## LARGE-APERTURE HIGH-FIELD Nb<sub>3</sub>Sn DIPOLE MAGNETS\*

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# work, publisher, and DOI. Abstract

Large-aperture high-field dipole magnets based on title of the Nb<sub>3</sub>Sn superconductor are necessary for various accelerator systems of future hadron and muon colliders. High magnetic fields and large apertures lead to large Lorentz forces and mechanical strains and stresses which can and success and incontained situates and success which can a damage brittle Nb<sub>3</sub>Sn coils. This paper describes concep- $\frac{1}{2}$  tual designs of 120-mm aperture dipoles with magnetic g fields up to 15 T based on cos-theta coils and a stress 2 management technique.

## **INTRODUCTION**

attribution Nb<sub>3</sub>Sn accelerator magnet technology has made signifiain cant progress in the past decade. Therefore, Nb<sub>3</sub>Sn dipoles and quadrupoles with a nominal field of 11-12 T are being ma considered in the near future for the LHC luminosity <sup>2</sup> upgrade (HL-LHC project) [1]. In the longer term, costeffective 15-16 T Nb<sub>3</sub>Sn magnets will be needed for the work LHC energy upgrade (HE-LHC) or the Future Circular Collider (FCC) [2].

this Large-aperture high-field dipole magnets based on of Nb<sub>3</sub>Sn superconductor are essential for various accelera-Nb<sub>3</sub>Sn superconductor are essential for various accelera-tor systems of future hadron and muon colliders. In had-ron colliders, they are needed for beam separation before and after the interaction points [3]. In a muon collider, sthey are considered for both the arc and the interaction Fregions to provide room for thick internal absorbers pro- $\hat{\infty}$  tecting magnets from muon decay products [4, 5]. Such  $\overline{\mathfrak{S}}$  magnets can also be used in test facilities to provide a <sup>©</sup> background magnetic field for testing conductor samples g or insert coils [6].

High magnetic fields and large apertures lead to large Lorentz forces and mechanical strains and stresses which can damage brittle Nb<sub>3</sub>Sn coils. Magnet design studies are being performed in the framework of the U.S. Magnet O Development Program (MDP) [7] to explore the limits of e the Nb<sub>3</sub>Sn accelerator magnet technology in general and the feasibility and limits of large-aperture magnets as outserts for a 20 T HTS/LTS dipole (MDP) in particular. acceptable for this purpose. This paper describes conceptual designs and parameters of 120-mm aperture Nb<sub>3</sub>Sn pur dipoles based on cos-theta coils with stress management. The stress management (SM) technique and limits are used also presented and discussed. þ

## MAGNET DESIGN AND PARAMETERS

work may The first design of a large-aperture dipole (Design 1), considered in this study, is based on the two outermost from this layers of the 17 T 60-mm aperture dipole described in [8]. The 2-layer outer coil has a 120 mm aperture and uses a stress management concept. Each layer consists of 5 blocks placed inside a slotted stainless steel structure with 5-mm thick radial and azimuthal walls. The coil is wound using a 40-strand Rutherford cable with 0.7 mm Nb<sub>3</sub>Sn strands.

The second design (Design 2) also has a 120-mm aperture and a 4-layer graded coil with a similar stress management concept as Design 1. All pole blocks are made of Ti alloy and are installed into the structure before winding. The 2-layer inner coil has 5 blocks in each layer wound from a 28-strand Rutherford cable with 1 mm Nb<sub>3</sub>Sn strands. The 2-layer outer coil has 5 blocks in the inner layer and 4 blocks in the outer layer. Both layers are wound as in Design 1 from the Rutherford cable with 40strands, each 0.7 mm in diameter. The cross-sections of 2layer (Design 1) and 4-layer (Design 2) coils with stress management are shown in Fig. 1.



Figure 1: Cross-sections of Design 1 and Design 2 coils with 120-mm aperture.

The 28-strand and 40-strand cables are 15.1 mm wide, have a keystoned cross-section with a keystone angle of 0.805 degree, and use Nb<sub>3</sub>Sn strands with a Cu/non-Cu ratio of 1.13 and critical current density J<sub>c</sub> at 15 T and 4.2 K of 1500 A/mm<sup>2</sup>. The mid-thickness of the 28-strand and 40-strand cables is 1.870 mm and 1.319 mm respectively.

Both designs were studied using ROXIE code [9]. The magnetic model included a cylindrical iron yoke with an outer diameter of 700 mm and a realistic B(H) curve. The main magnet parameters are summarized in Table 1 for a Nb<sub>3</sub>Sn cable J<sub>c</sub> of 1.5 kA/mm<sup>2</sup> at 15 T and 4.2 K, which corresponds to 3.0 kA/mm<sup>2</sup> at 12 T and 4.2 K.

