

Muon Beam Dynamics and Spin Dynamics in the g-2 Storage Ring

D. L. Rubin, Cornell University

A.Chapelain, Cornell University, F. Gray, Regis University

S.Charity, J. Price, University of Liverpool

Joe Mott, Boston University

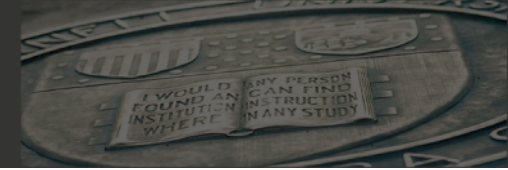
J.D.Crnkovic, W. Morse, V. Tishchenko, Brookhaven National Laboratory

W. Wu, University of Mississippi

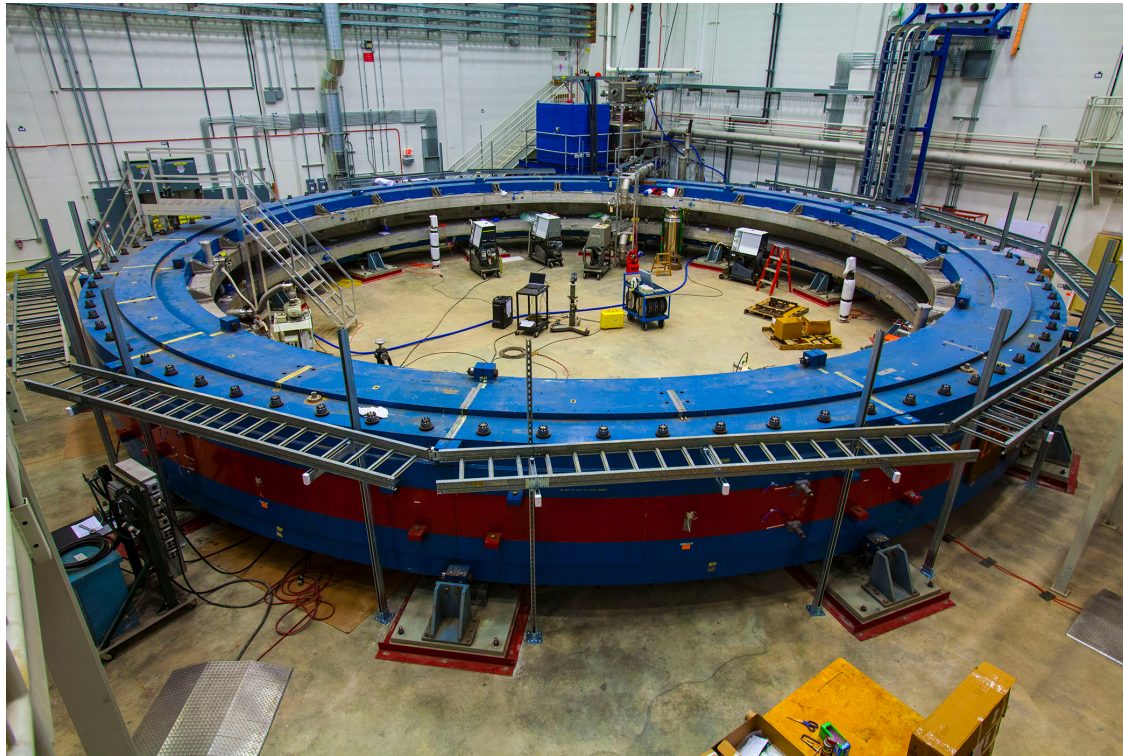




- Goal
 - measure the anomalous magnetic moment of the muon (g-2) with 0.14 ppm uncertainty (*BNL 0.54 ppm*)
- Method
 - Circulate polarized muons at the ‘magic’ momentum, in uniform magnetic field and measure precession frequency
- Polarized muons
 - select of highest energy muons in $\pi^+ \rightarrow \mu^+ \nu_\mu$
- Measure muon polarization
 - Correlation with energy of decay positrons



- Commissioning run summer 2017
- Production run is just beginning
- Statistics in excess of BNL by summer 2018



- *Upgrades of several systems planned for summer 2018*



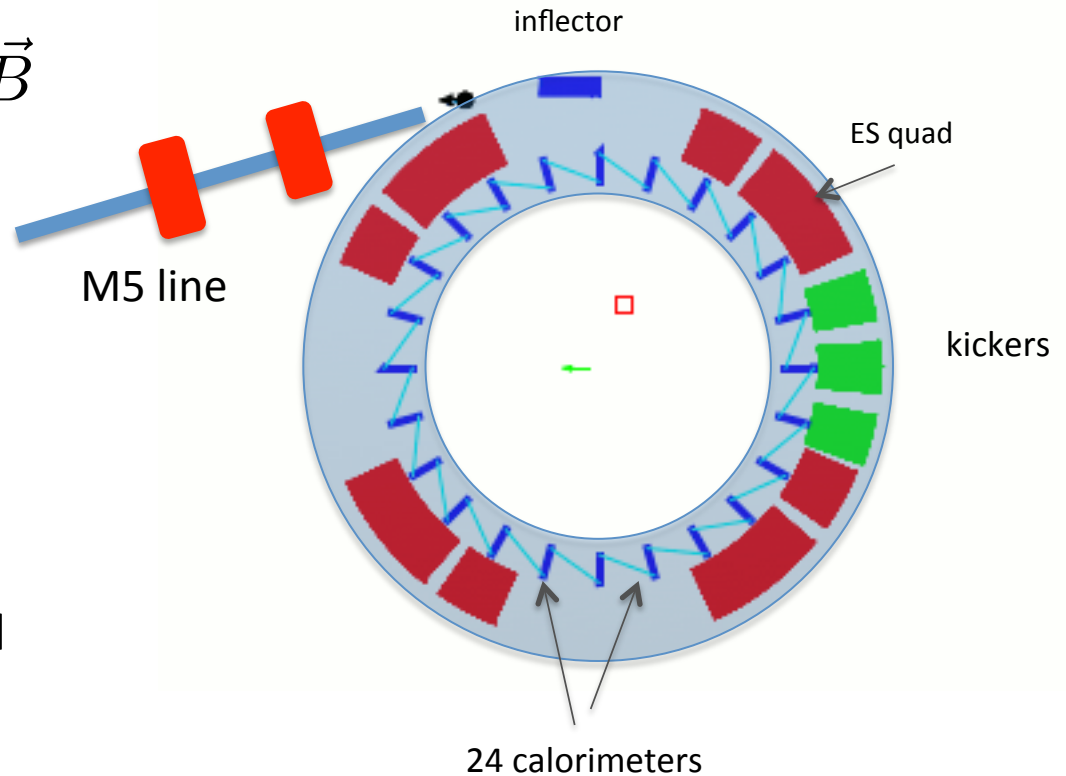
Muon storage ring

$$a_\mu = \frac{g - 2}{2}$$

$$\vec{\omega}_a = \vec{\omega}_S - \vec{\omega}_C = -\frac{q}{m} a_\mu \vec{B}$$

(ideal field and trajectory)

Requirement of uniform B-field
=> electrostatic focusing



Circumference = 44.69 m



Muon storage ring

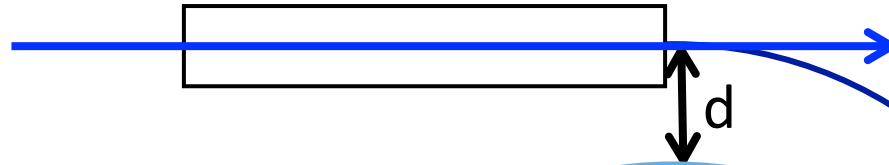


- Circumference = 44.69 m
- Revolution period = 149.1 ns
- $B_y = 1.45$ T
- $P_\mu = 3.09$ GeV/c
- $\gamma = 29.3 = \left(\frac{1}{a_\mu} - 1\right)^{1/2}$
- Aperture = 9 cm
- $\tau_\mu = 64.4$ μ s (at 3.09 GeV/c) = 432 turns



Electrostatic quadrupoles

inflexor



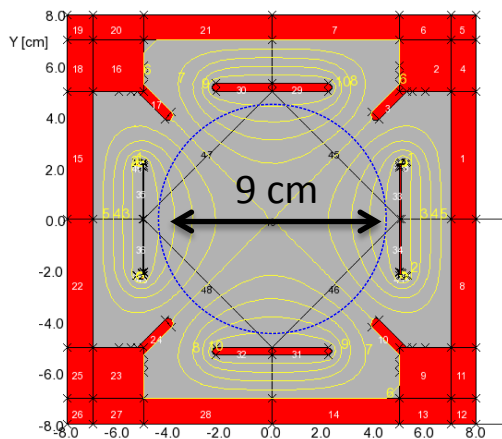
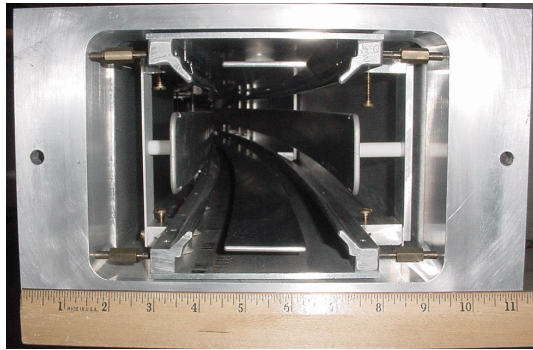
quadrupoles

711.2 cm

kickers

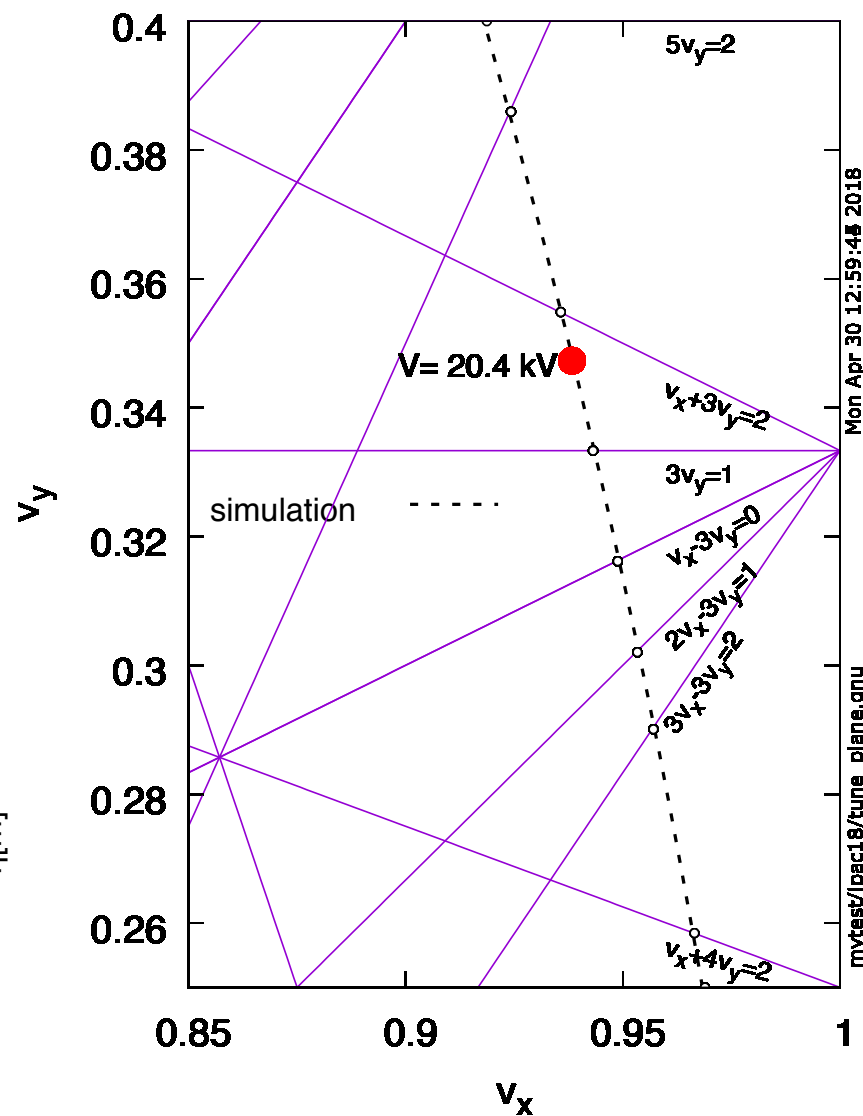
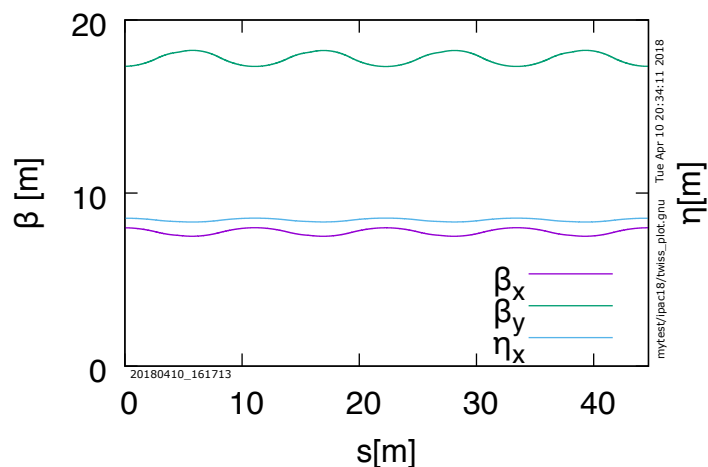
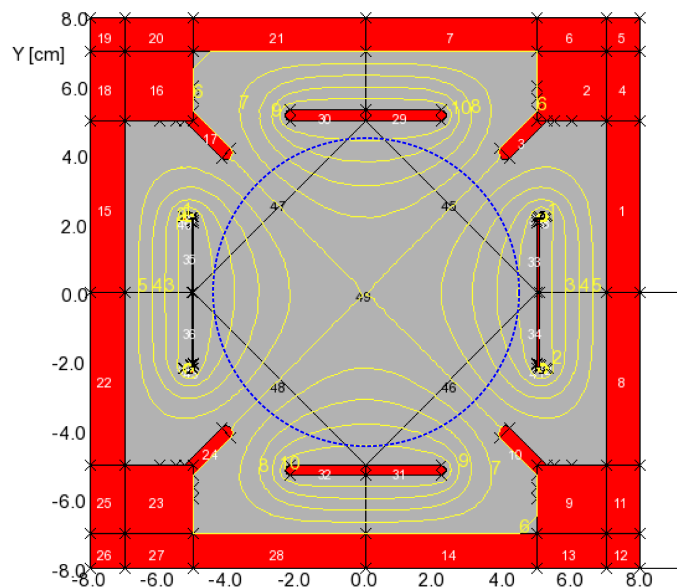
10.8 mrad

