



Elettra Sincrotrone Trieste

# Amplified Emission of a Soft-X Ray Free-Electron Laser Based on Echo-Enabled Harmonic Generation

*E. Allaria*  
on behalf of the **FERMI** team  
and **EEHG** collaboration

## FERMI Free-Electron Laser

- Spectral and temporal control
- Coherence and coherent control

## EEHG at FERMI

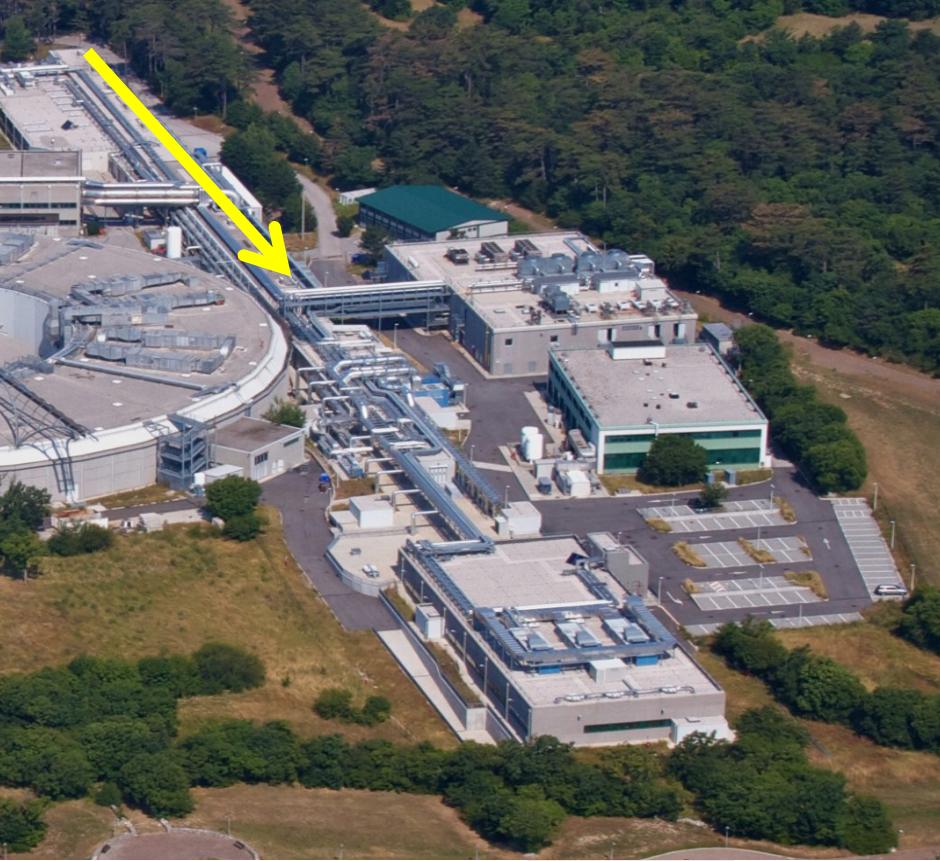
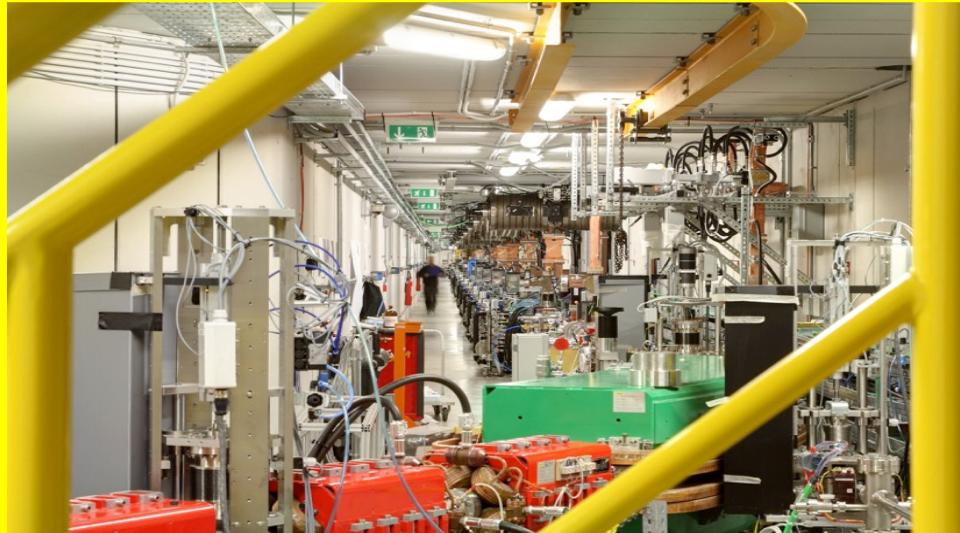
- Evidence of EEHG
- FEL gain
- Spectral quality