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Parameter

Bore field, T

Peak field, T

Current, A

Number of layers

Inductance, mH/m

Stored energy, MJ/m

F<sub>x</sub>, MN/m/quadrant

Fy, MN/m/quadrant

and 4.2 K	The goal of
Design 2	maximum field stress in the co brittle Nb <sub>3</sub> Sn
4	
15.42	tures, and b) the
15.88	structure in th
12.50	length of sepa
64.94	field is 11 T o
5.07	Calculated d
12.38	Design 1 and 1
-8.01	corresponding
	T for Design 2
4.2 K is 12.1 T,	signs, the peal
t since the coil	down is in the
uality, the bore	in Design 1 an

Table 1: Magnet Parameters at SSL and 4.2 K

Design 1

2

12.10

14.18

13.06

22.47

1.92

7.59

-3.17

The maximum bore field in Design 1 at 4.2 K is 12.1 T, whereas in Design 2 it is 15.4 T. Note that since the coil cross-section was not optimized for field quality, the bore field in Design 1 is 18% lower than the peak field. In Design 2, the bore field is only 3% lower than the peak field due to a more optimum coil geometry. The sensitivity of bore field to magnet operation temperature and superconductor critical current density  $J_c$  at 12 T and 4.2 K is shown in Fig. 2. It can be seen that variations of the conductor critical current density within 50% changes the bore field by only ~10% for Design 1 and by ~8.5% for Design 2.



Figure 2: Maximum field in magnet aperture vs. conductor critical current density  $J_c$  at 12 T and 4.2 K.

### **MECHANICAL ANALYSIS**

To investigate the limitations related to the large Lorentz forces, a mechanical analysis was performed using ANSYS code and parametric models. The coil blocks are integrated into the stainless steel structure with 5 mm thick radial and azimuthal spacers and separated pole blocks. The coil blocks are allowed to separate from the structure to capture the effect of unloading under Lorentz forces. In addition, each layer could slide with respect to the adjacent layers and to the iron yoke. It was also assumed that the vertical gap in the iron yoke remains closed at all fields. The coils were pre-compressed during assembly by placing appropriate radial shims between the inner and outer coils and between the outer coil and the iron yoke. The goal of the mechanical analysis was to find the maximum field for each design at which a) the maximum stress in the coil does not exceed the acceptable stress for brittle Nb<sub>3</sub>Sn coils of  $\sim 200$  MPa at operation temperatures, and b) the separations of the coil blocks from the structure in the azimuthal direction are small, and the length of separation areas is less than the cable width. Results of the analysis have shown that for Design 1 this field is  $\sim 11$  T and for Design 2 it is  $\sim 15$  T.

Calculated distributions of the equivalent stress in the Design 1 and Design 2 coils after cool-down and at the corresponding maximum field (11 T for Design 1 and 15 T for Design 2) are shown in Figs. 3 and 4. In both designs, the peak equivalent stress in the coil after cool-down is in the pole blocks of layer 1. It reaches 185 MPa in Design 1 and 190 MPa in Design 2. In both designs, due to the azimuthal component of the Lorentz force, at the maximum field turns in each block shift towards the coil midplane. In Design 1 at a bore field of 11 T the maximum equivalent stress in the coil is ~155 MPa, and in Design 2 at 15 T it is ~190 MPa.

Lorentz forces push turns in each block towards the coil mid-plane. Under certain conditions they may separate from the support structure. In Design 1 at a bore field of 11 T, small gaps up to 4  $\mu$ m with lengths smaller than 30% of the cable width are seen in the two coil blocks of the outer-layer next to the pole. In Design 2, at 15 T the larger gaps up to 25  $\mu$ m and with lengths up to 50% of the cable width are in the 2-3 pole blocks in each layer.



Figure 3: The coil equivalent stress (Pa) after cool-down (top) and at the bore field of 11 T (bottom) of Design 1.

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Figure 4: The coil equivalent stress (Pa) after cool-down (top) and at the bore field of 15 T (bottom) of Design 2.

Calculated distributions of the equivalent stress in the coil structure after cooling-down and at the maximum bore field in the two designs are shown in Figs. 5 and 6.



from this work may be used under the terms of the CC BY 3.0 licence (@ 2018). Any distribution Figure 5: Equivalent stress (Pa) in the coil support structure of Design 1 after cool-down (top) and at 11 T bore field (bottom).



Figure 6: Equivalent stress (Pa) in the coil support structure of Design 2 after cool-down (top) and at 15 T bore field (bottom).

In both designs the maximum equivalent stress is located near the midplane of the innermost radial wall of the coil support structure. In Design 1 the maximum equivalent stress is ~550 MPa after cool-down and ~630 MPa at the maximum bore field of 11 T. In design 2 it is 435 MPa and 505 MPa respectively.

#### **CONCLUSION**

Nb<sub>3</sub>Sn dipole designs with 2-layer and 4-layer coils with 120-mm free aperture and stress management have been analyzed. Although the magnet conductor limit for state-of-the-art Nb<sub>3</sub>Sn composite wires is quite high, up to 13 T in the 2-layer design and up to ~17 T in the 4-layer design at 1.9 K, the magnet design field is limited to 11 T and 15 T respectively by mechanical considerations. Optimization of coil and stress management structure will continue to extend the operation field range in 120-mm aperture dipole magnets.

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