Quads span 156° of arc



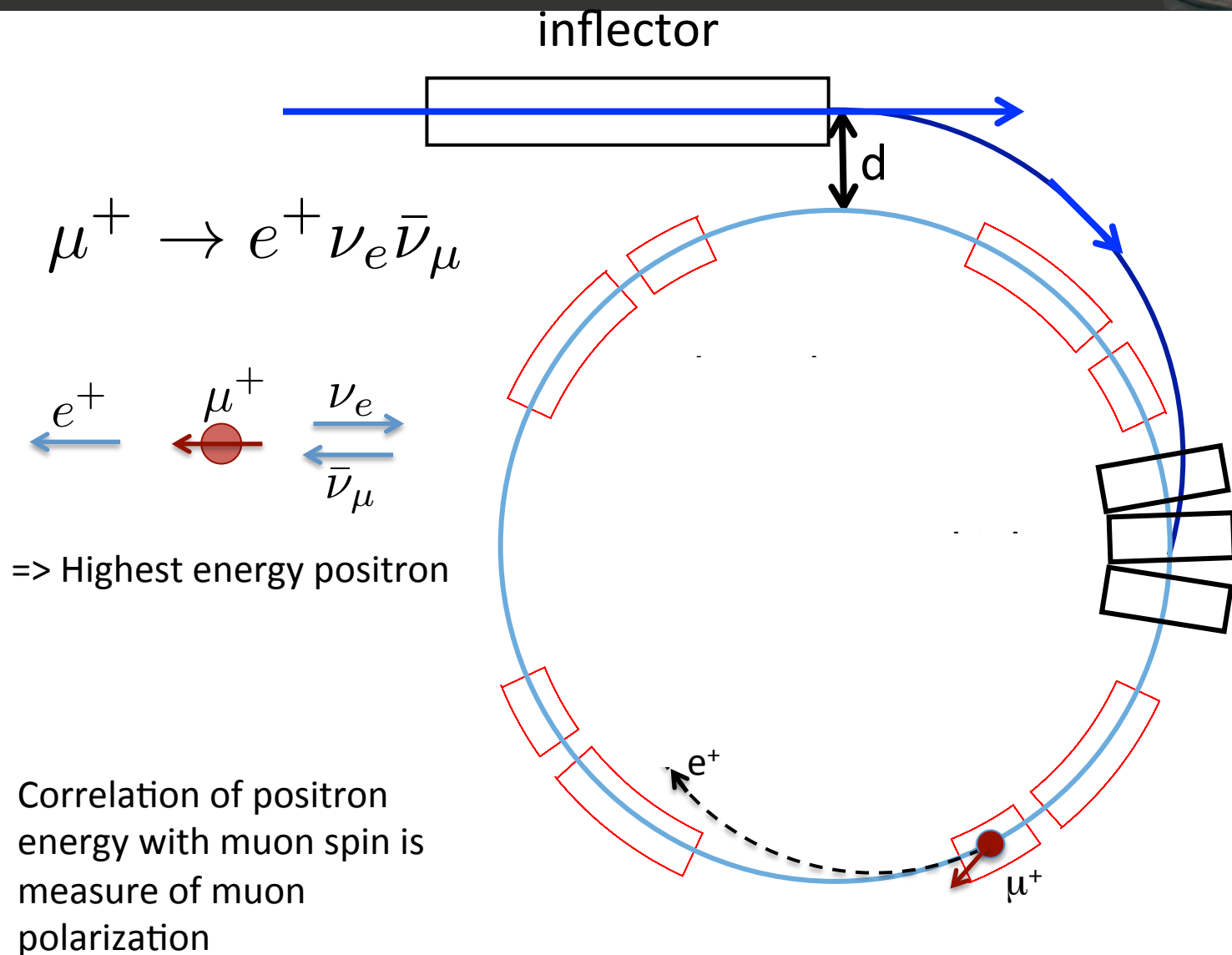


Weak focusing lattice



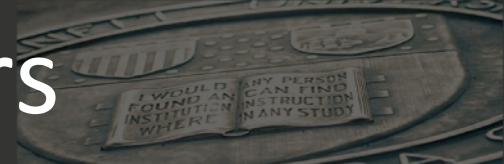


Muon decay

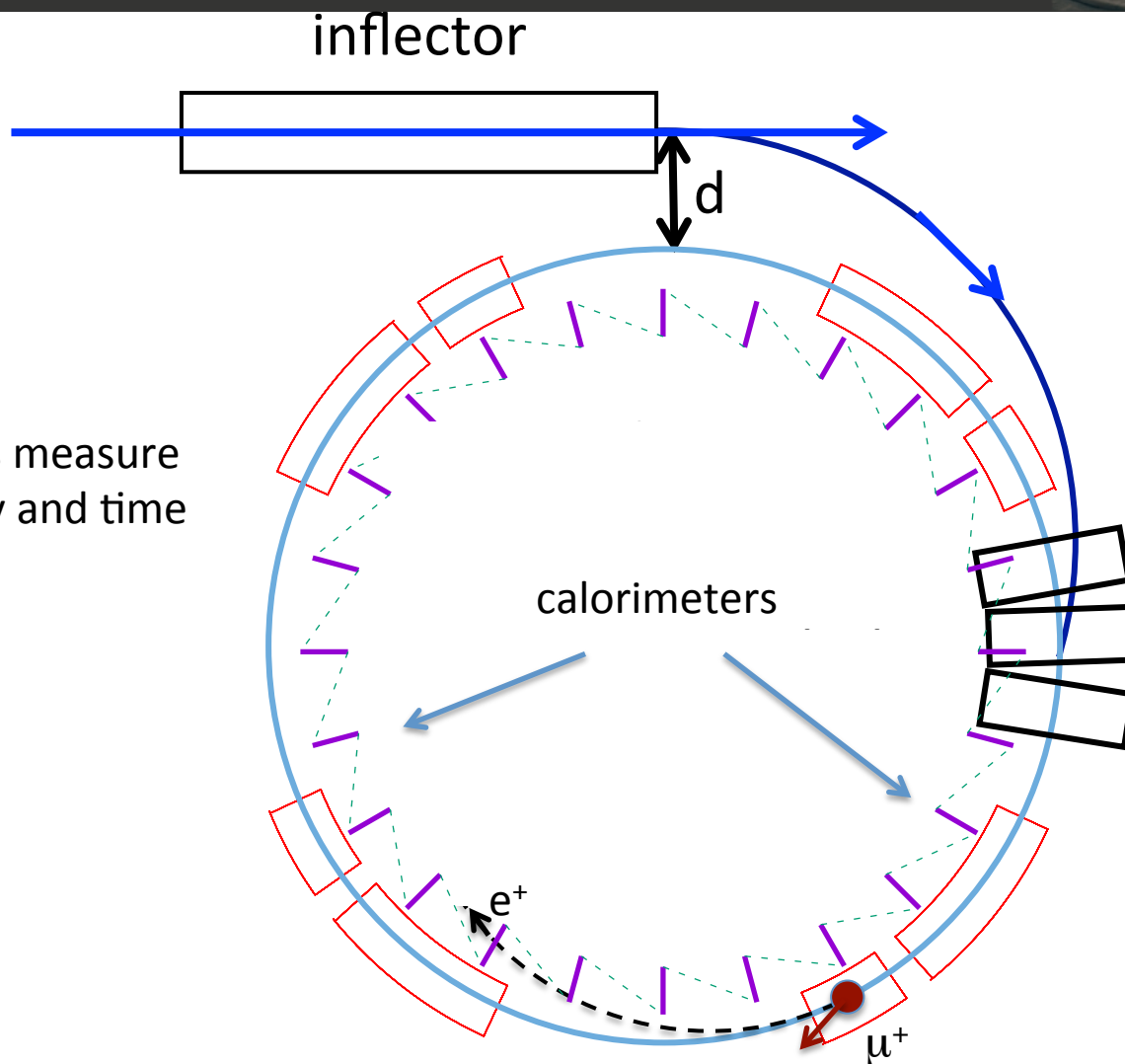




Calorimeters

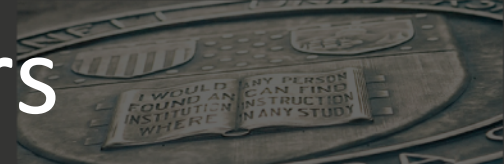


24 calorimeters measure
positron energy and time

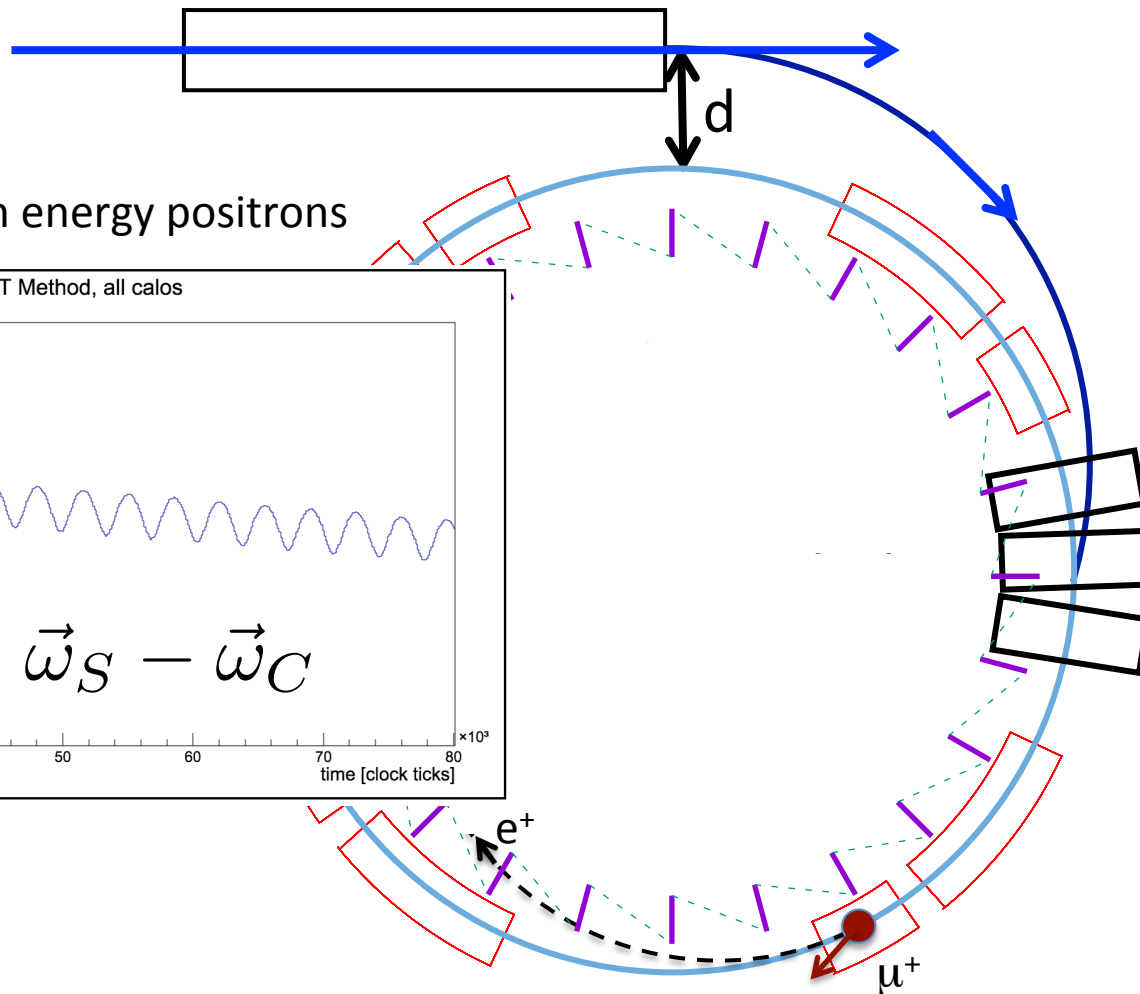




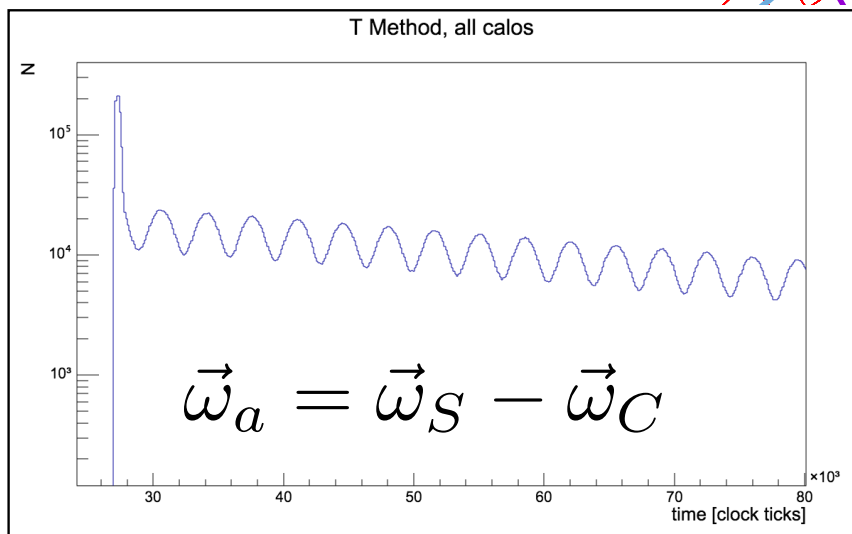
Calorimeters



inflexor

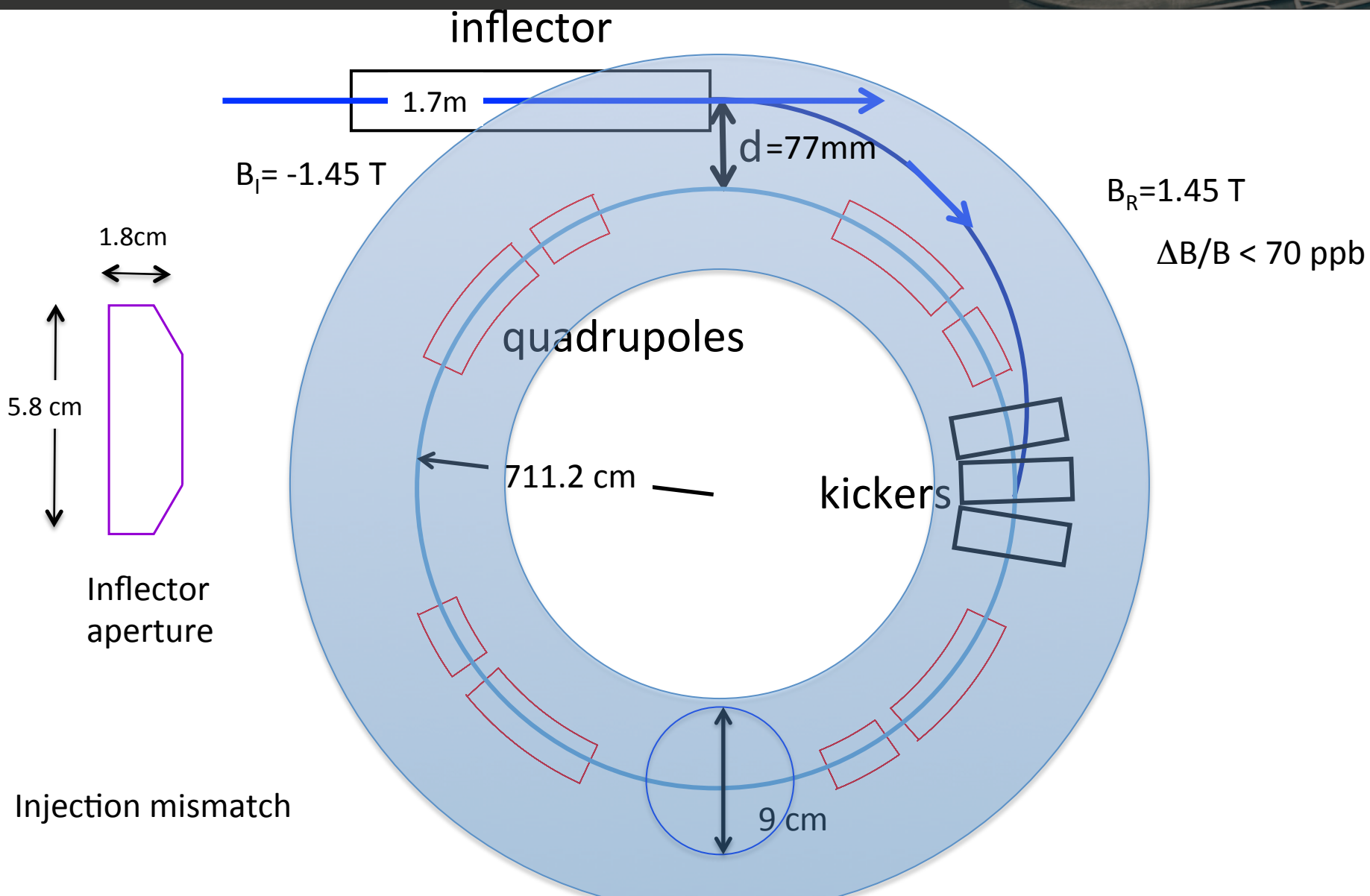
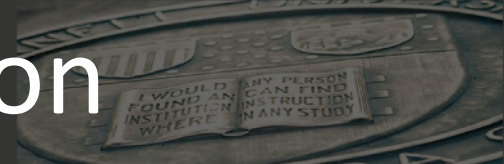


