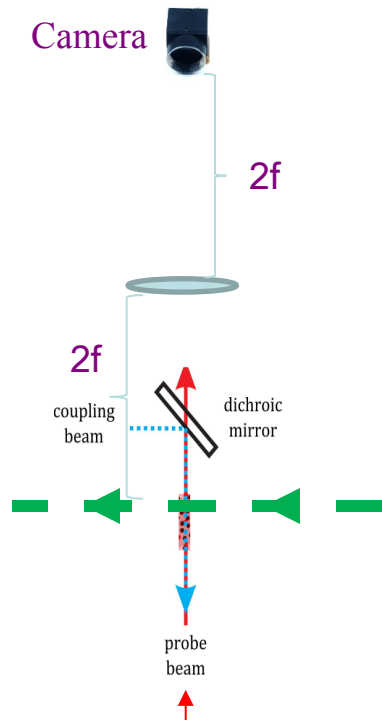
Pump beam detuning (MHz)

DC stark shift on Rydberg state due to e-beam



# Beam Imaging by EIT

$$\text{EIT} = \text{Probe [With Pump]} - \text{Probe [Without Pump]}$$



# Monitoring e-Beam Profile

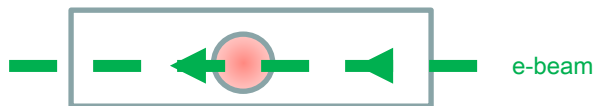
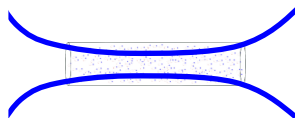
E-beam setup/focus 1

FWHM  $\sim 1$  mm

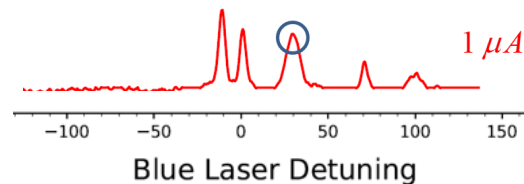
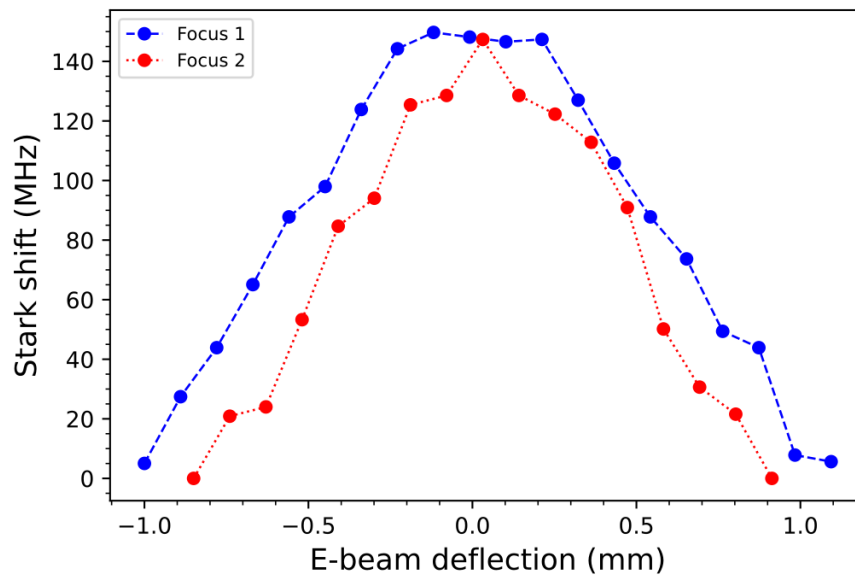