## Conclusions

# FERMI Free Electron Laser

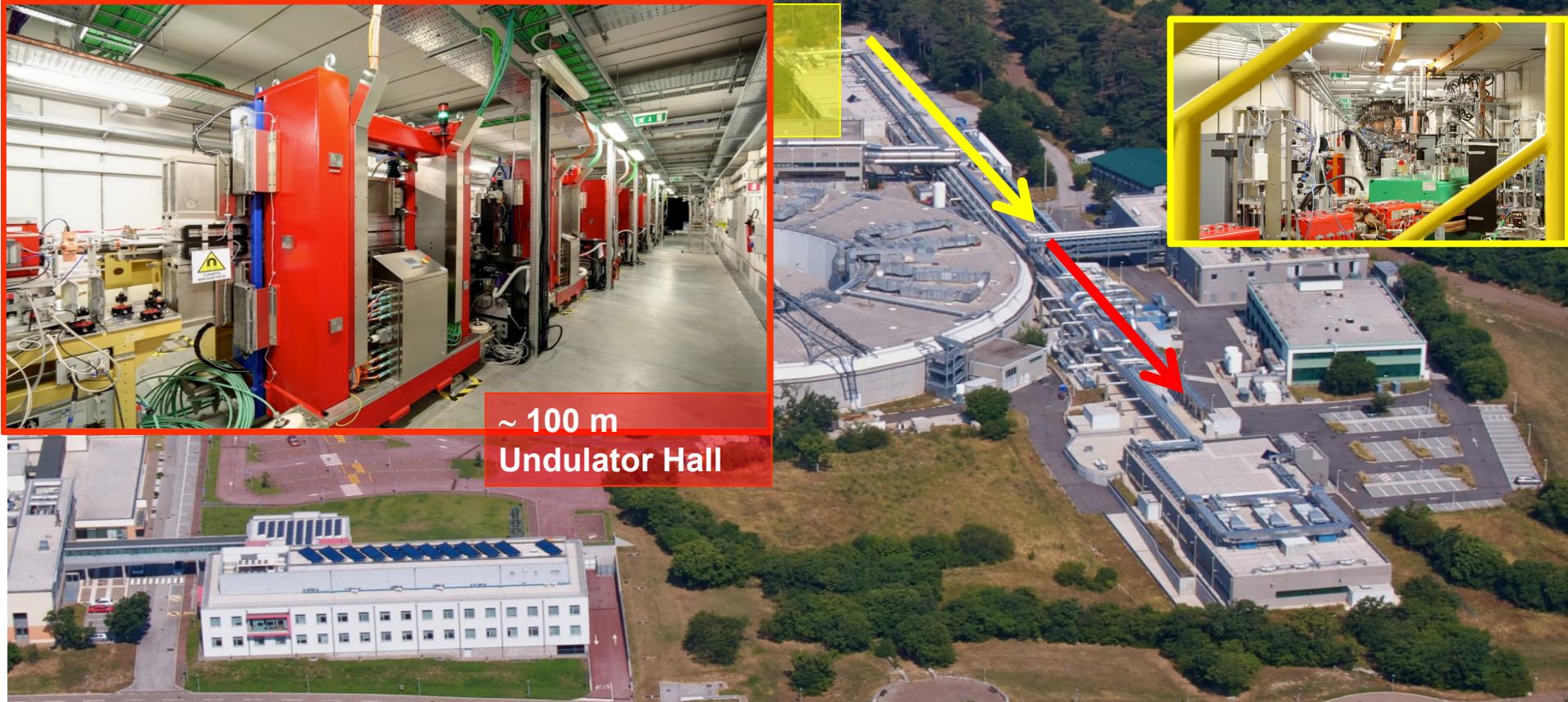


# FERMI Free Electron Laser

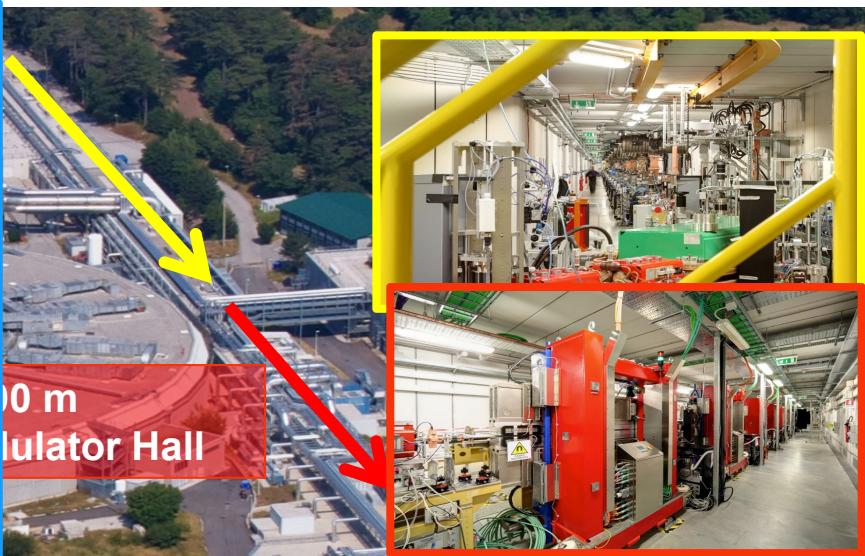


Linac Tunnel +  
Injector Extension

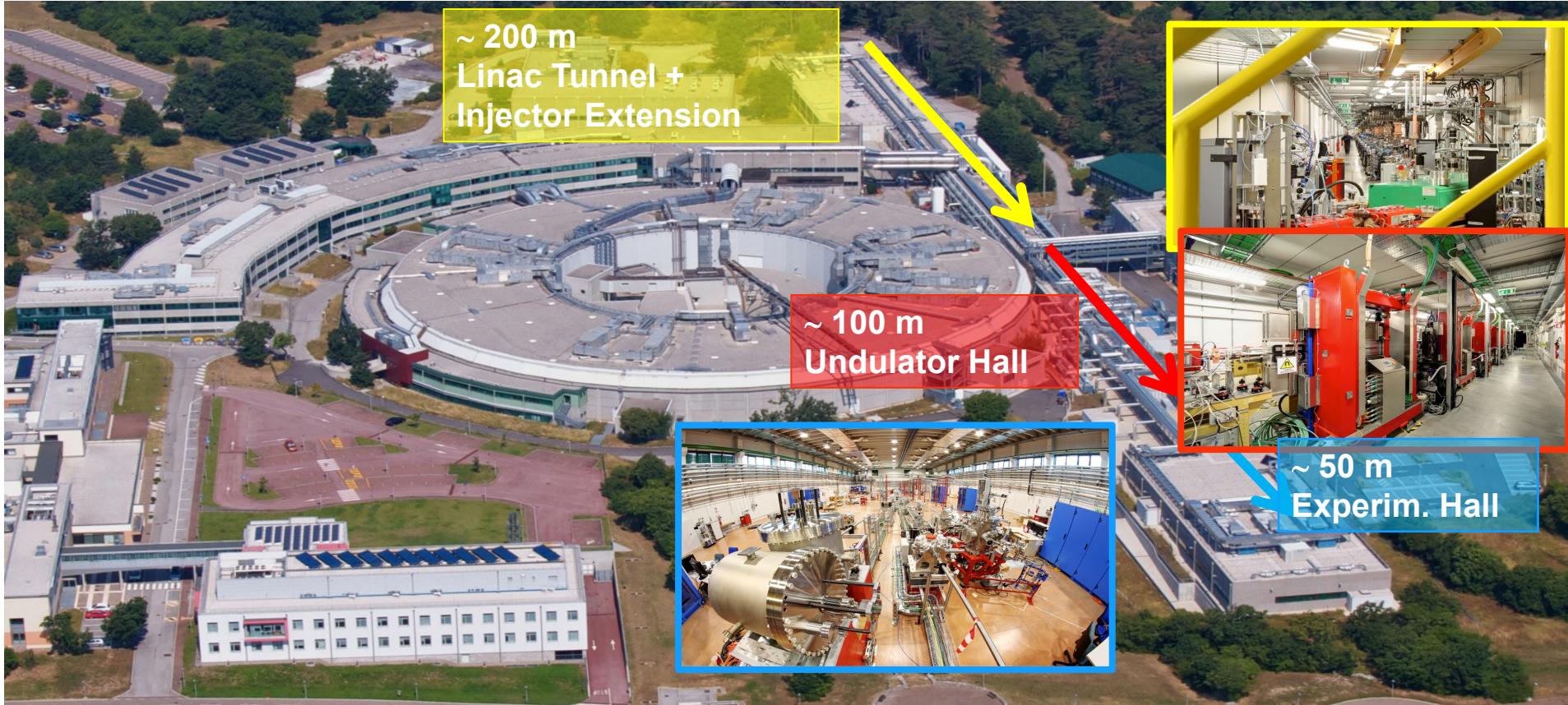
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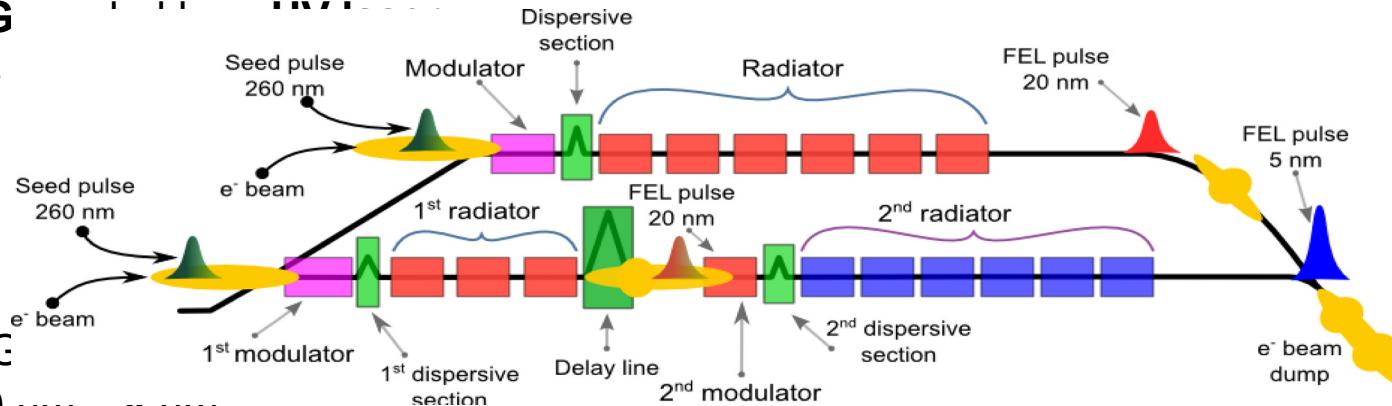


# FERMI Free Electron Laser



# FERMI FELs: FEL-1 & FEL-2

**FEL-1: single stage HGHG**  
covers the range **100 nm** –



**FEL-2: double cascade HG**  
to reach the wavelength **20 nm** –

**FEL-1 (Nat. Photon. 6, 699 (2012))**

Tuning range	100-20 nm (12-60eV)
Relative bandwidth	$1 \times 10^{-3}$ (FWHM)
Pulse length	<100 fs
Pulse energy	20-100 $\mu$ J

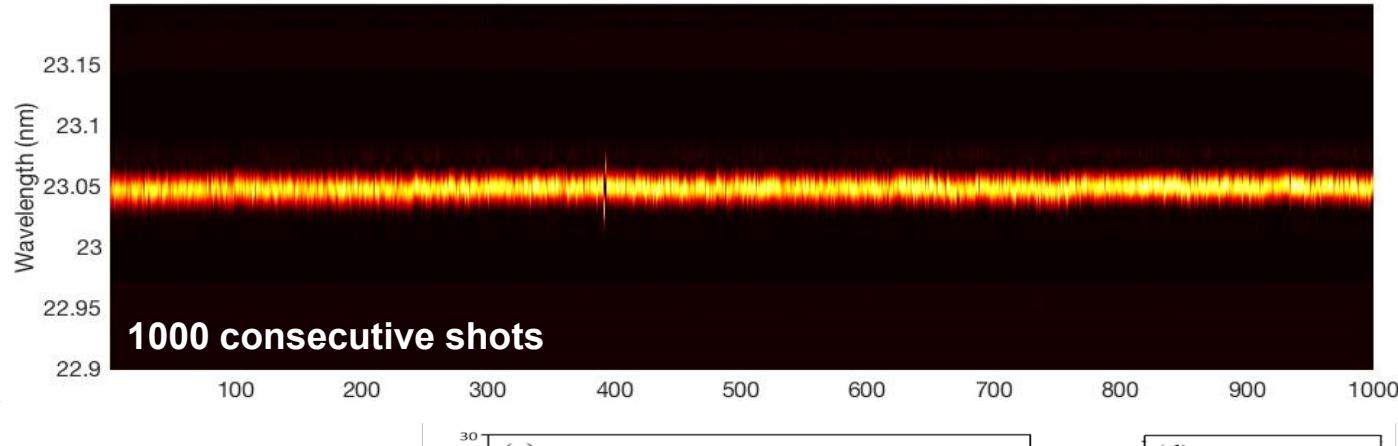
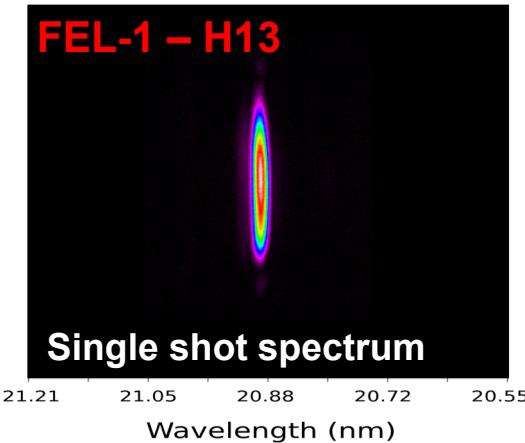
**FEL-2 (Nat. Photon. 7, 913 (2013), Journal of Synchrotron Radiation 22 (2015))**

Tuning range	20-4 nm (60-300eV)
Relative bandwidth	$1 \times 10^{-3}$ (FWHM)
Pulse length	~50 fs
Pulse energy	10-70 $\mu$ J

Both FELs have APPLE-II undulators in the final radiator allowing **polarization control**.