Histogram of high energy positrons



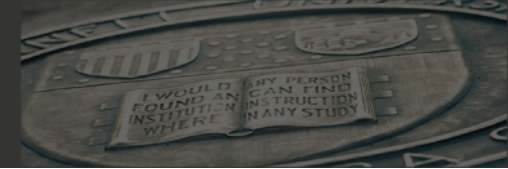


Injection





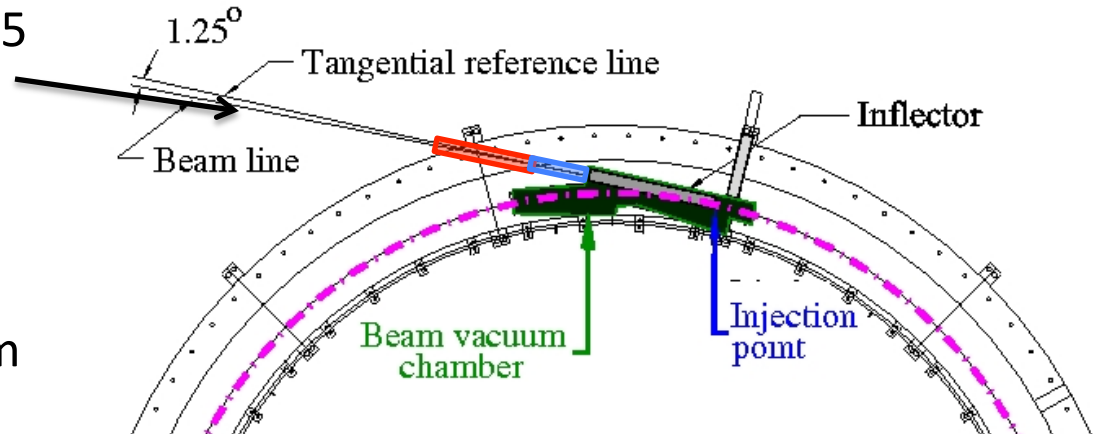
Injection



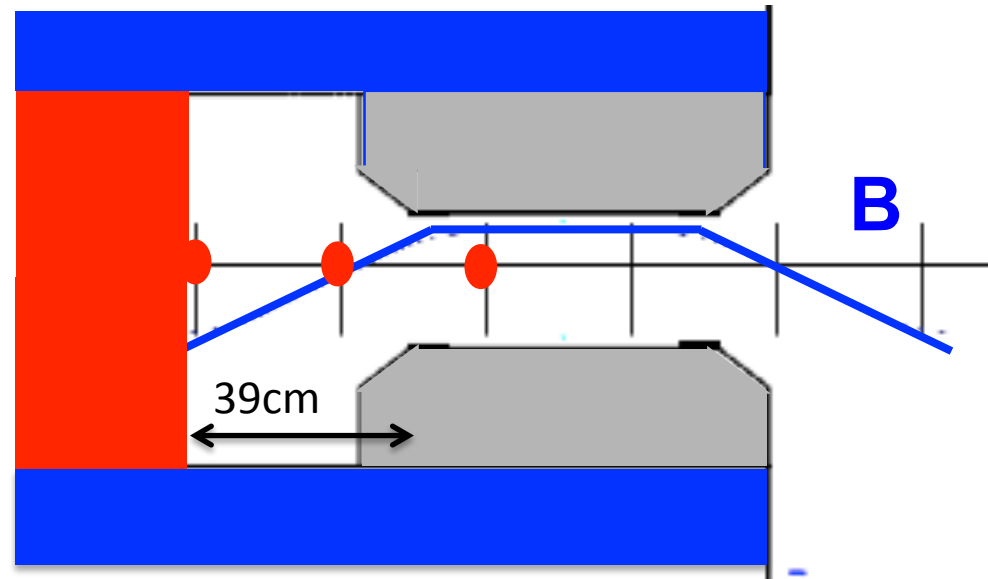
Muons come to the end of the M5 line and then propagate through:

- Hole in magnet yoke
- Dipole fringe field
- Inflector

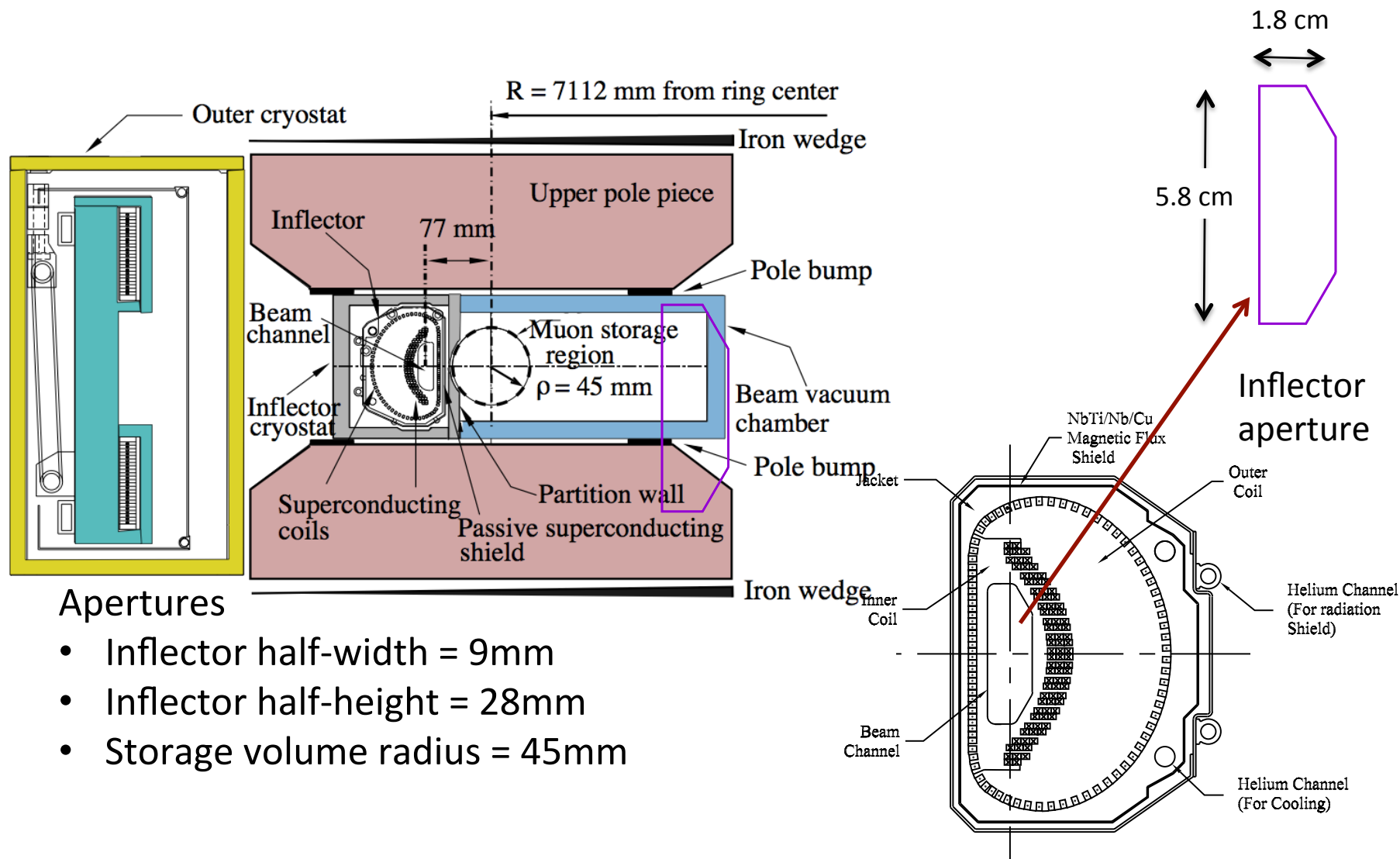
And exit the inflector 77 mm from the center of the dipole aperture



The magnetic field is near zero at the inner surface of the yoke, and rises to 1.45T between the magnet poles, over a distance of ~39cm



Inflector



Apertures

- Inflector half-width = 9mm
- Inflector half-height = 28mm
- Storage volume radius = 45mm

Fig. 7. Inflector cross-section.



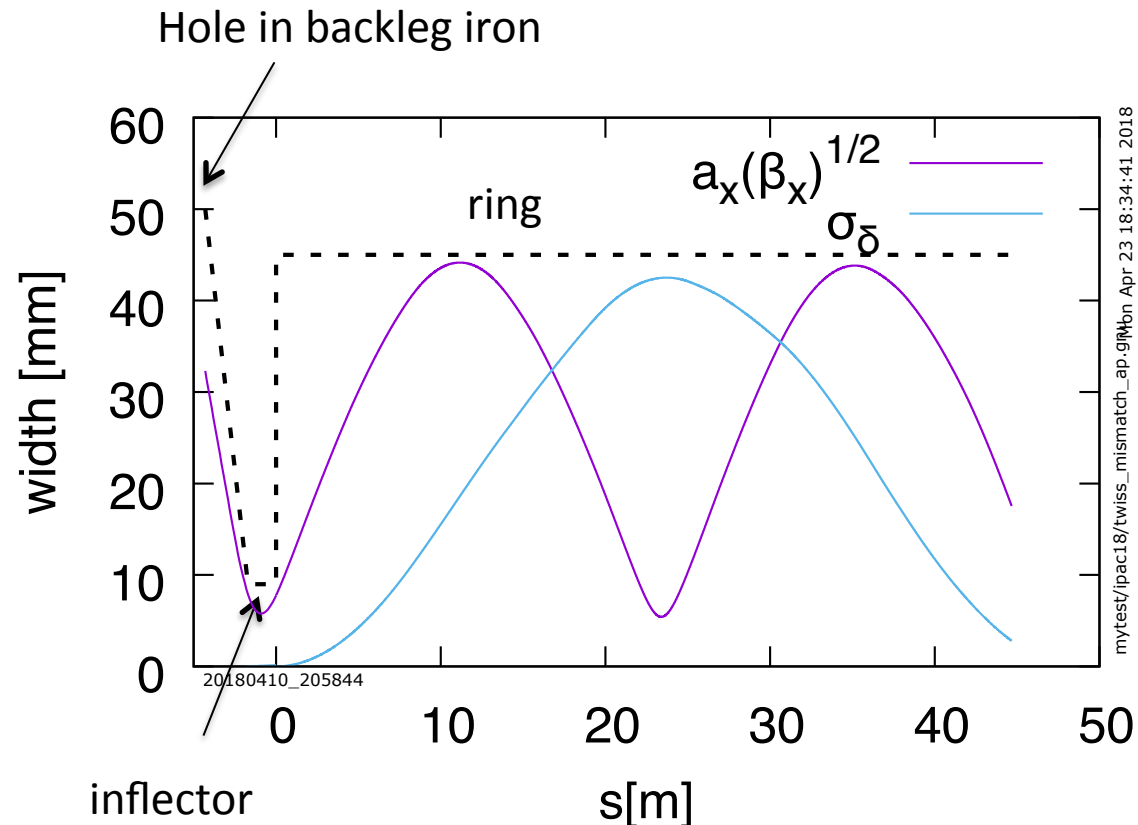
Injection phase space

Injected beam

$$\varepsilon = 40 \text{ mm-mrad (95\%)}$$

$$\Delta p/p = 2\%$$

- Beam is focused to a waist to clear inflector aperture
- (Muons are scattered in coil overlapping inflector aperture)
- Dispersion is zero in inflector

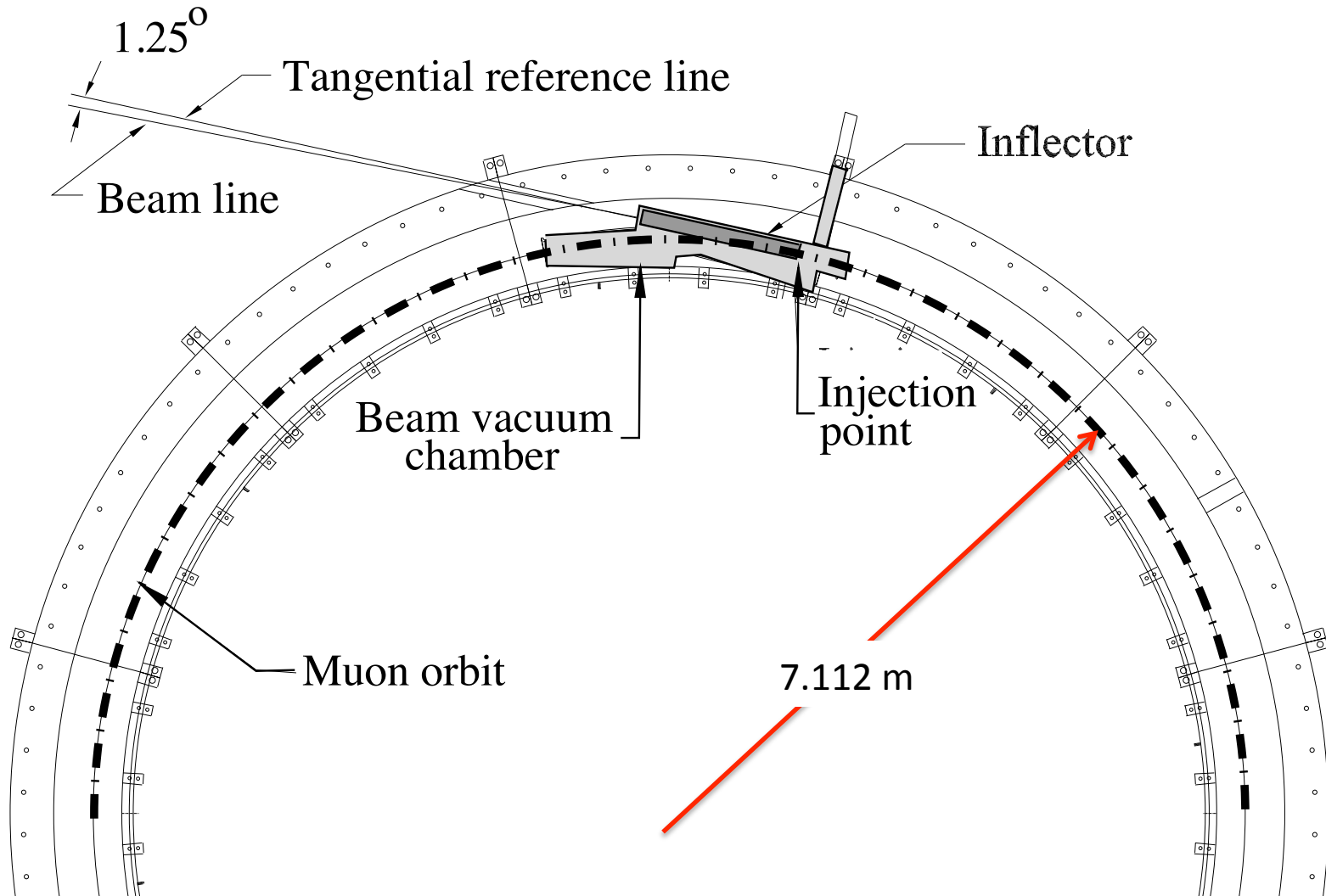


=> Dispersion and phase space mismatch

Ring
 $\beta_x = 7.8 \text{ m}$
 $\eta = 8.2 \text{ m}$



Injection channel



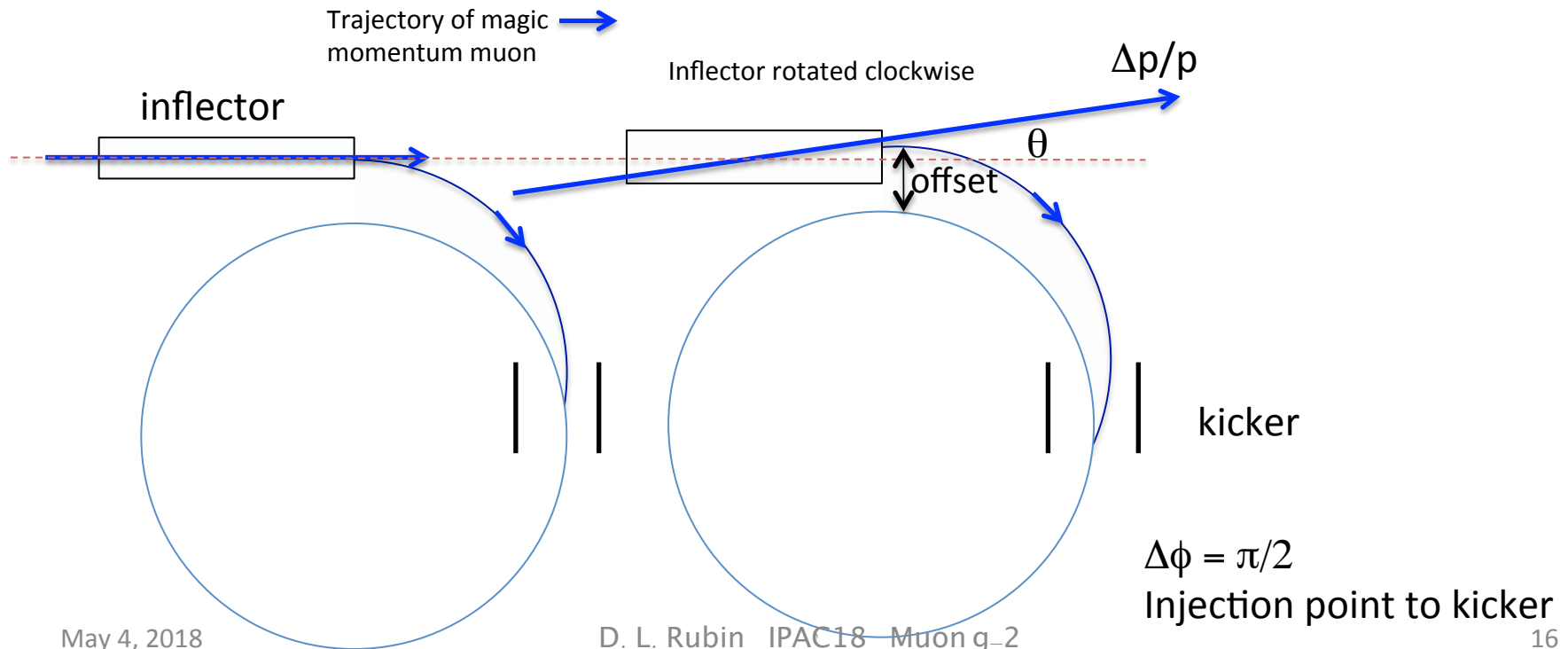


Injection Dependencies

Injected particles oscillate about the closed orbit with amplitude that depends on;
momentum, injection angle, offset, and kick angle