E-beam setup/focus 2

FWHM  $\sim 1.4$  mm

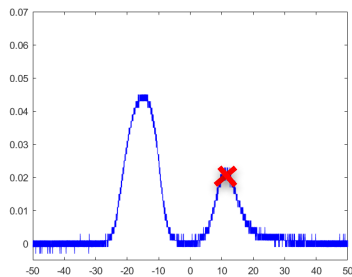
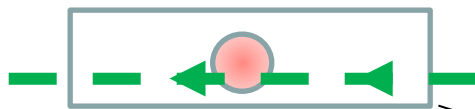
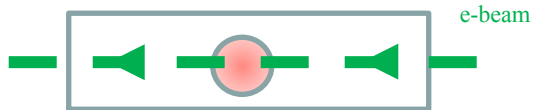


Extract e-beam width through DC Stark shift



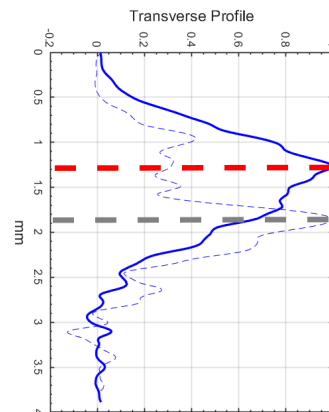
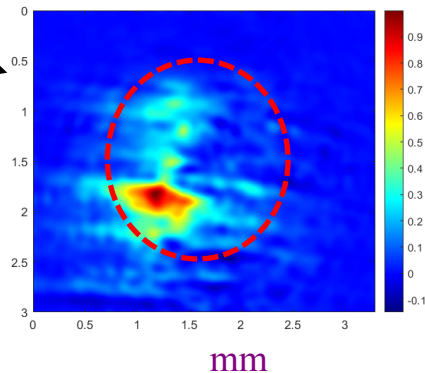
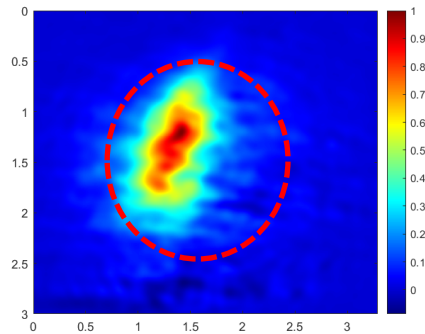
# Imaging electron beam

Transverse Cross Section

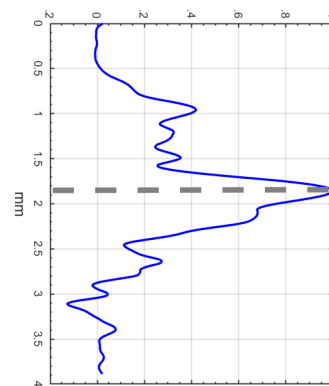


Pump laser ( MHz)

$I_{\text{E-beam}} - I_{\text{No E-beam}}$



$\Delta Y_{\text{measured}} = 0.58 \text{ mm}$

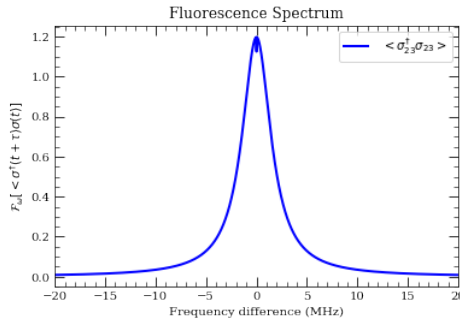


- Still under investigation, charging on cell wall may have affected the result!

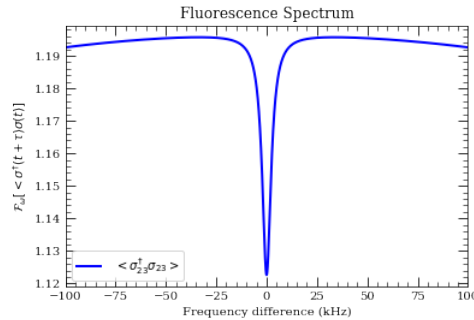
# Semi-classical Theory Study

## Demonstration of the effect of relativistic charged particles on EIT measurements

- AC Stark shifts due to fields from relativistic charged particles calculated using semiclassical formalism with the minimal coupling Schrodinger equation.
- Effects of interest include phase shift to EIT wavefunction and the scattering spectrum including coherent and incoherent fluorescent effects.