# Spectral and temporal properties

Single spectral mode up to **H13** on **FEL1** and **H65** on **FEL2** (*linewidth*  $\sim 2 \cdot 10^{-4}$  rms, depending on wavelength & seed). High wavelength stability set by the seed ( $10^{-5}$  rms)

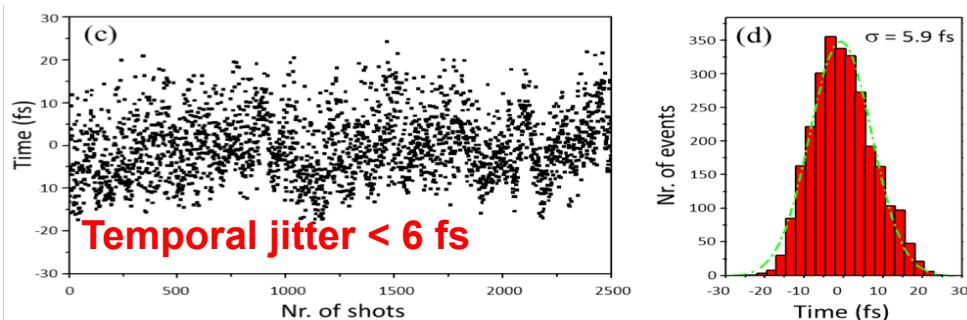


Thanks to the the seeding, the temporal jitter of the FEL with the external laser is very low.

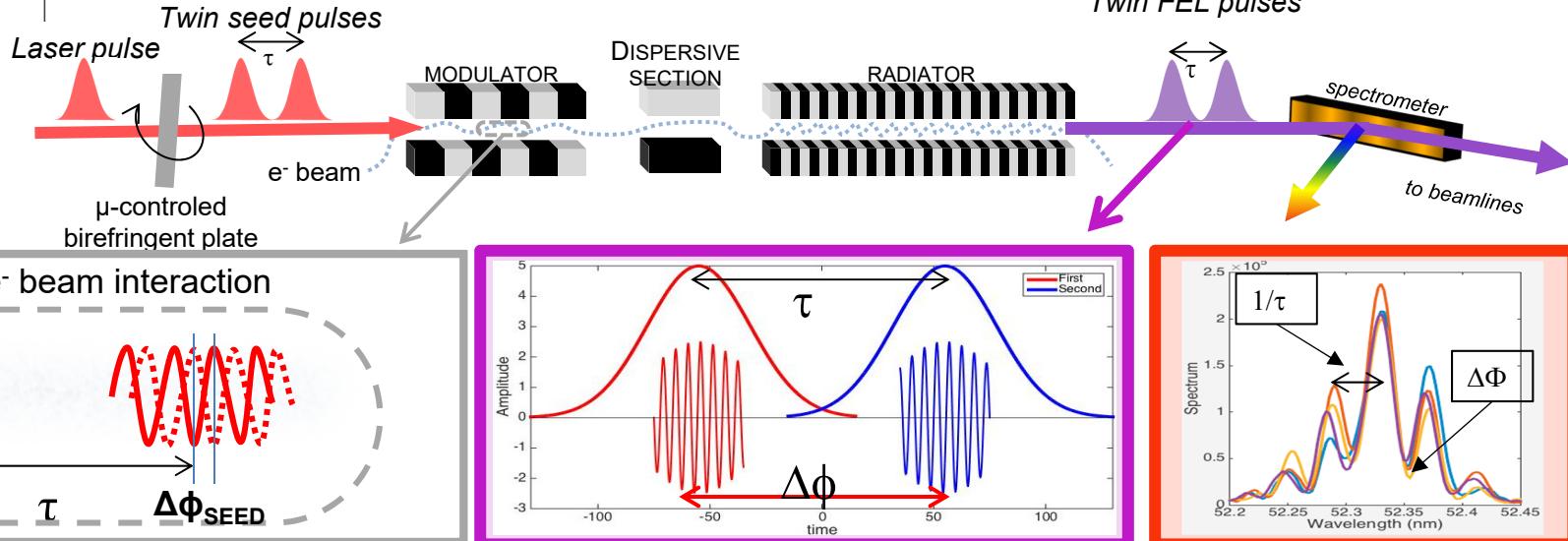
M. Danailov *et al.* Optics Express, Vol. 22, Issue 11, 12869 (2014)



10<sup>th</sup> International Particle Accelerator Conference



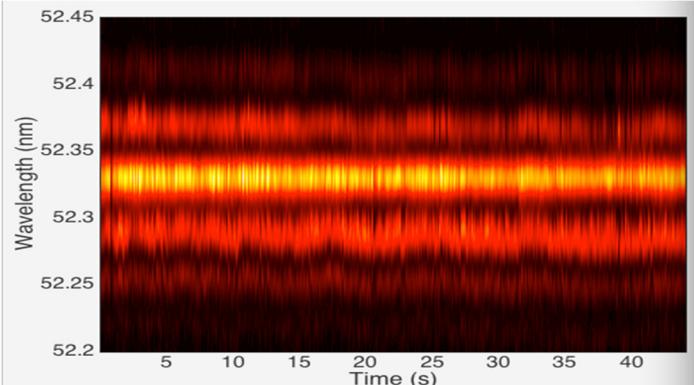
# Phase-locked FEL pulses



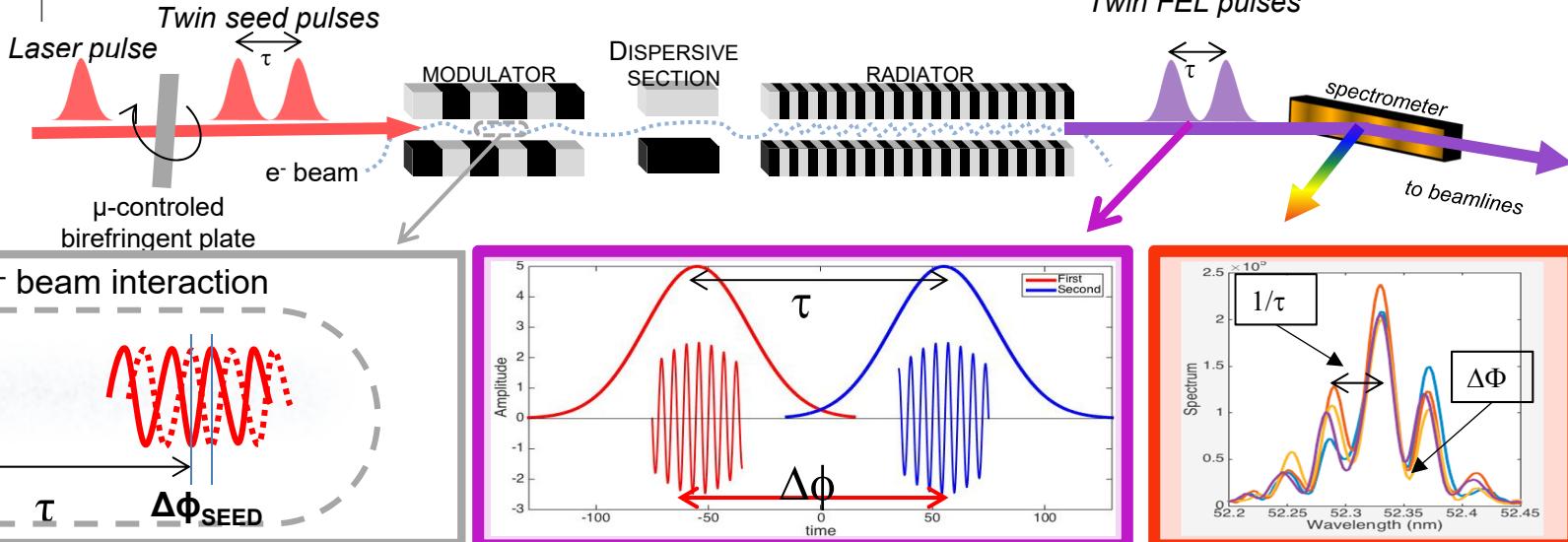
Using two **seed** lasers we can control and change the **relative time** between two **FEL** pulses. For coherent pulses, a fine tuning **control** the **relative phase** between the two **FEL pulses**.