Amplitude and envelope of betatron oscillation of stored beam depends on

- injection angle θ and
- Kicker angle
- Momentum

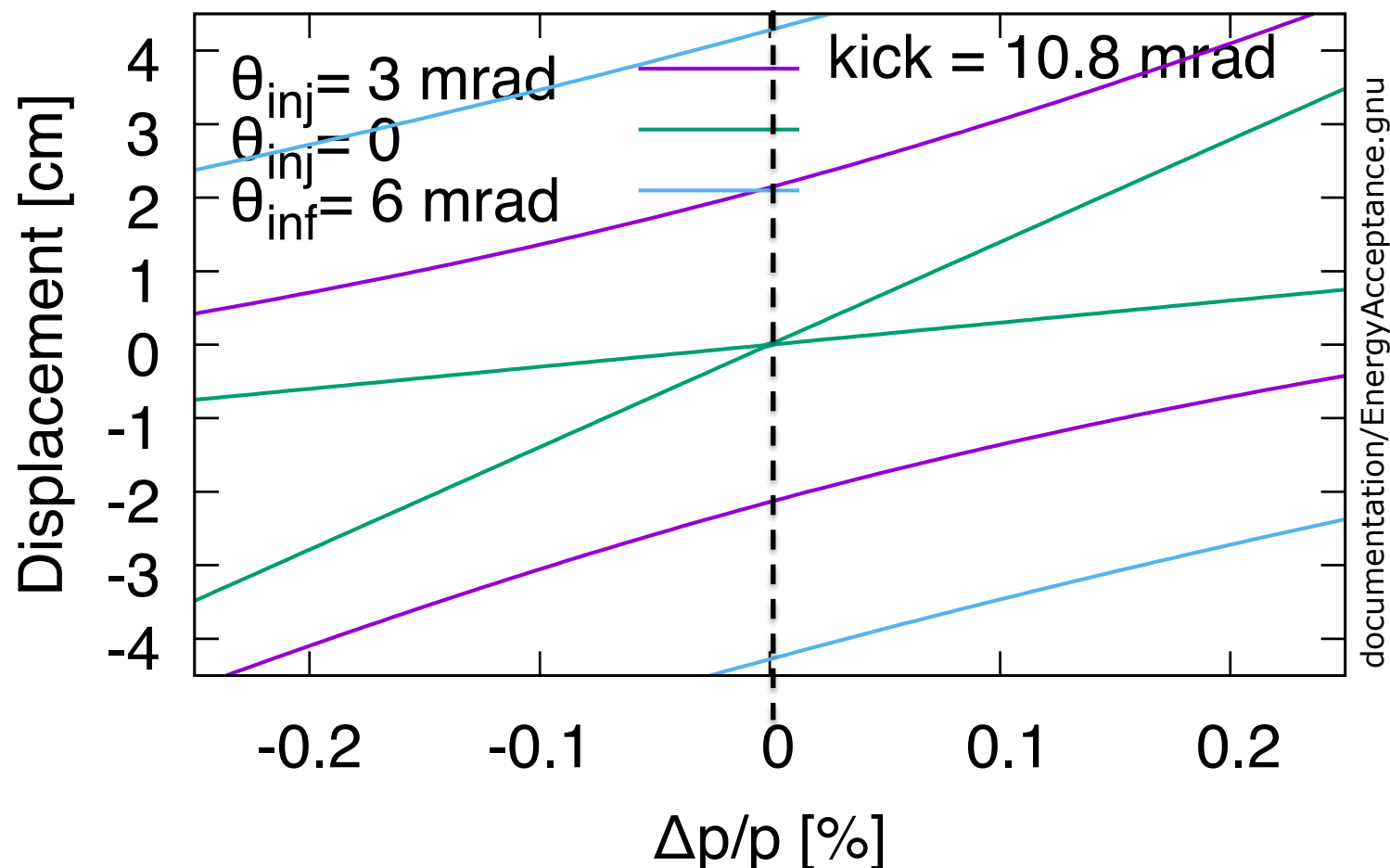




Muon capture



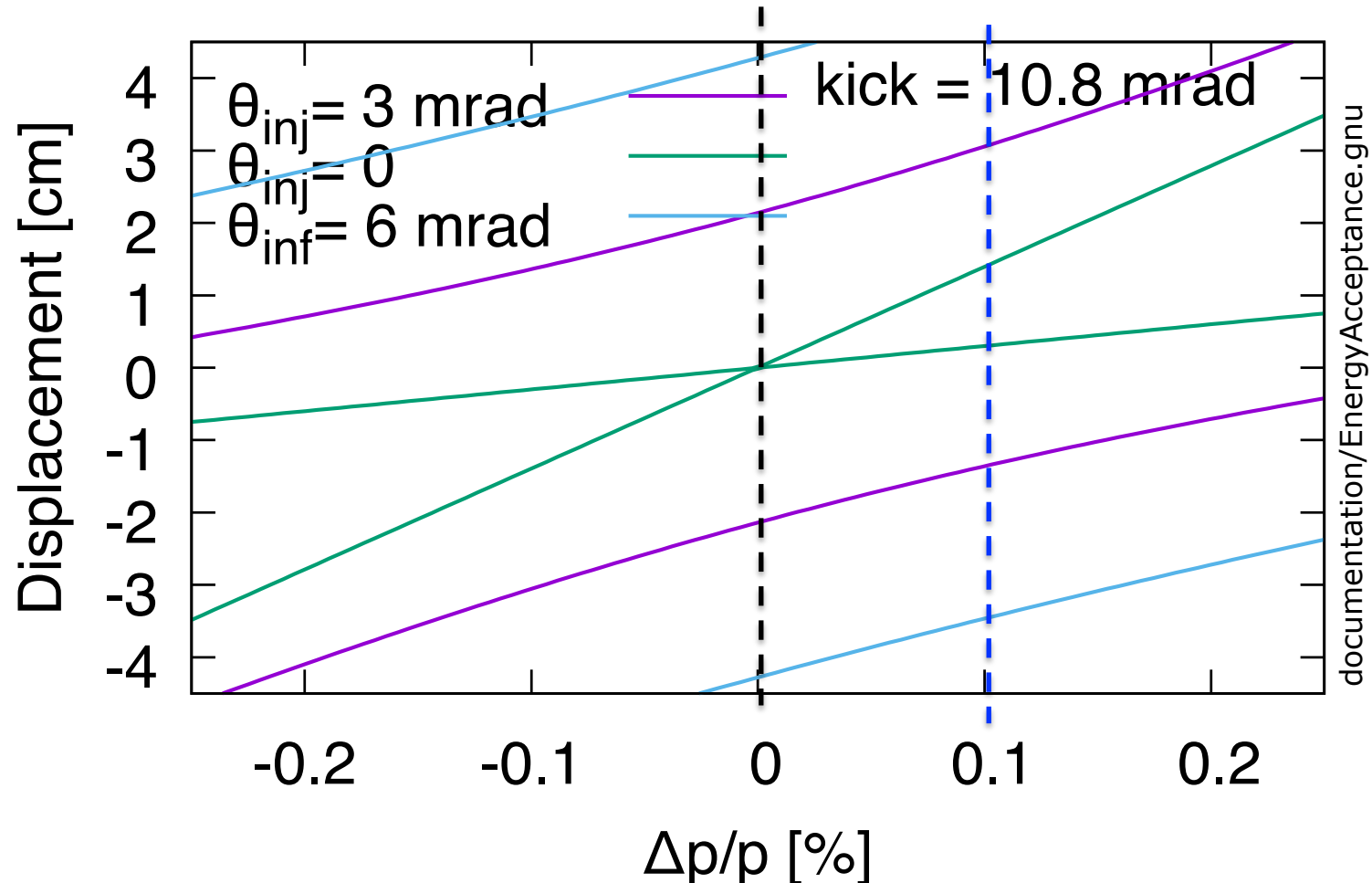
Envelope of betatron motion of stored muon
(10.8 mrad kick angle puts on momentum muon on axis)





Muon capture

Envelope of betatron motion of stored muon
(10.8 mrad kick angle puts on momentum muon on axis)

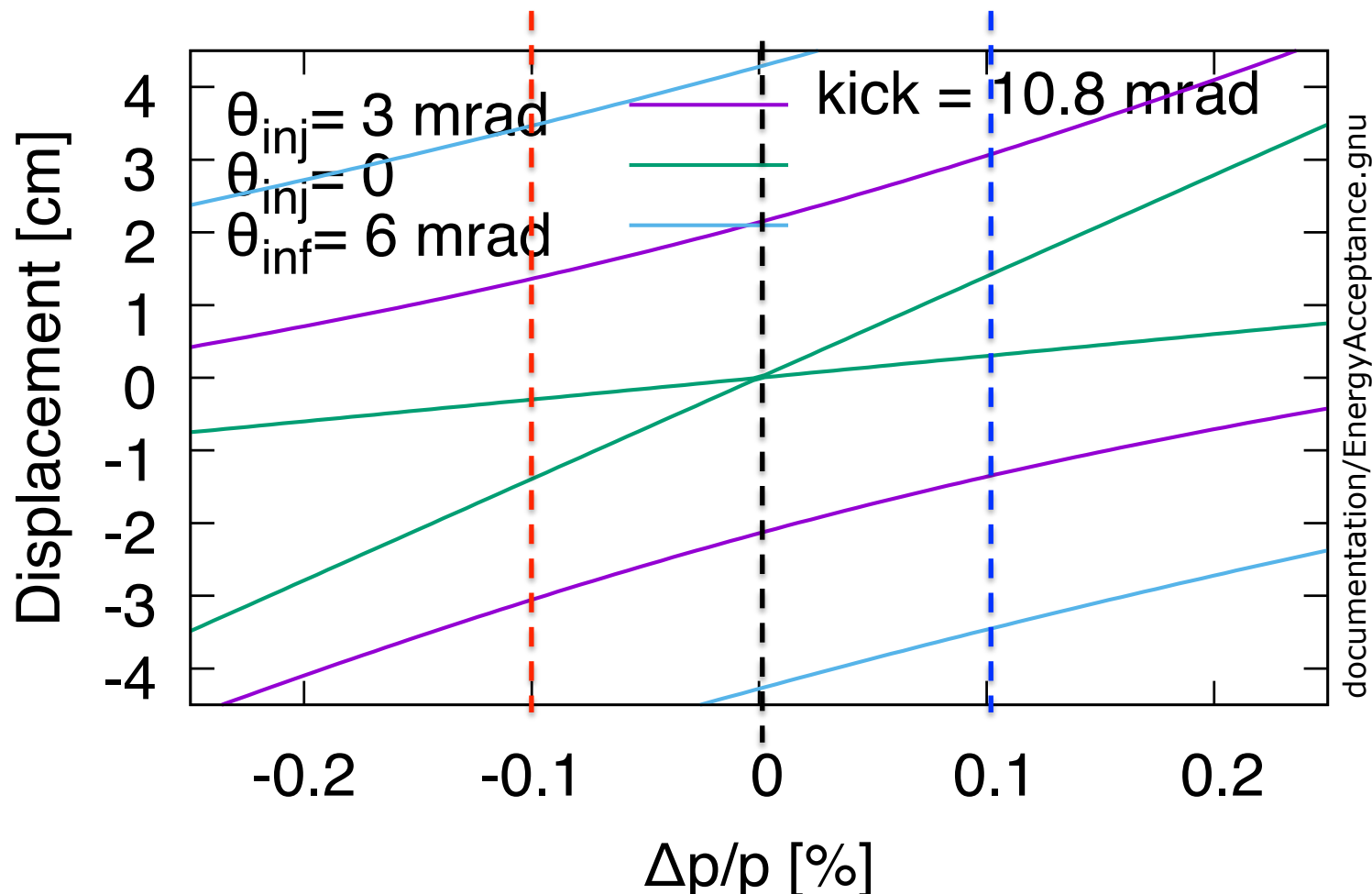




Muon capture

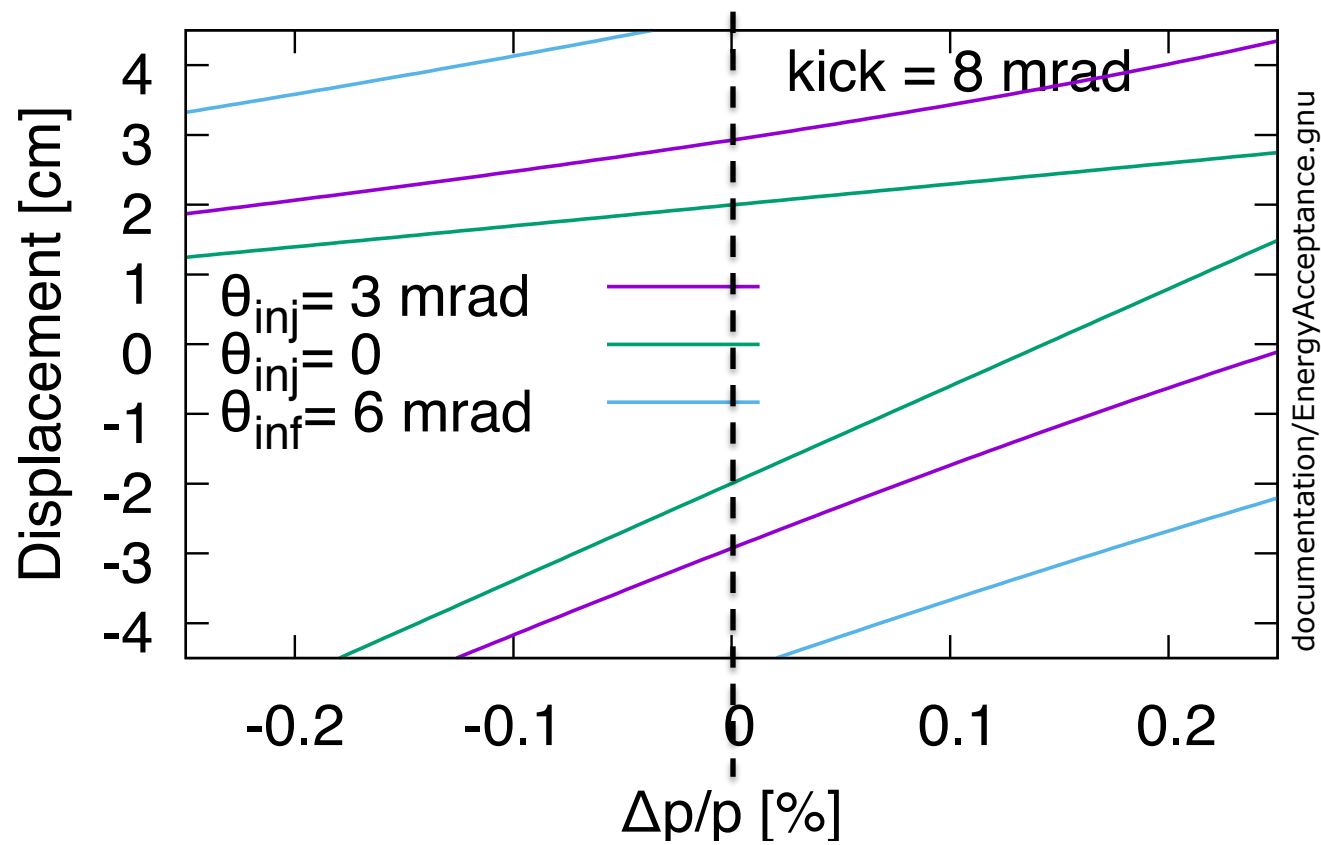


Envelope of betatron motion of stored muon





Envelope of betatron motion of stored muon (8 mrad kick = 75% nominal – under kicking)