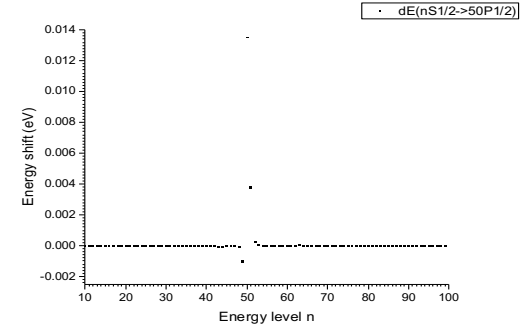
$\Omega_p = 0.1$  MHz,  $\Omega_S = 10$  MHz,  $\Delta_1 = \Delta_2 = 0$



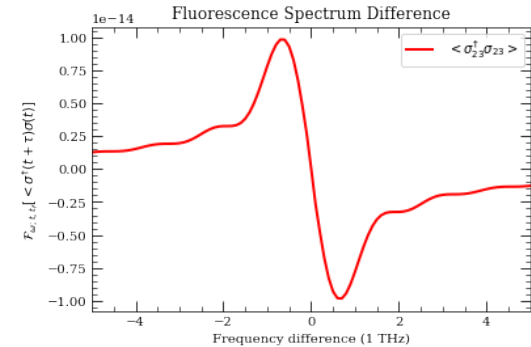
$\Omega_p = 0.1$  MHz,  $\Omega_S = 10$  MHz,  $\Delta_1 = \Delta_2 = 0$

The EIT steady state fluorescence spectrum from a driven single atom is proportional to the spectral transform of the two-time correlation function for the atomic dipole operator  $g_{\lambda}^{(1)}(\vec{r}', t'; \vec{r}, t) = \left\langle \left( \sum_i \sigma_i^{\dagger} \left( \vec{r}', t' - \frac{|\vec{r}' - \vec{r}_i|}{c} \right) \right) \left( \sum_j \sigma_j \left( \vec{r}, t - \frac{|\vec{r} - \vec{r}_j|}{c} \right) \right) \right\rangle$ .

External perturbations will introduce Stark shifts and change the spectrum for the duration where the perturbation is active. A short-term Fourier transform of  $g_{\lambda}^{(1)}(\vec{r}', t'; \vec{r}, t)$  is used to evaluate the perturbed spectrum



Lowest order perturbations for the energy of the  $n=50$  Rydberg  $P1/2$  state from  $S$ -state virtual transitions due to Coulomb field of relativistic charged particle with  $\beta = 0.999$  at  $1 \mu\text{m}$  distance perpendicular to motional axis. No dipole approximation and inhomogeneity of field is considered.



$\Omega_p = 0.1$  MHz,  $\Omega_S = 10$  MHz,  $\Delta_1 = \Delta_2 = 0$ . Introduced gaussian DC Pulse with duration  $\tau = 10$  fs causing effective AC Stark shift 20.0 MHz for the Rydberg level, 100 kHz for other levels, .

# Summary & Future Perspectives

## Studied e-beam detection by two approaches

- *NL magneto-optical rotation* to measure magnetic fields
  - Beam sensed in both longitudinal and transverse configurations
  - Beam image captured on camera via polarization rotation
  - Minimum detection of  $1\ \mu A$  particle beam
- Rydberg atoms attempted for high sensitivity detection, under further study

## Theoretical demonstration of the effect of relativistic charged particles on EIT measurements

### In Future:

- Study across wider energies range, including *relativistic electron* source
- Demonstrate *3D-imaging* capability
- Improve/Optimize system/explore *beam halo & Single-particle detection*

**Acknowledgement:** This material is based upon work supported by the U.S. Department of Energy, Office of Science, Office of Nuclear Physics under contract DE-AC05-06OR23177. This project is supported by JLAB LDRD program and College of W&M.