**Interference** between two **coherent** and **phase-locked** pulses is evident in the spectral domain.

D. Gauthier et al., Phys. Rev. Lett. **116**, 024801 (2016)



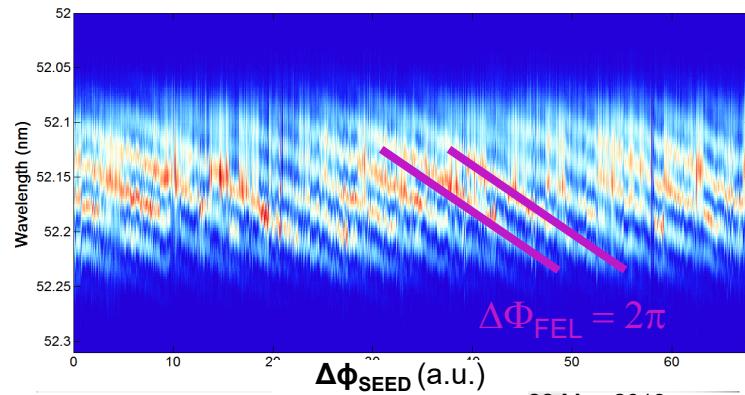
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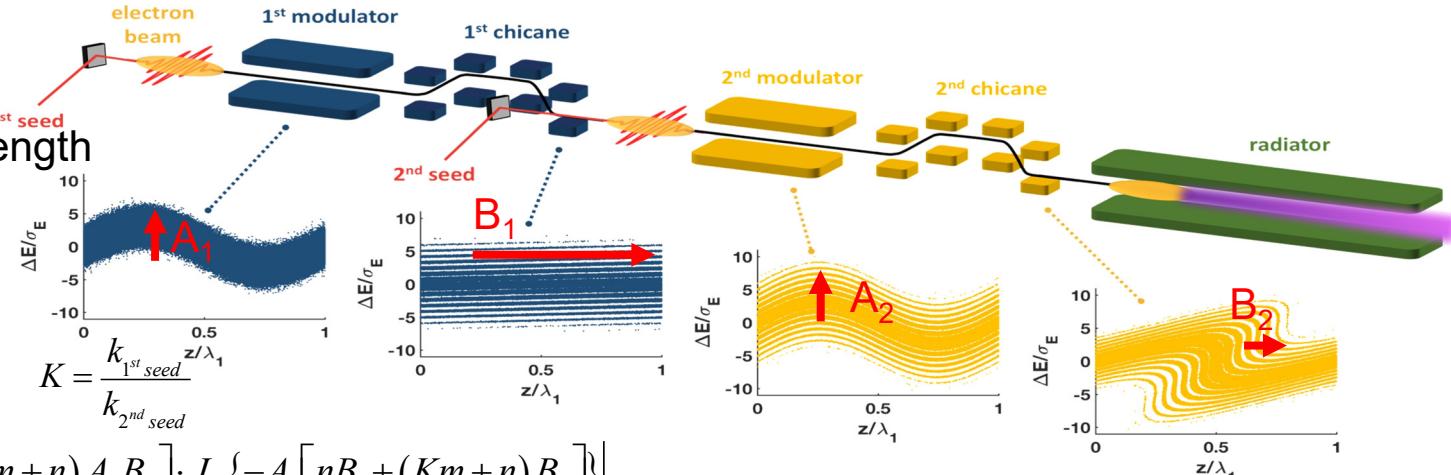
# EEHG scheme

G. Stupakov, PRL, 2009

A new scheme allows reaching FEL-2 wavelength with a single stage harmonic generation seeded in the UV.

$$k_{EEHG} = n \cdot k_{1^{st} seed} + m \cdot k_{2^{nd} seed}$$

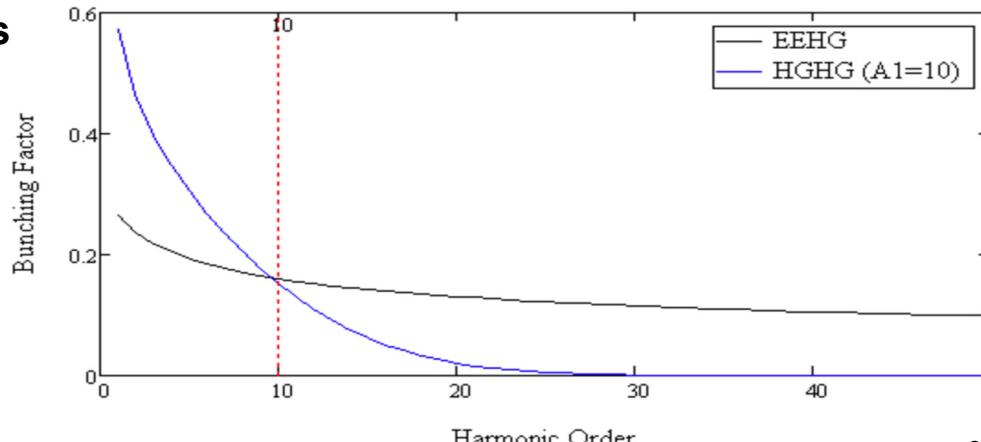
$$b_{n,m} = \left| e^{\frac{j}{2} [nB_1 + (Km+n)B_2]} \cdot J_m [-(Km+n)A_2 B_2] \cdot J_n \{ -A_1 [nB_1 + (Km+n)B_2] \} \right|$$



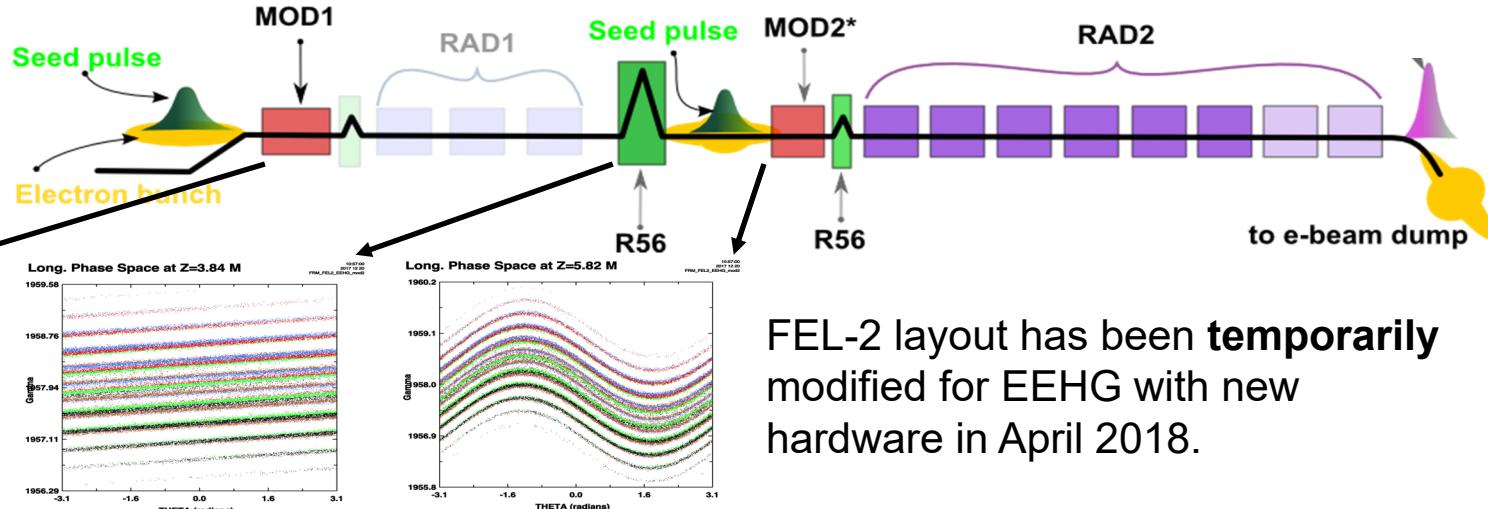
**EEHG bunching at the desired harmonic depends on 4 parameters and slowly decreases with harmonic number.**

Bunching is maximized for  $B_2 \approx \frac{|n|}{H} B_1$ .

Higher bunching requires  $n = -1$   
and slightly decreases for  $n = -2, -3$



# FEL-2: from HGHG-FB to EEHG

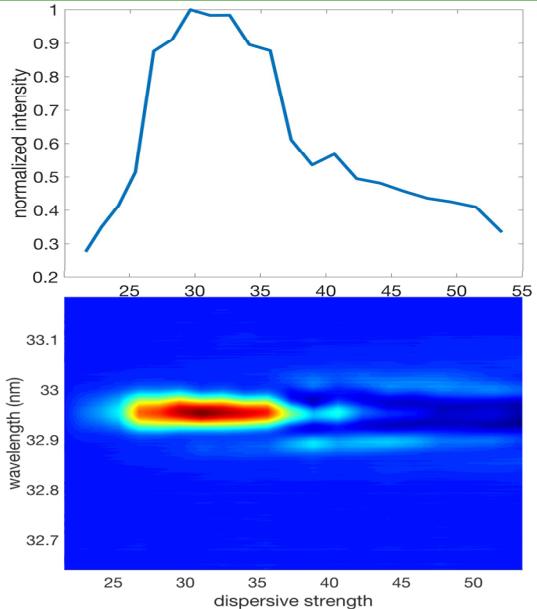


FEL-2 layout has been **temporarily** modified for EEHG with new hardware in April 2018.

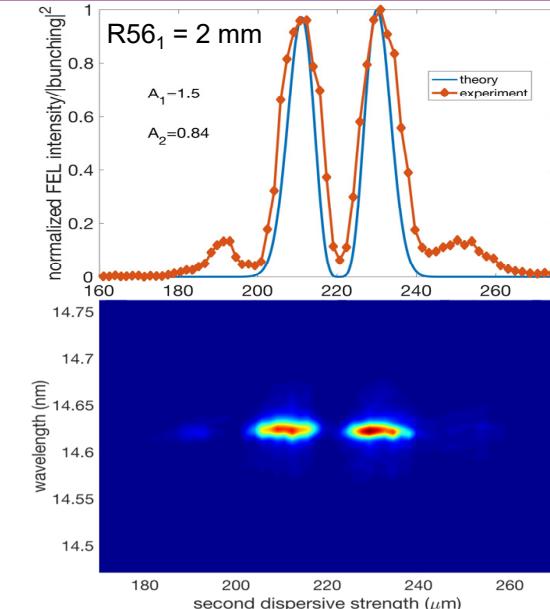
- Delay line:** changed position of the magnets and new power supply for increasing R56.
- Second modulator:** changed the undulator.
- Seed laser:** new seed laser for EEHG.
- Diagnostic:** new optical table for seed and electron diagnostic.
- Photon diagnostic:** VUV CCD on PRESTO

# Signature of EEHG

**HGHG**



**EEHG**



FEL intensity and spectra show **different response** to the dispersive strength  $R_{56}$  in **HGHG** and **EEHG**.