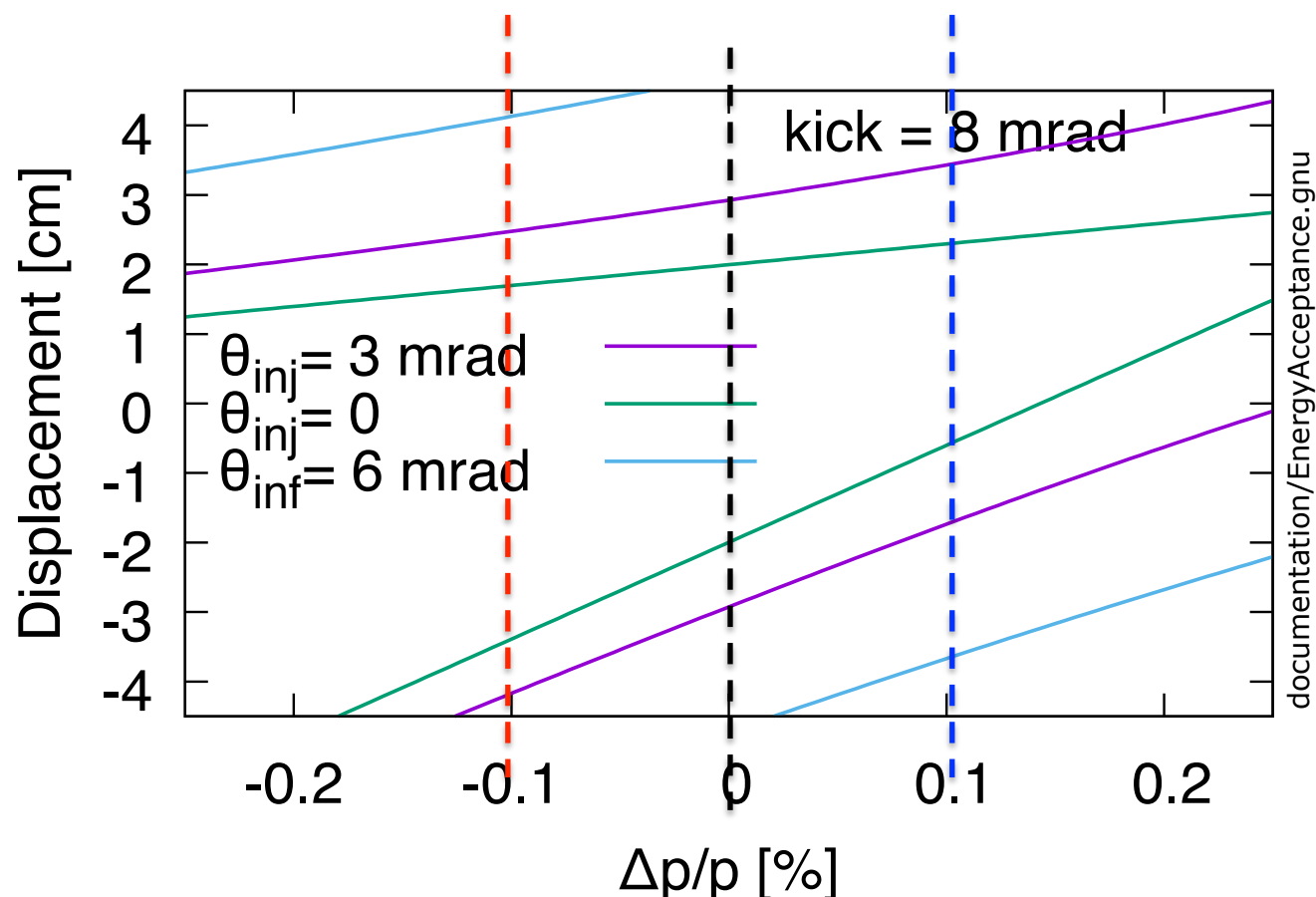




Muon capture

In the event of *under kicking*

- Acceptance is skewed to higher momenta





Systematics

$$\vec{\omega}_a = -\frac{q}{m} \left[a_\mu \vec{B} - a_\mu \left(\frac{\gamma}{\gamma + 1} \right) (\vec{\beta} \cdot \vec{B}) - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{\beta} \times \vec{E}}{c} \right]$$

In general

- Finite momentum distribution
- Transverse motion
- Non zero E-field



Pitch correction

$$\vec{\omega}_a = -\frac{q}{m} \left[a_\mu \vec{B} - a_\mu \left(\frac{\gamma}{\gamma + 1} \right) (\vec{\beta} \cdot \vec{B}) - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{\beta} \times \vec{E}}{c} \right]$$

Pitch correction:

- Contribution depends on vertical phase space distribution

$$\beta_y = p_y / (\gamma m c)$$

- Correction requires measurement of

vertical phase space



E-field correction

$$\vec{\omega}_a = -\frac{q}{m} \left[a_\mu \vec{B} - a_\mu \left(\frac{\gamma}{\gamma + 1} \right) (\vec{\beta} \cdot \vec{B}) - \left(a_\mu - \frac{1}{\gamma^2 - 1} \right) \frac{\vec{\beta} \times \vec{E}}{c} \right]$$

E-field correction:

- Choose momentum.

$$\text{At } \vec{p}_{\text{magic}} = 29.3 \text{ GeV}/c \rightarrow a_\mu = \frac{1}{\gamma^2 - 1}$$

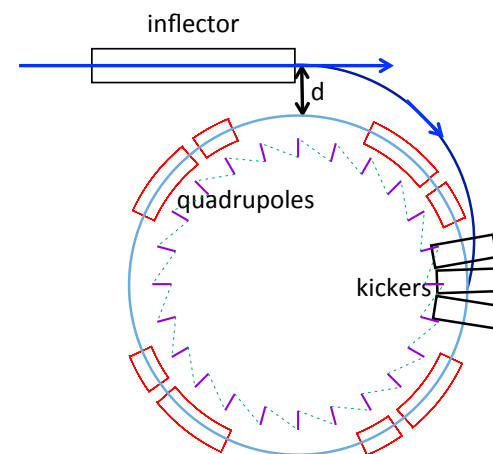
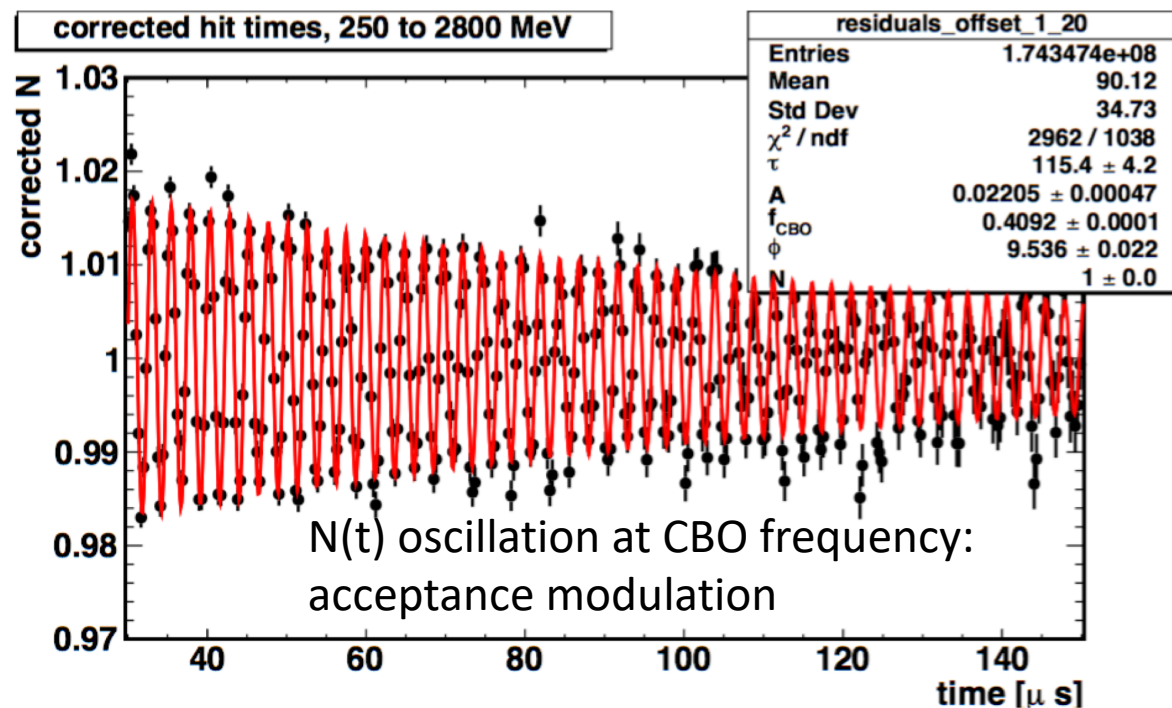
Correcting for contribution from non-magic momentum muons requires

- Measurement of *momentum distribution*
- Measurement of *radial distribution* $\vec{E} = k(x\hat{i} - y\hat{j})$

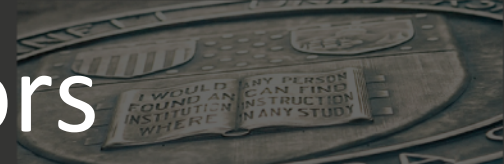
Acceptance correction

The acceptance of the calorimeters depends on the radial offset of the decay muon

- Radial offset varies with modulation of beam width
- And with oscillation of the centroid



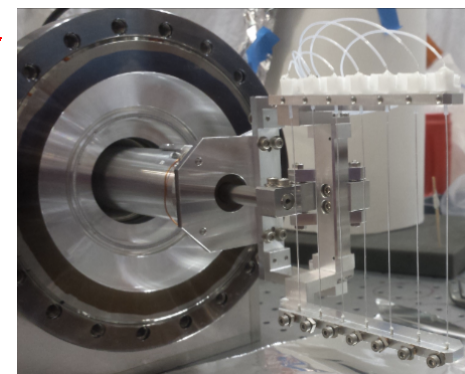
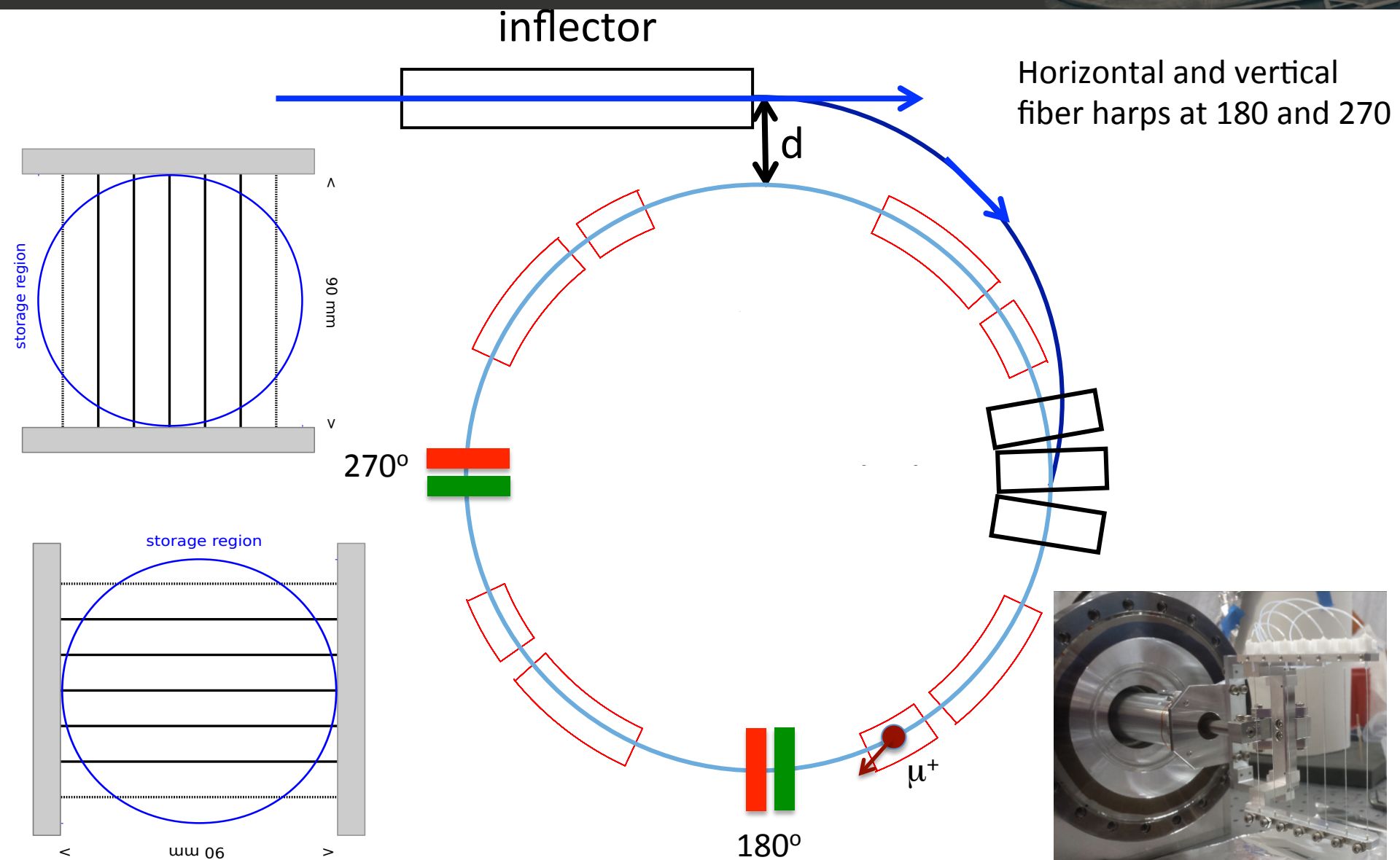
Correction requires measurement of
evolution of radial distribution



Fiber Harps
Traceback trackers
Calorimeters

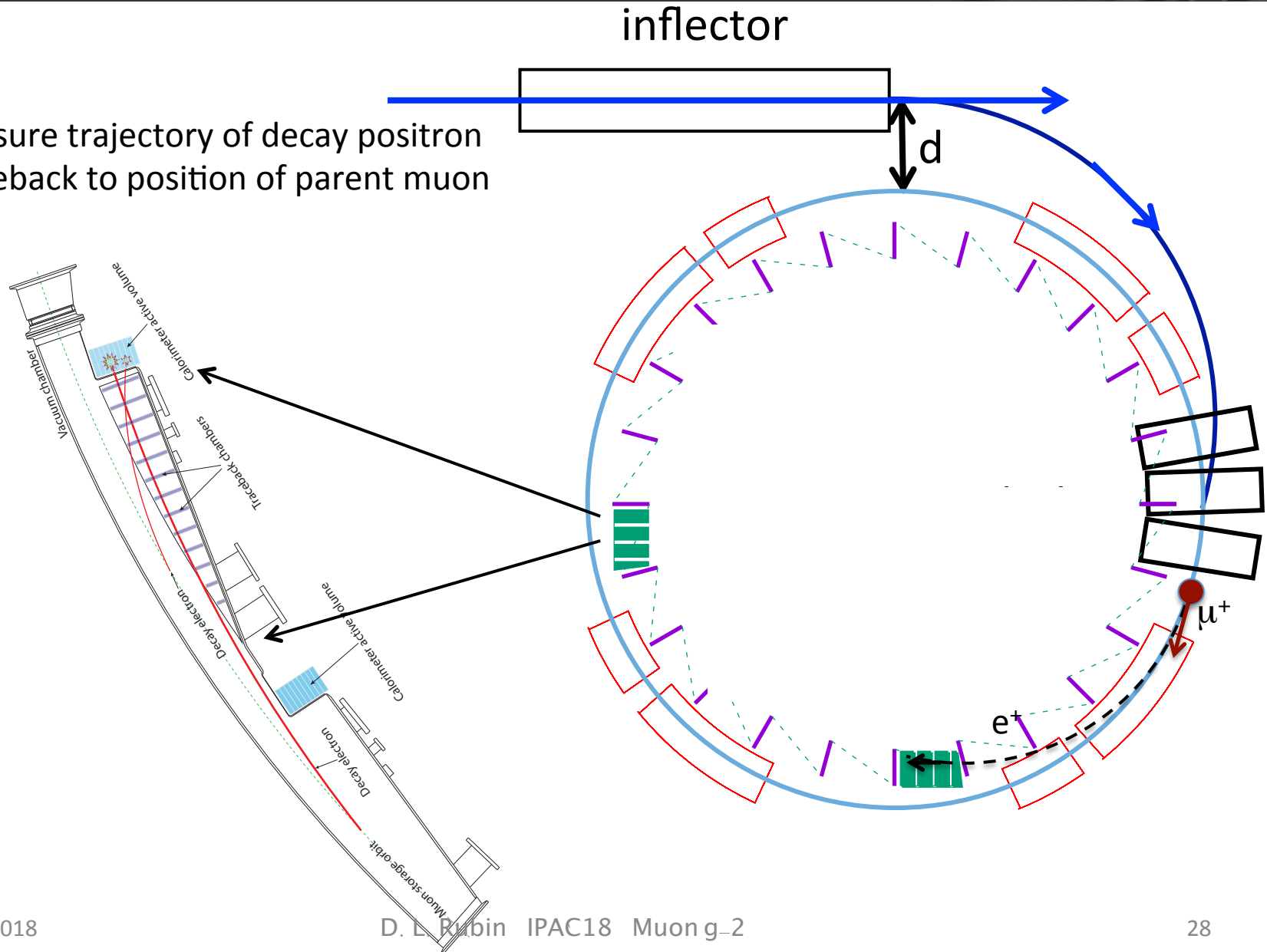


Fiber Harps



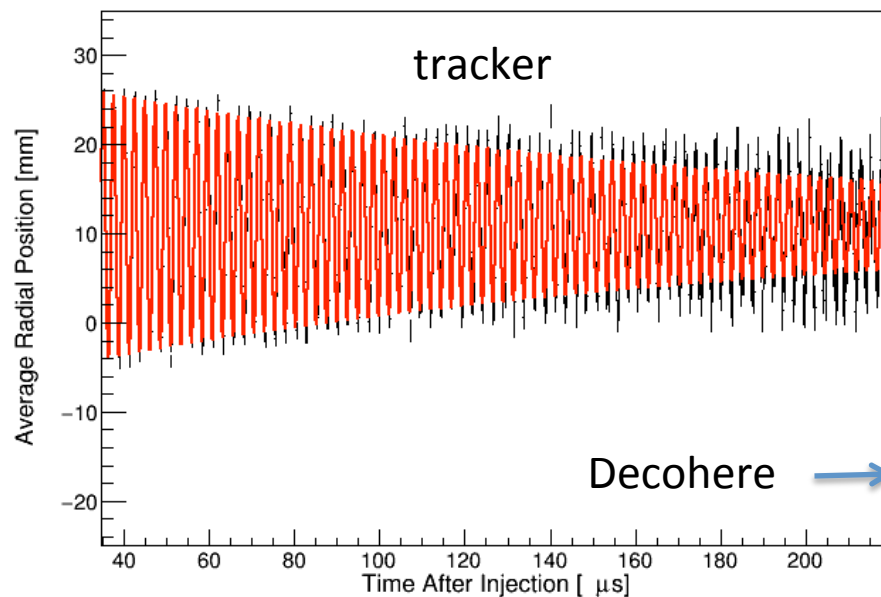
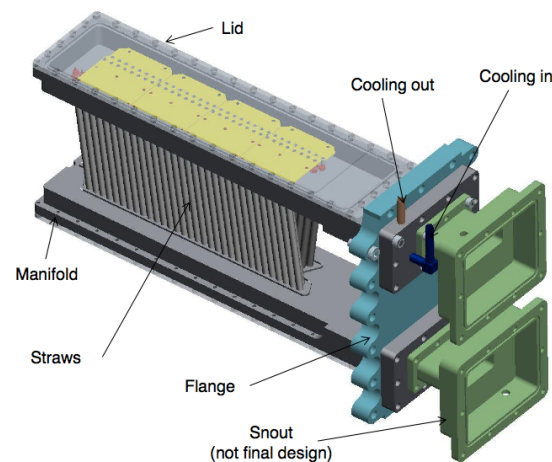
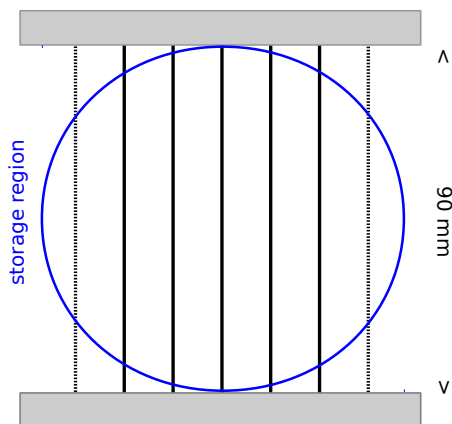
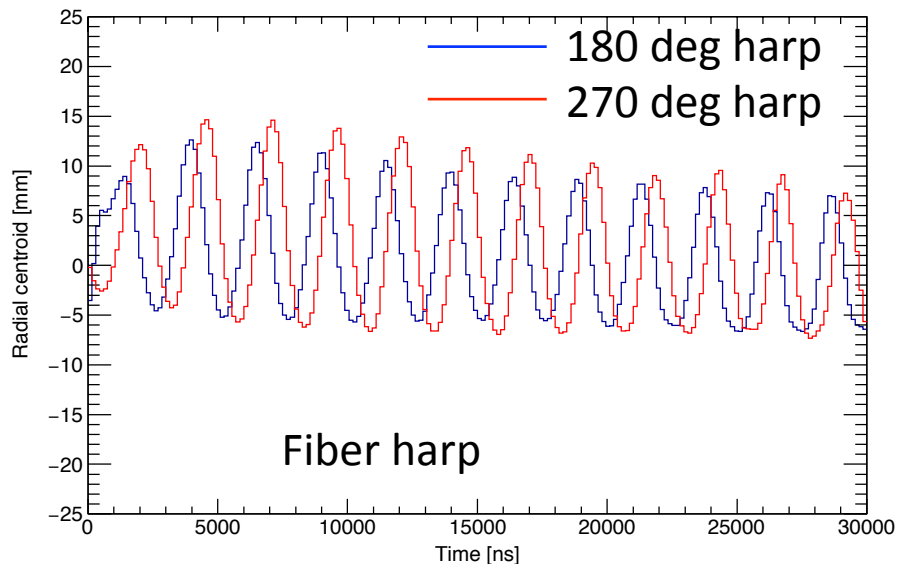
Traceback trackers

- Measure trajectory of decay positron
- Traceback to position of parent muon



Beam centroid

Measurement of evolution of centroid



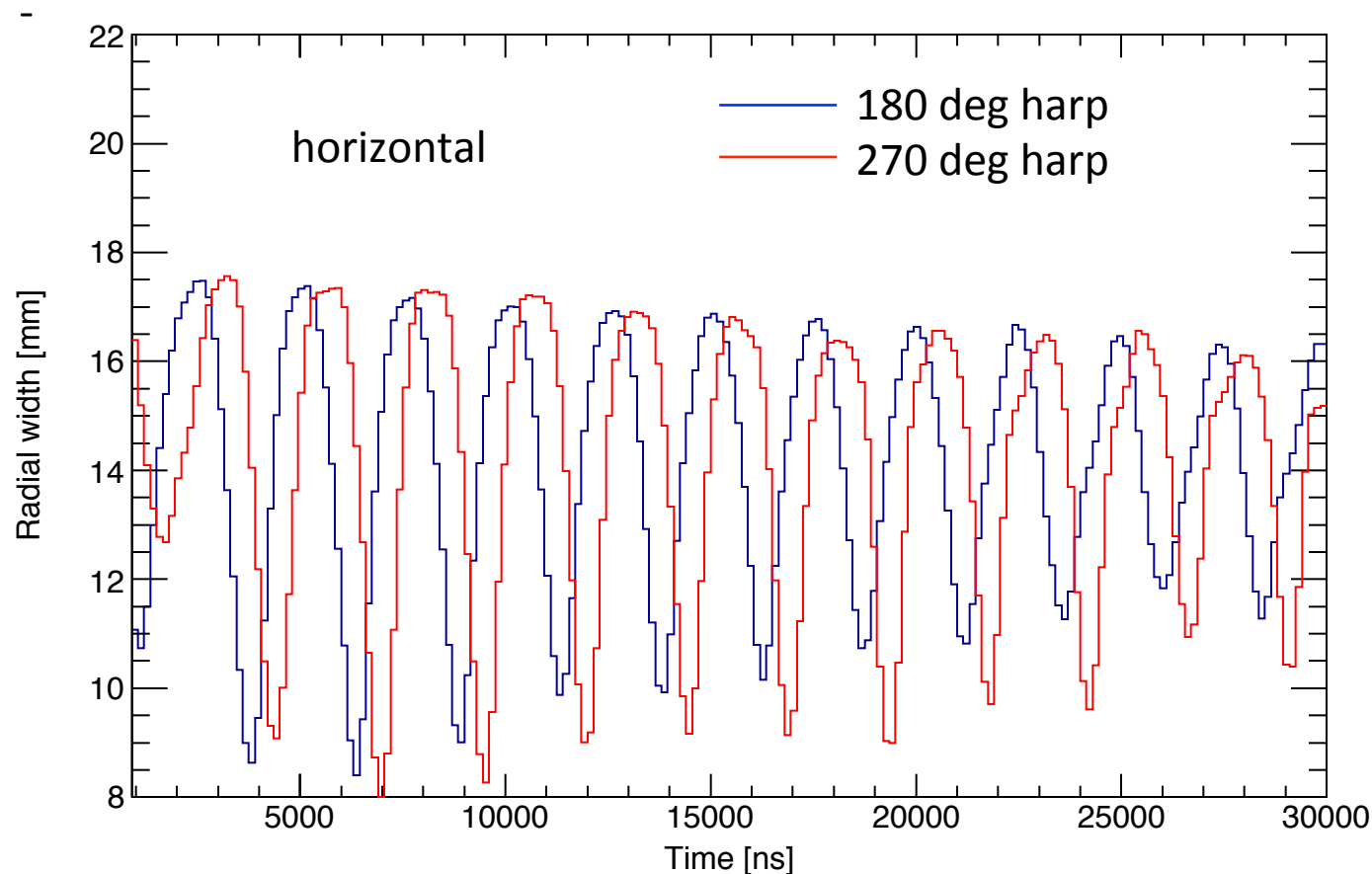


Fiber harp



Radial width / 150 ns

- Phase space and dispersion mismatch at injection
- Width modulated at Q_x and $2Q_x \Rightarrow$ dispersive and betatron components