**HGHG** presents **one maximum**; for larger  $R_{56}$ , power slowly **decreases** and **spectrum splits** as a result of the over-bunching.

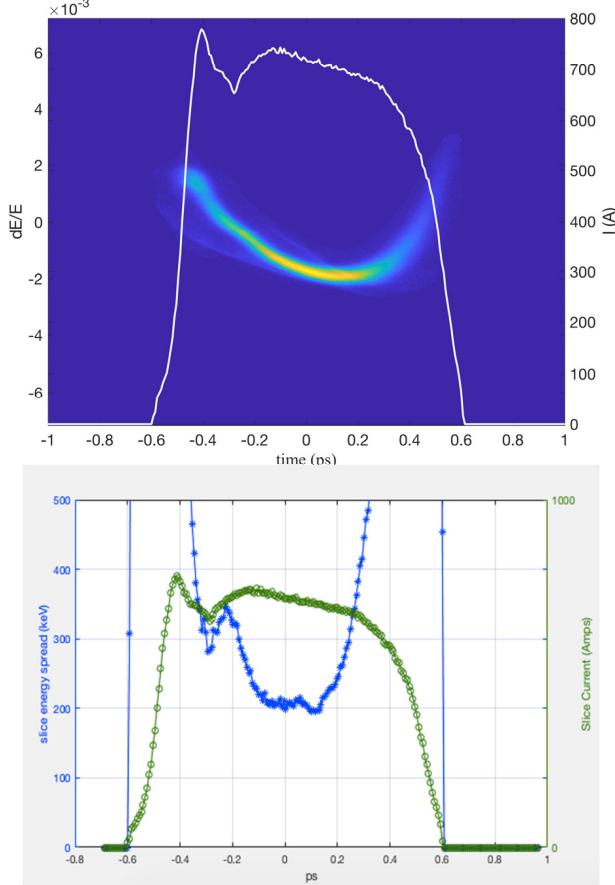
**EEHG** is maximized in **two very narrow regions** determined by the bunching equations.

For this case at  $H = 36$  optimized for  $n=2$ , a minimum is

found for  $R_{56_2} \equiv \frac{2}{18} R_{56_1}$  with two adjacent maxima.



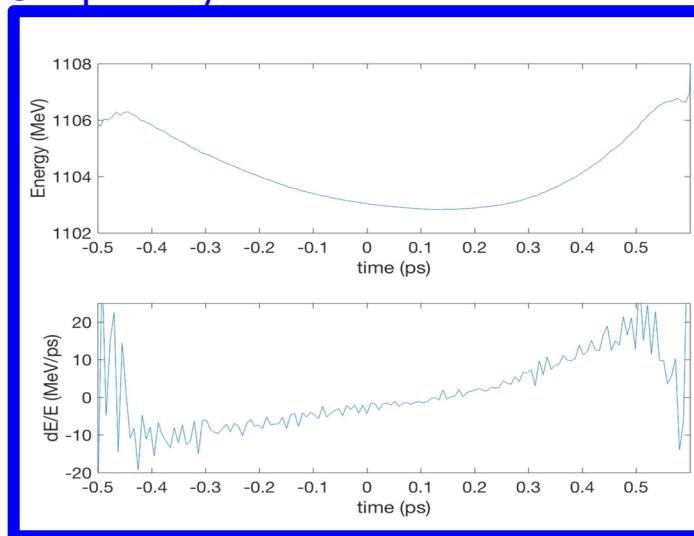
# E-beam phase space at E = 1.1 GeV



The **same** electron beam has been used for **both EEHG and HGHG** to investigate the sensitivity to electron beam properties.

E-beam properties are measured at the end of the linac.

**Chirp analysis:**

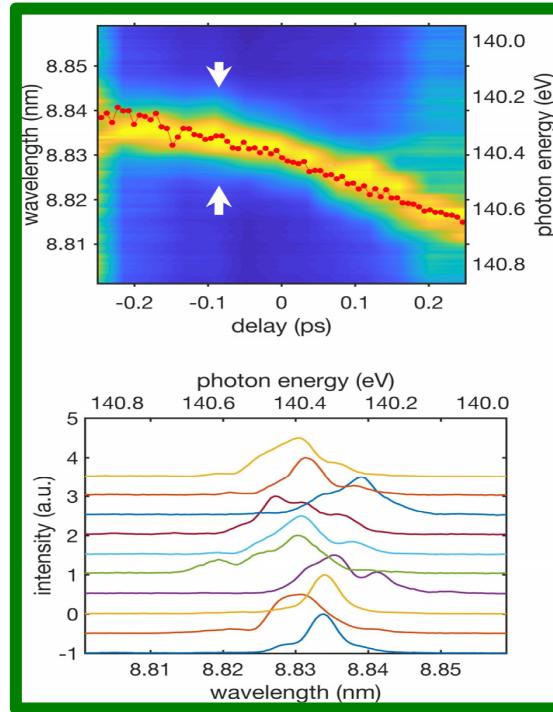
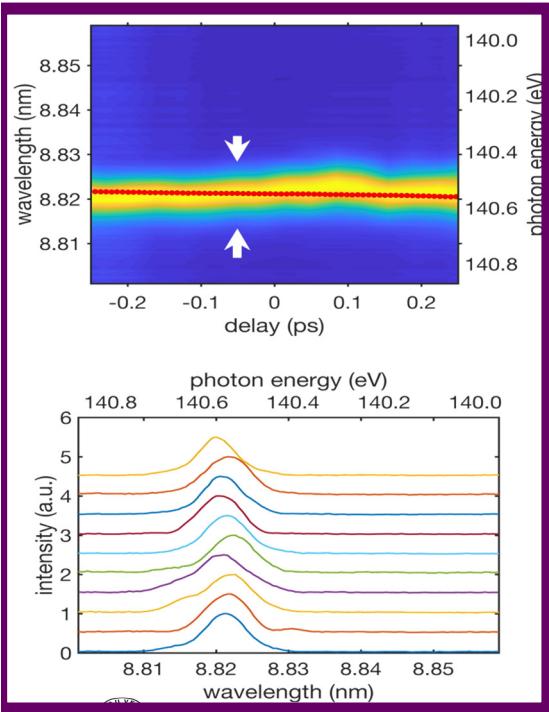


A not negligible **quadratic chirp** characterizes the phase space of FERMI electron beam.

# FEL behavior vs. e-beam properties

From bunching equation one gets that **FEL wavelength is affected by e-beam energy chirp** according to:  $\frac{d\lambda}{\lambda} \approx \frac{1}{E} \left( R_{56_2} + \frac{n}{H} R_{56_1} \right) \frac{dE}{dz}$

**HGHG**



The value of  $R_{56_2}$  is **comparable** for **EEHG** and **HGHG**.

**EEHG** takes **advantage** of the **combined** effects of **two R56**.

Maximum EEHG bunching when  $R_{56_2} \approx \frac{|n|}{H} R_{56_1}$  with negative  $n$ .

**Sensitivity to e-beam phase space almost cancelled.**

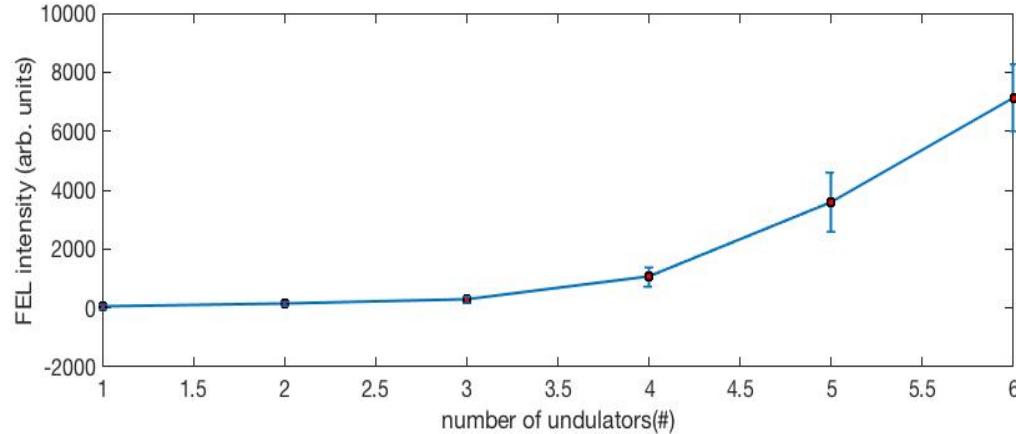
Much more stable FEL central wavelength and spectra with **less structures**.

# Amplification of EEHG

**EEHG bunching** at high harmonics limited to **less than 10%**.

Generation of **usable** radiation in the soft-X ray region critically **depends** on the **exponential growth** of FEL radiation in the **long radiator**.

For the first time **amplification** of EEHG has been measured and characterized down to **~5nm**.