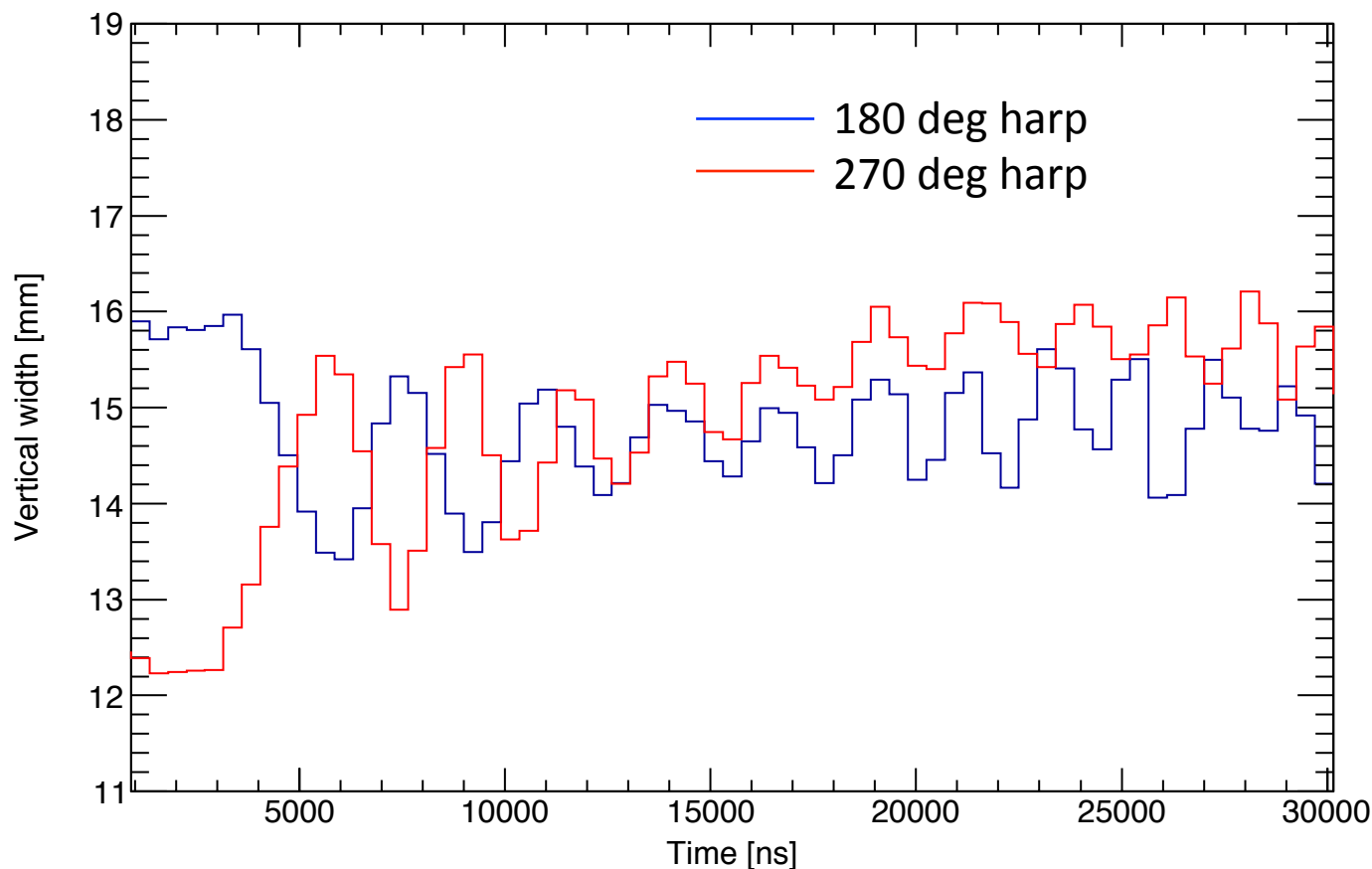


Fiber Harp



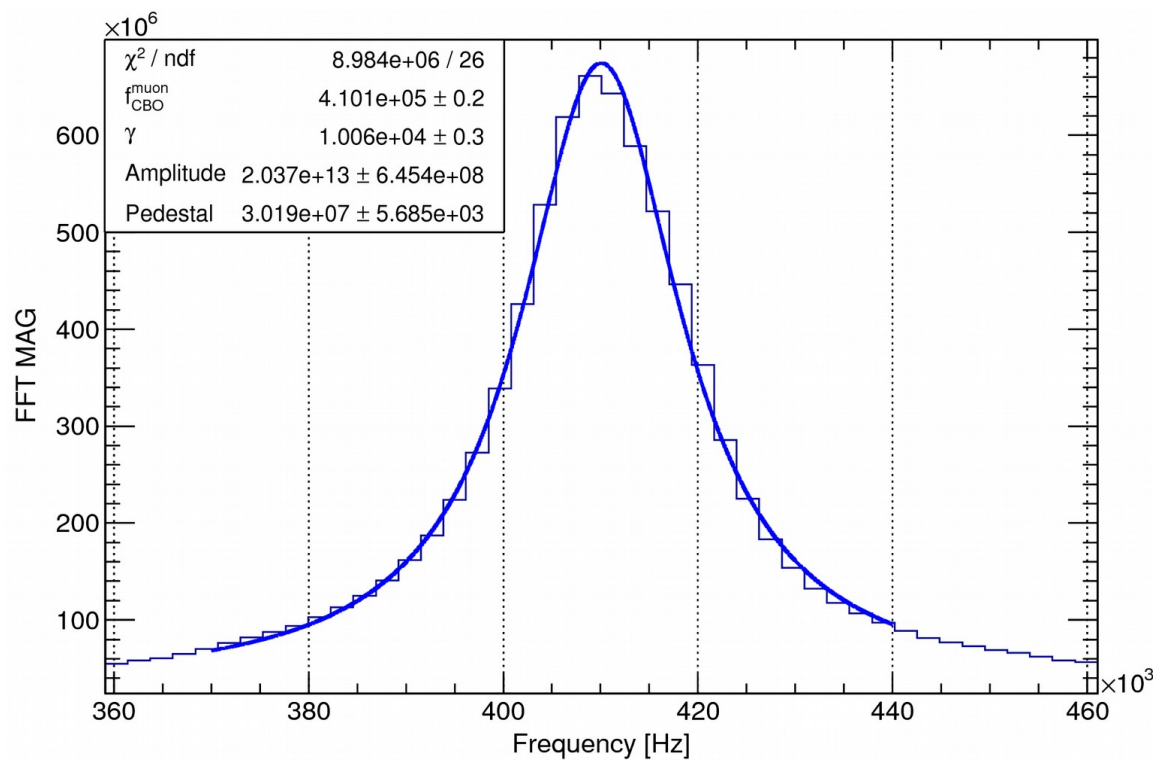
Vertical height / 150 ns

- Modulation at $2Q_y \Rightarrow$ phase space mismatch at injection
- Informs pitch correction



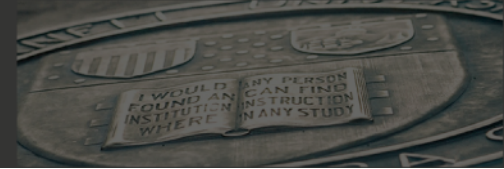


FFT of fiber harp centroid

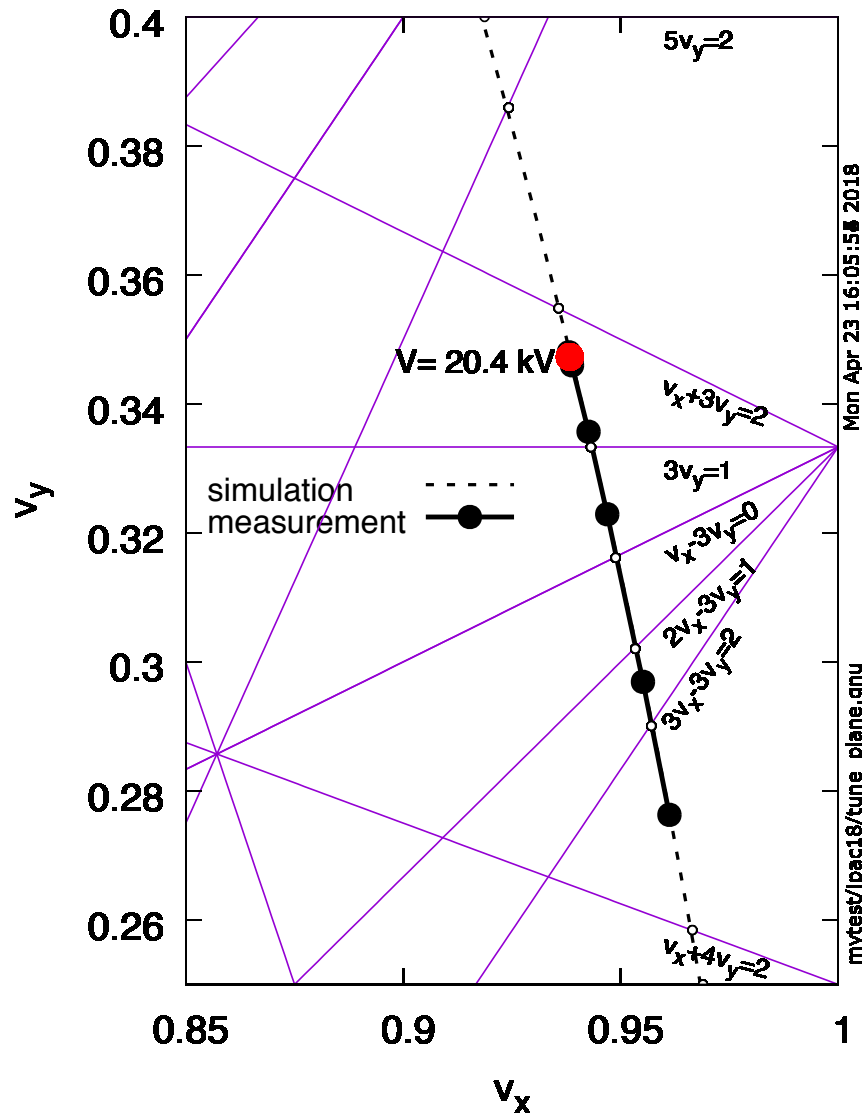




Tune scan

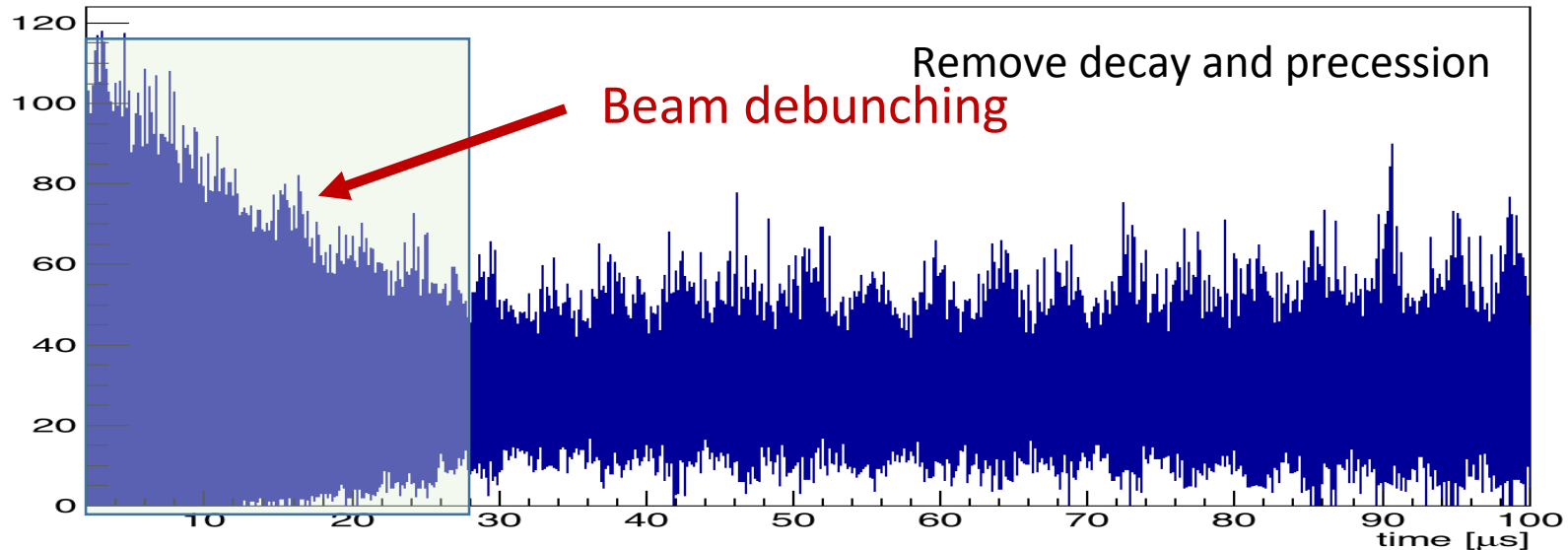
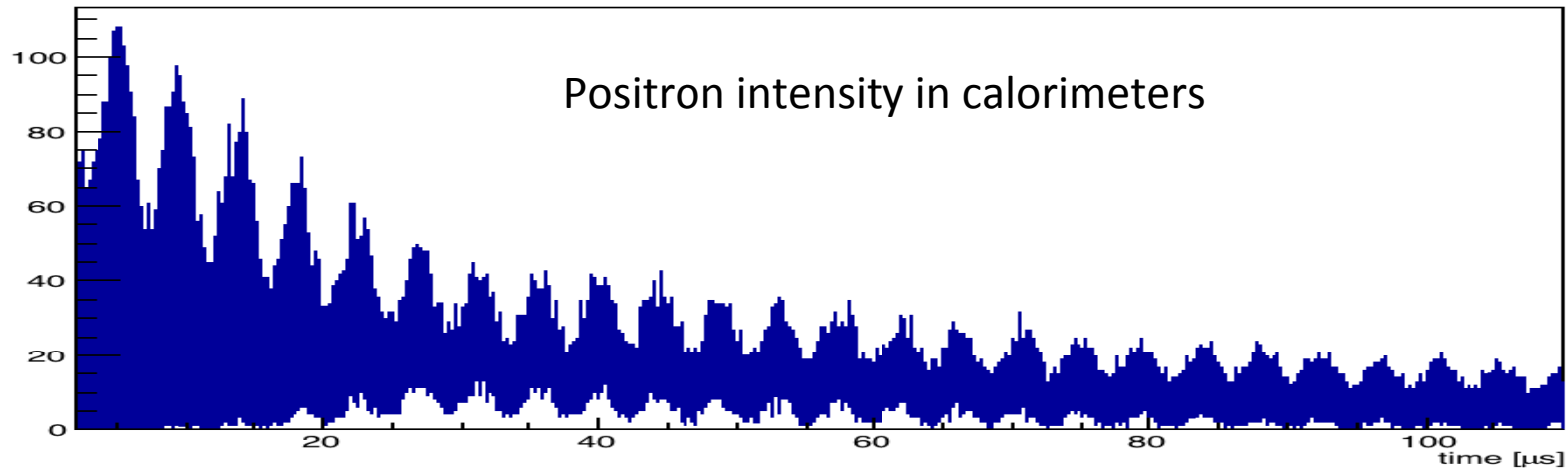


Tune along quad
voltage contour





Momentum distribution



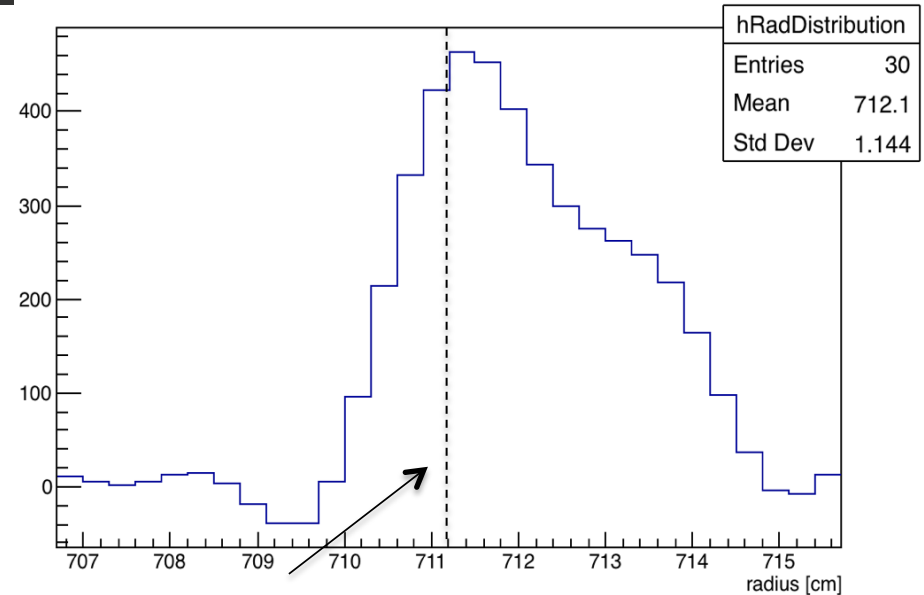


Momentum Distribution

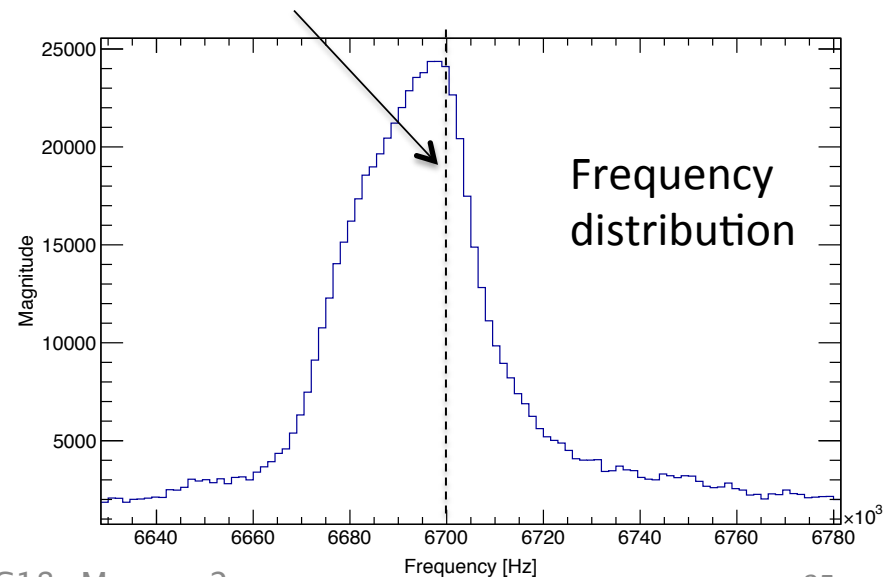
*Fit momentum distribution to
debunching signal*

Fourier analysis of fast rotation signal

Distribution of revolution frequencies **is** the
distribution of momenta. Fourier transform
of the fast rotation signal => frequency
distribution



Magic momentum

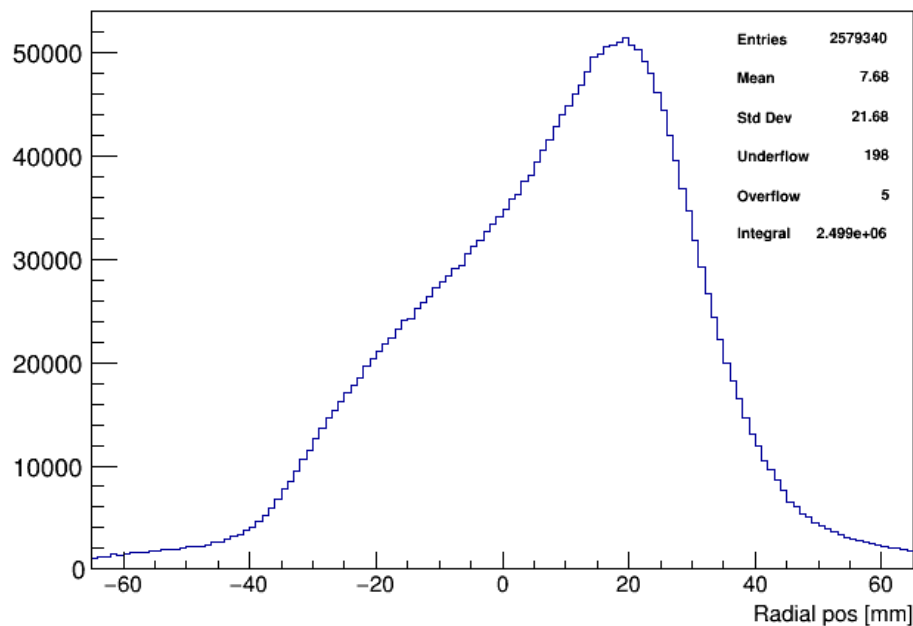




Radial distribution



Average radial distribution measured with traceback tracker
(Radial distribution convoluted with detector acceptance)

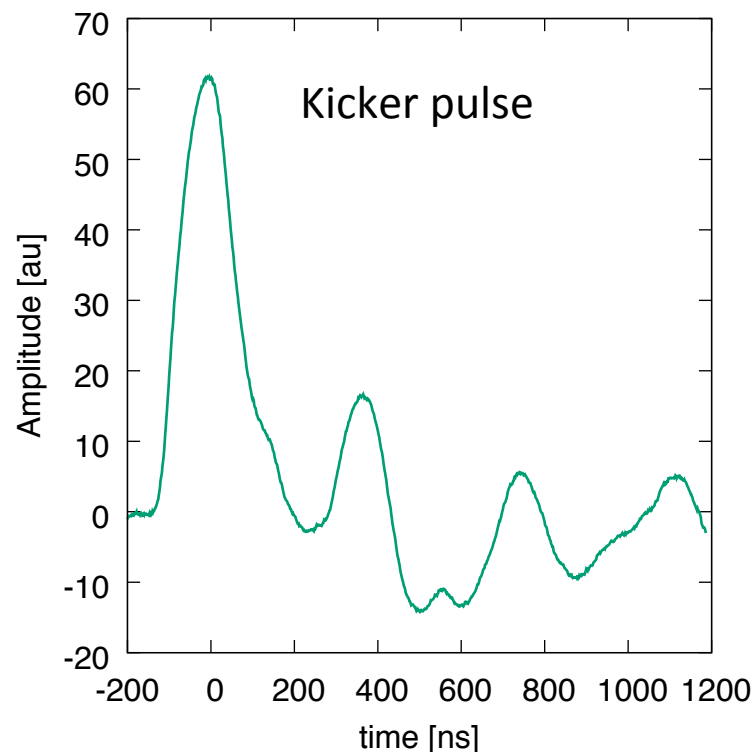
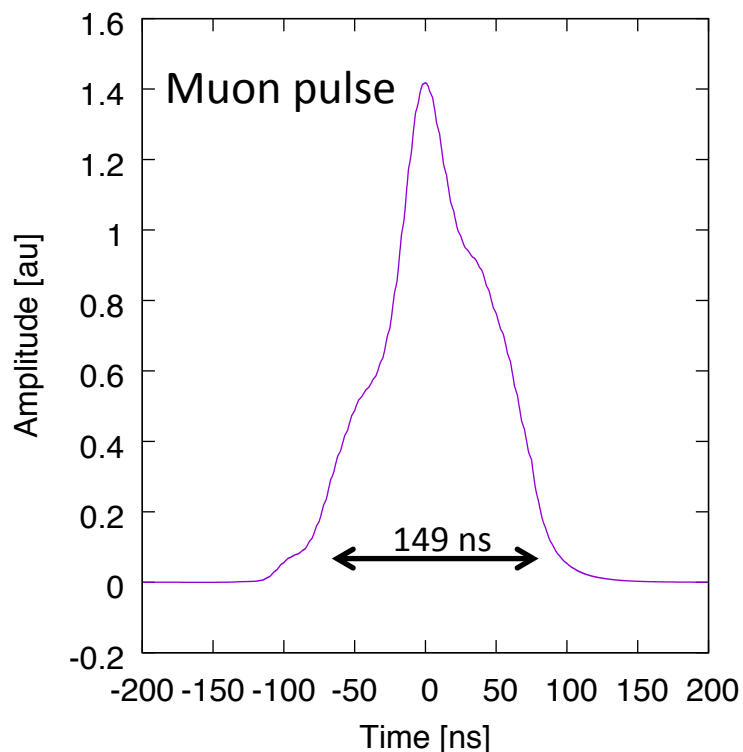




Momentum distribution

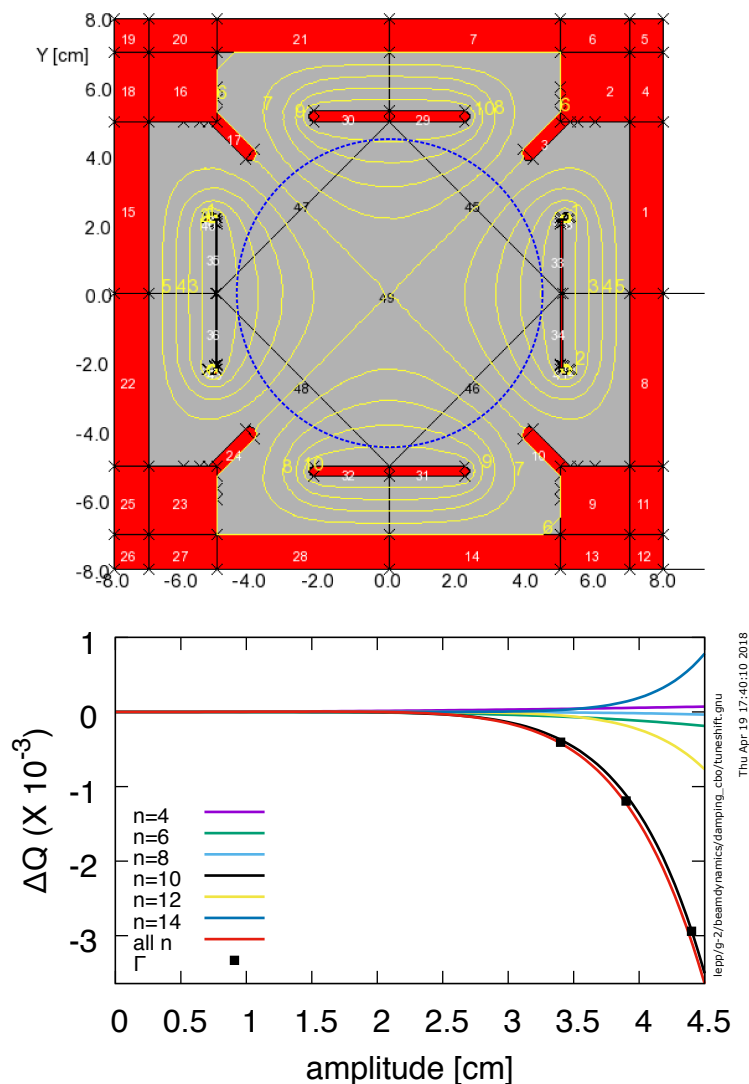
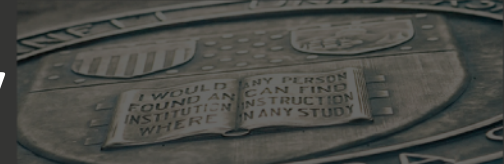
Momentum distribution is 'high'

⇒ Injection kicker is under-kicking and/or tails of the muon pulse
are in the rising and falling tails of the kick

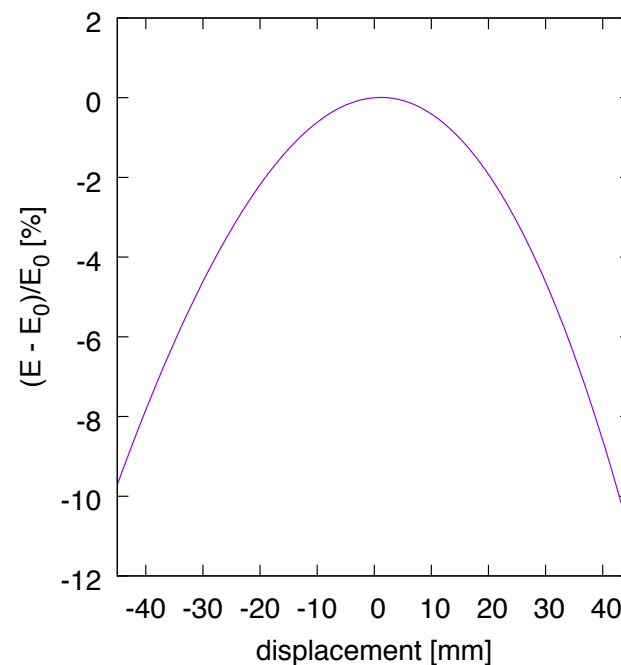




Nonlinearity



Quad field error



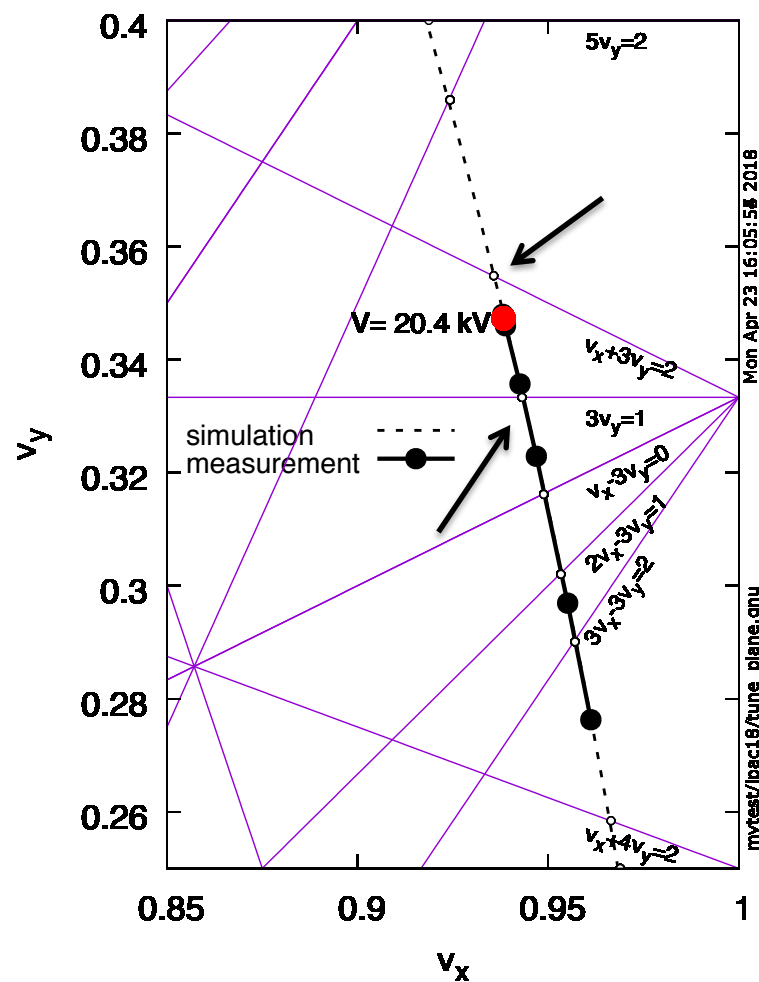
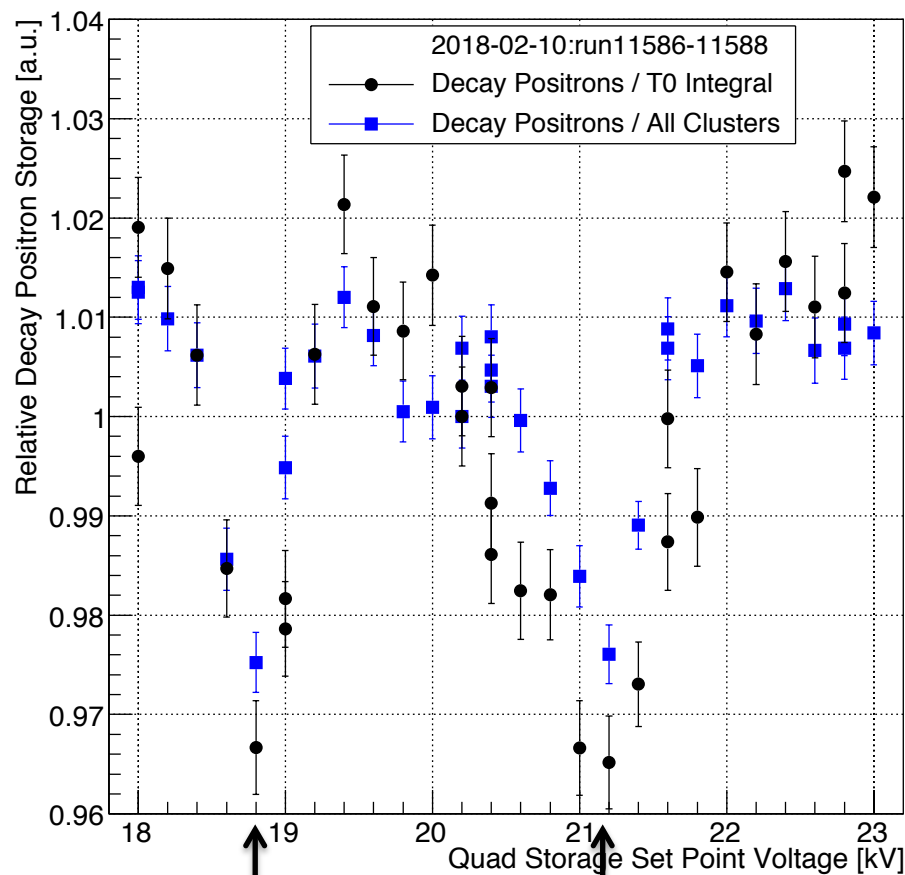


Resonances



Tune scan

Stored muons



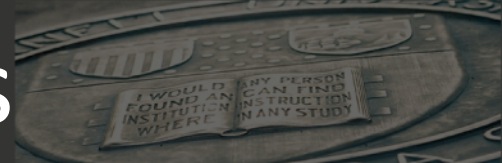
Quad voltage



Summary



- A quantitative understanding of the evolution of the muon distribution over the course of the fill is essential to limiting the systematic uncertainty in the measurement of the anomalous magnetic moment to the 70 ppb (systematic) target.
- The Muon g-2 experiment is equipped with detectors that can measure phase space and momentum distribution in some detail.
- Beam dynamics simulations informed by the measurements complete the description.
- The preliminary data presented was collected during the commissioning phase of the experiment.
- Based on that data it is clear that the fiber harp and tracker systems are an effective window on the behavior of the circulating distribution.



A.Chapelain, Cornell University, F. Gray, Regis University
S.Charity, J. Price, University of Liverpool
Joe Mott, Boston University
J.D.Crnkovic, W. Morse, V. Tishchenko, Brookhaven National Laboratory
W. Wu, University of Mississippi
K. S. Khaw, A.T. Feinberg, University of Washington

All of the Muon g-2 Collaboration (E989)

- Commissioning the Superconducting Magnetic Inflector System for the Muon g-2 Experiment – N.S.Froemming *et al.* WEPAF014
- Commissioning the Electrostatic Quadrupole System for Muon g-2 - J.D. Crnkovic *et al.* WEPAF015
- Application of Quad-Scan Measurement Techniques to Muon Beams in the Muon g-2 Experiment – J. Bradley *et al.* WEPAF016
- Lost Muon Function for the Muon g-2 Experiment – S.Ganguly & J.D.Crnkovic THPAK139
- New Fast Kicker Results from the Muon G-2 Experiment – A.P.Schreckenberger *et al.* THPML093
- Using Time Evolution of the Bunch Structure to Extract the Muon Momentum Distribution in the Muon g-2 Experiment – W. Wu, B. Quinn, J.D.Crnkovic
- Initial Studies into Longitudinal Inization Cooling for the Muon g-2 Experiment - J. Bradley *et al.*
- Correction of the Effect of the Coherent Betatron Oscillation with RF Electric Field for Muon g-2 experiment – O. Kim *et al.* THPAF010