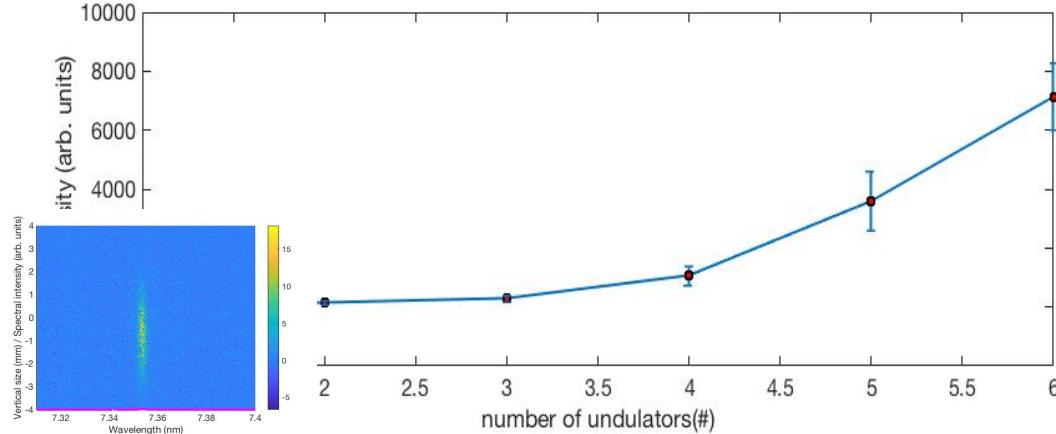


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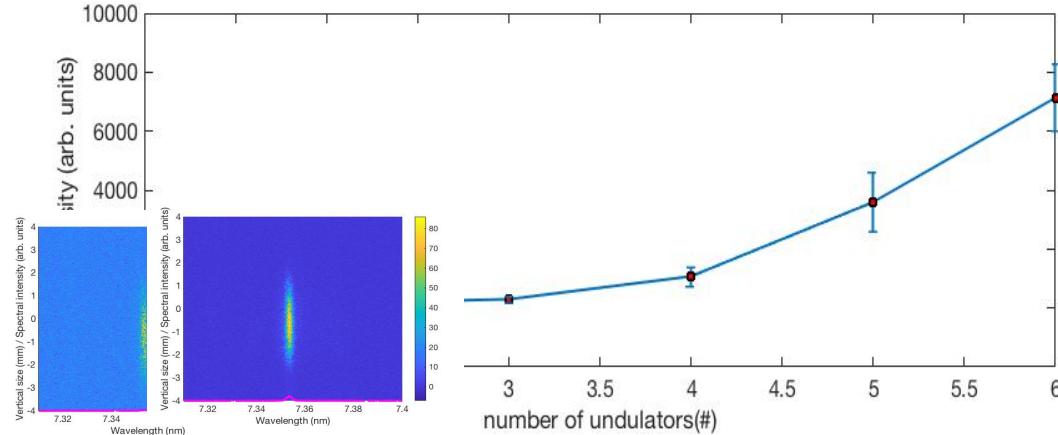


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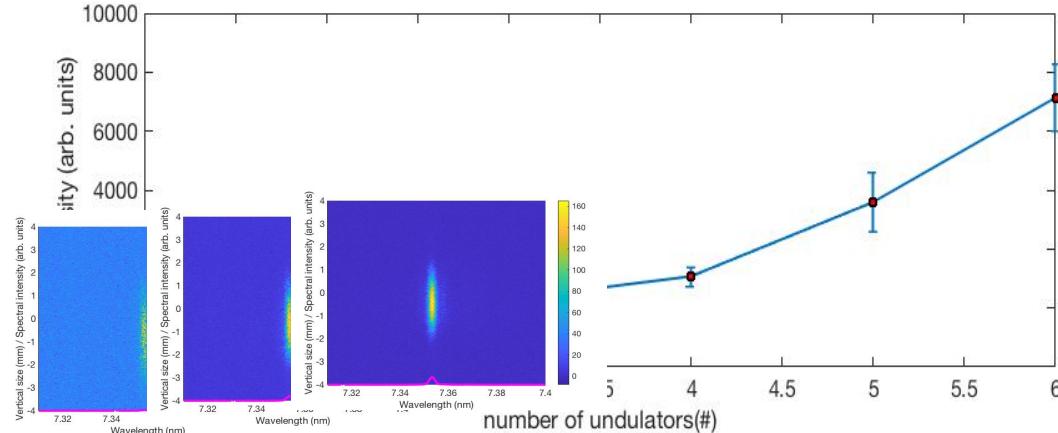


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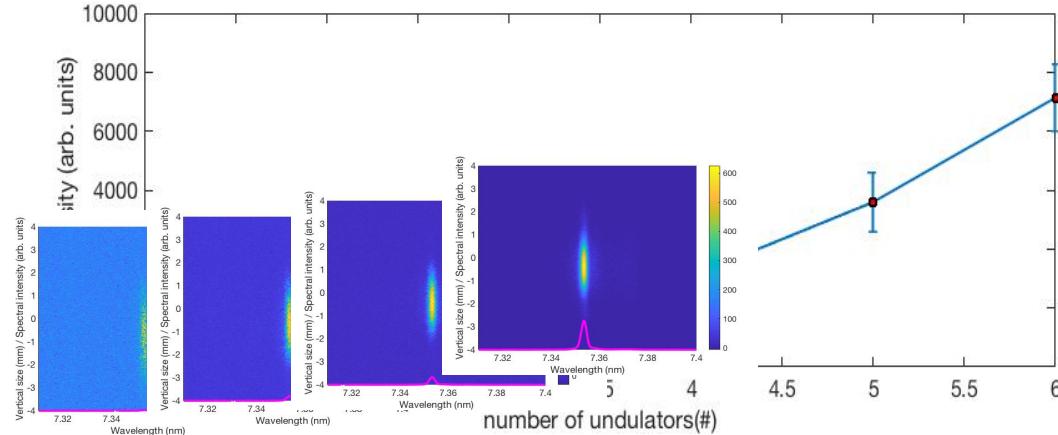


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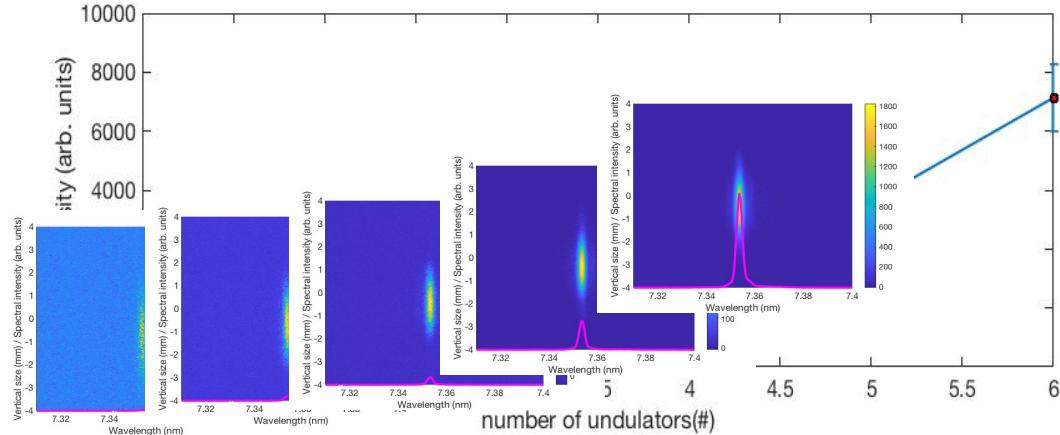


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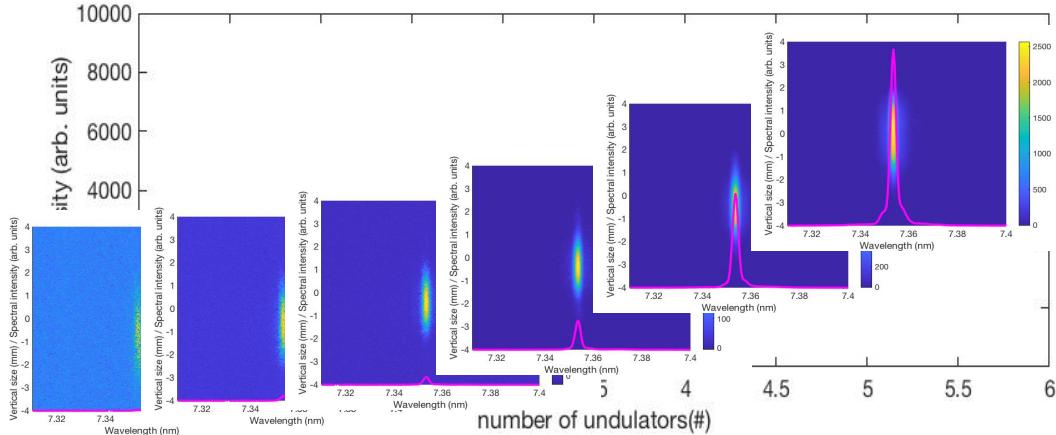


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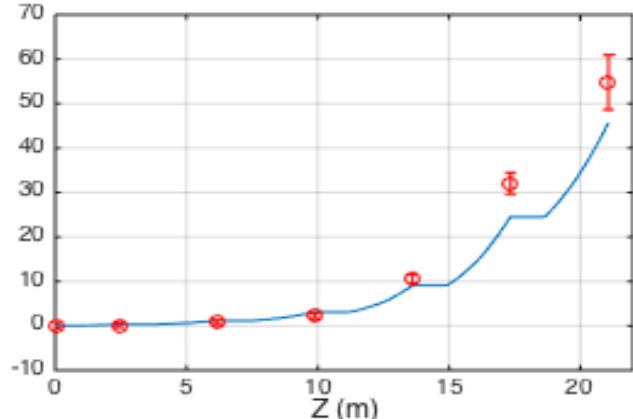
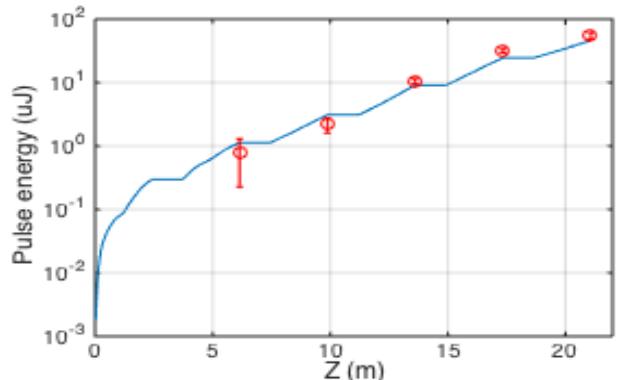
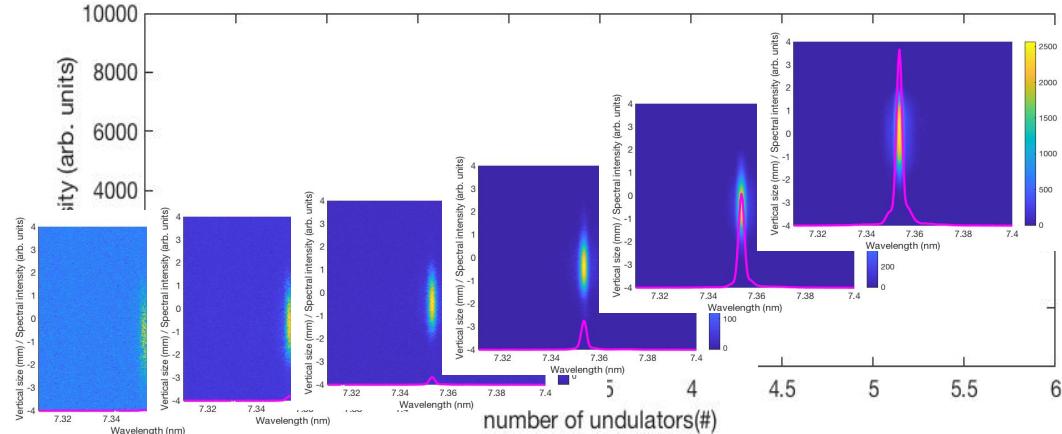
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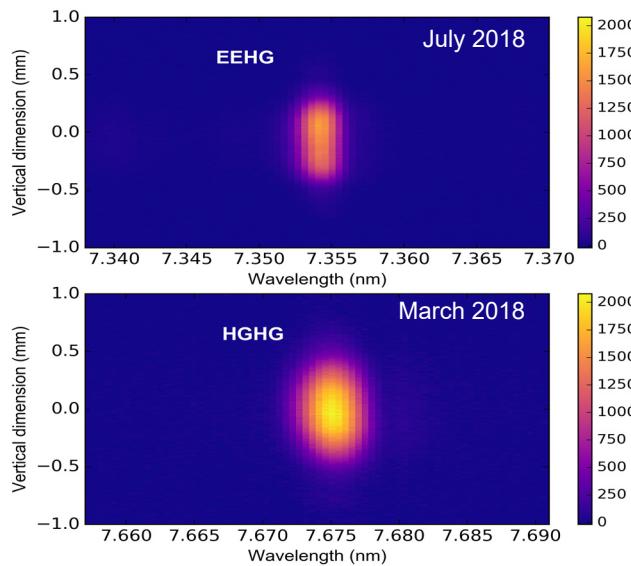
**Experimental exponential growth rate matches** results of **numerical simulations** with our parameters.



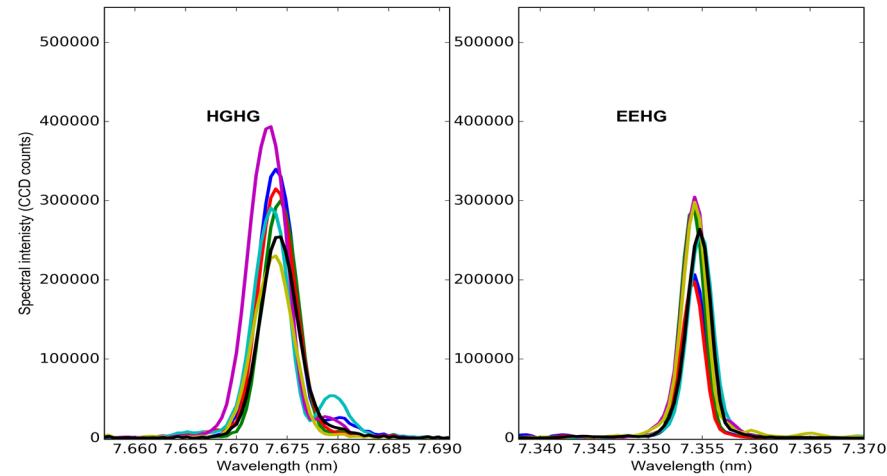
Red dots: rescaled experimental FEL intensity  
Blue line: numerical simulation

# EEHG pulse energies

**Measurements** with calibrated diagnostics (photodiode, CCDs) **suggests** that  $\sim 20\mu\text{J}$  per pulse are generated at  $\sim 7\text{nm}$ .



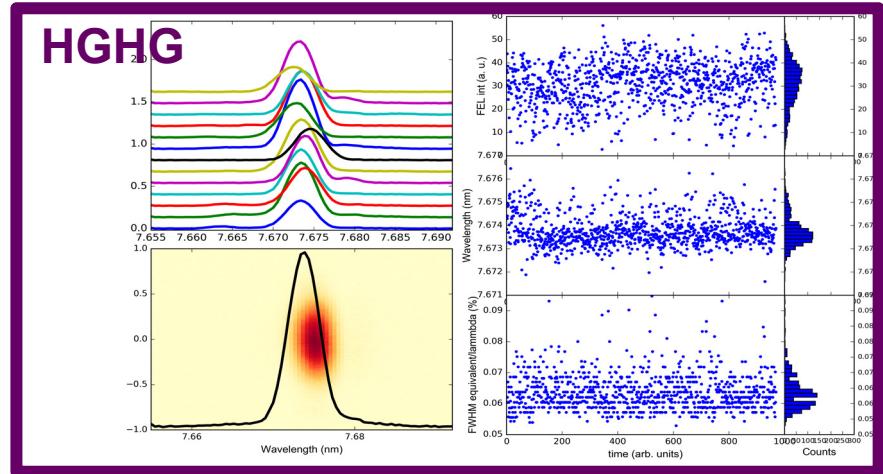
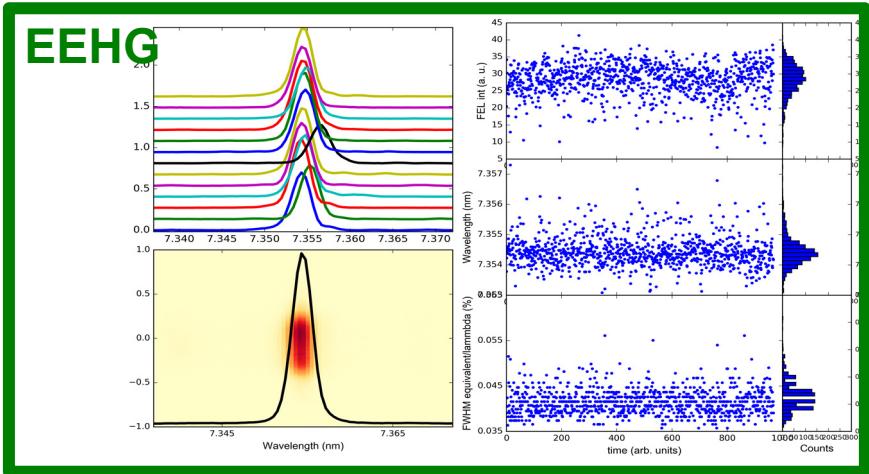
At 7.3 nm direct comparison of the CCD intensity from the **best EEHG results** and a **good FEL** from **HGHG** shows a **similar level of intensity**.



Relaxing requirements on **spectral purity**, **HGHG** can generally **increase the energy** per pulse. At **shorter wavelengths** **HGHG** has **higher energy** per pulse due to the higher flexibility in maximizing bunching (stronger seed).

# EEHG spectral control (7nm)

Best EEHG vs. a very good HGHG-FB (**not same beam not same seed laser**):



Both cases FEL optimized for **clear spectra**.

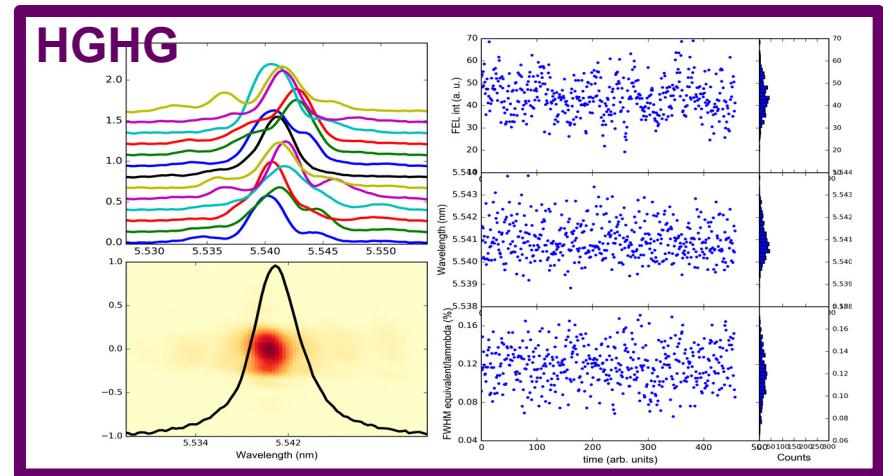
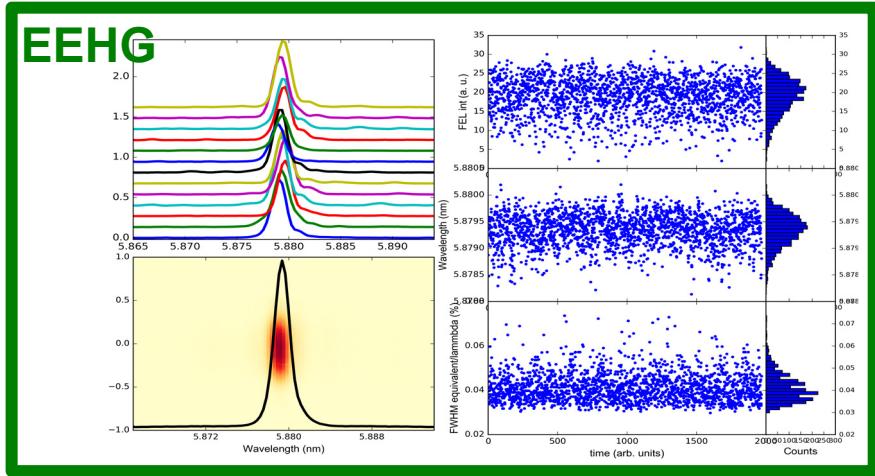
FWHM spectral width **might** be related to slightly **different seed** laser pulses.

FW<sub>76%</sub> (width containing 76% of the energy) is affected by **sidebands** that are more evident in HGHG.

Spectral **reproducibility** is higher in **EEHG**.

# EEHG spectral control (5nm)

Best EEHG vs. a very good GHG-FB (**not same beam not same seed laser**):



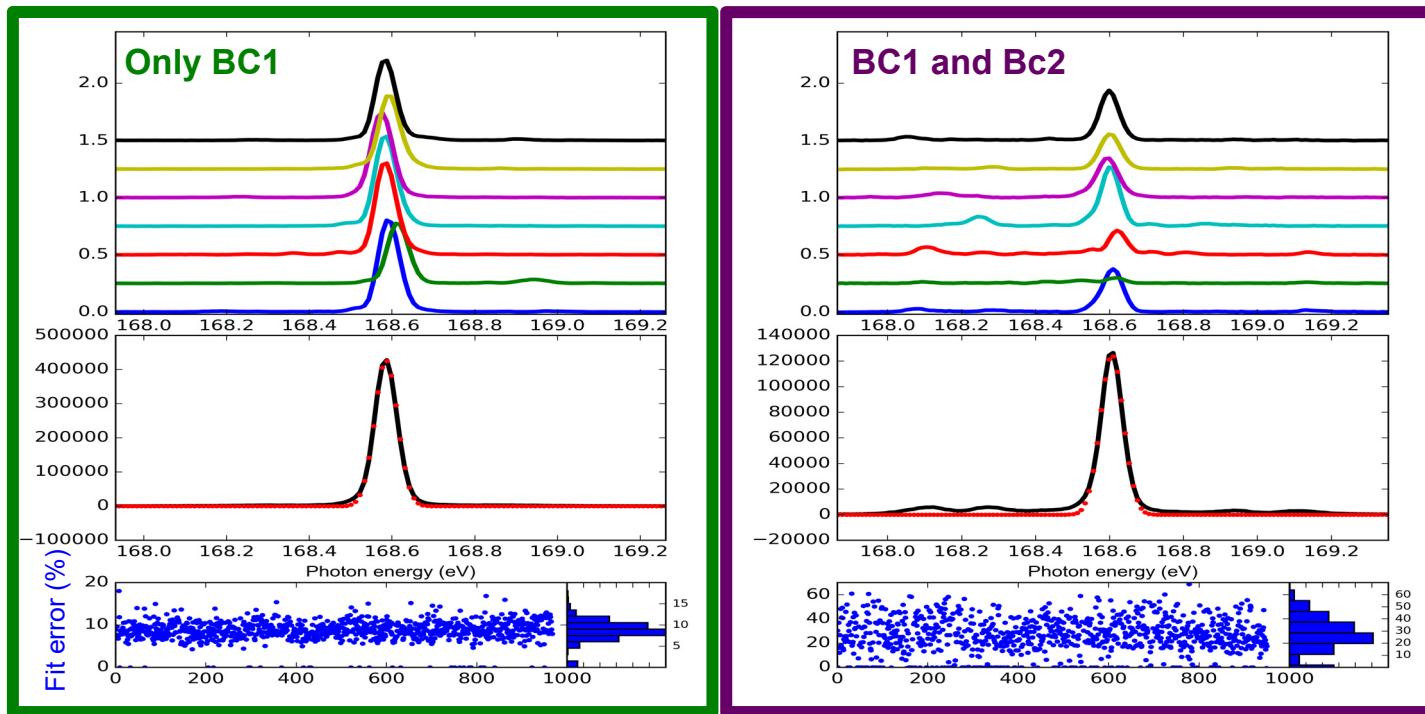
Difference is more pronounced at higher harmonics.

For **GHG cleaner spectral** can generally be obtained by **reducing** the overall FEL intensity.

At harmonics ~45 also **EEHG starts** to show **some structure** in the spectrum.

# EEHG sensitivity to $\mu\mathbf{B}$

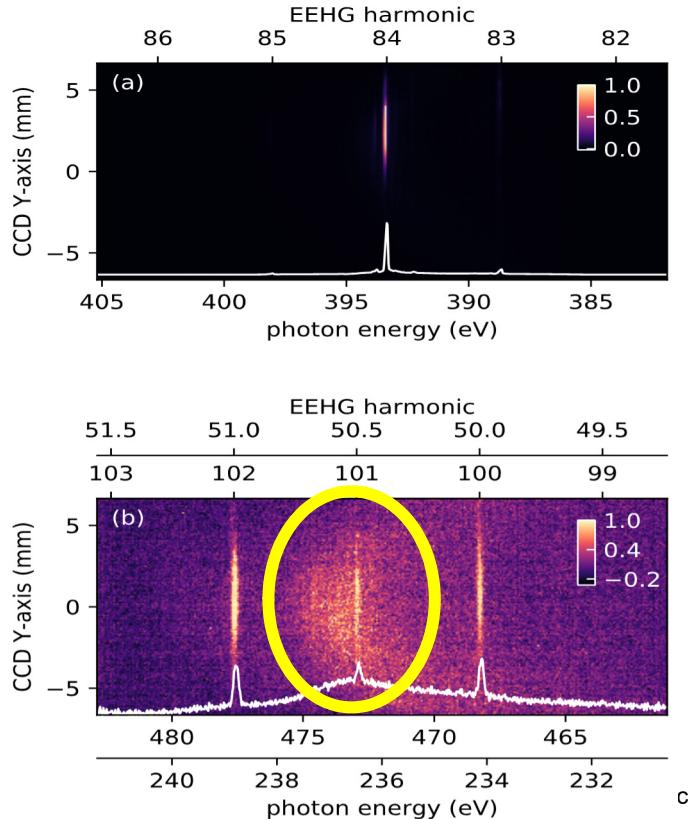
Using **2 bunch compressors** in the linac has an impact on the  $\mu\mathbf{B}$ .



Higher  $\mu\mathbf{B}$  has a clear **impact** on the **spectral purity** also for **EEHG** as seen from the Gaussian fit error.

# EEHG at very high harmonics

With the e-beam energy at **1.5 GeV** harmonic 84 (**3.1 nm**) and  $h=101$  (**2.6nm**) have been measured.



**Limits in the experimental setup forced us to operate at  $n=-3,-4$  and low undulator  $k$ .**

Nevertheless we have acquired **single** shot spectra with the EUV CCD (Andor) at **3 nm** (shortest wavelength for first order of diffraction)

Measurements at **2.6 nm** required second order of diffraction and few seconds CCD integration time.

## Conclusions

- ◆ Capabilities of EEHG have been fully demonstrated down to harmonic 45 (5.8 nm).
  - Good agreement between simulations and experiment, work to be continued.
- ◆ From the limited experience EEHG operations at short wavelengths (limited gain) appears more critical than HGHG-FB.
  - EEHG strongly relies on exponential amplification.
- ◆ EEHG confirms promises of producing bunching at very high harmonics
  - $H = 101$  clearly visible in a non optimized setup.
- ◆ New FEL capabilities can become available with EEHG in the 20 – 4 nm spectral range.

An optimized EEHG configuration **appears** a viable **option** for extending the tuning range of **FEL-2 to shortest wavelengths** if a **suitable FEL gain** can be **provided**.

Work started at FERMI for a possible upgrade based on EEHG.

# Acknowledgements



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# Thank you!