TIME-OF-FLIGHT, BEAM-ENERGY MEASUREMENT OF THE LANSCE 805-MHz LINAC*

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Abstract

Control of the beam-energy ramp along the length of a proton linear accelerator is required to keep the accelerator tuned according to design. Historically, the values of the field amplitudes and phases of the sidecoupled. 805-MHz LANSCE linac modules are maintained using a well-known delta-t tuning procedure [1]. Time-of-flight measurements of the proton beam energy are now also being used to confirm and improve the overall control of the energy ramp along the linac. The time-of-flight method uses absolute measurements of beam energy using direct signals from beam at an oscilloscope, as well as the difference in RF phases measured as the beam passes installed delta-t pickup loops. A newly developed BPPM data acquisition system is used. Details of the procedure and results of measurements are presented.

INTRODUCTION

The LANSCE accelerator facility is equipped with two independent injectors for H⁺ and H⁻ beams, merging at the entrance of a 201.25 MHz Drift Tube Linac (DTL). The DTL performs acceleration up to the energy of 100 MeV. After the DTL, the Transition Region beamline directs a 100 MeV proton beam to the Isotope Production Facility, while the H⁻ beam is accelerated up to the final energy of 800 MeV in an 805 MHz Coupled Cavity Linac (CCL). The H- beams, created by different time structure of a low-energy chopper, are distributed in the Switch Yard (SY) to four experimental areas: Lujan Neutron Scattering Center equipped with Proton Storage Ring (PSR), Weapon Neutron Research Facility, Proton Radiography Facility, and Ultra-Cold Neutron Facility.

The Drift Tube Linac consists of 4 tanks, which amplitudes and phases are selected through absorbercollector phase scans. Coupled Cavity Linac includes 44 accelerating modules (modules 5-48), which amplitudes and phases are tuned using the "delta-t" method [1]. This classical method is based on measurement of time of flights between two pairs of delta-t pickup loops while accelerating module is on and off. Delta-t method uses only a small phase range ($\sim 10^{\circ}$) for tuning because the procedure is based on a linear model. Delta-t tuning procedure works well when particles perform significant longitudinal oscillations within RF tanks. Accuracy of the method unavoidably drops with energy. In order to independently control tuning of the machine, the time of flight method for beam energy measurement was added.

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BEAM ENERGY MEASUREMENT BY TIME OF FLIGHT

TIME OF FLIGHT Energy measurement of the beam is usually performed through measurement of time of flight (TOF) of beam is $t = L / \beta c$ between 2 pickup loops (see Fig. 1), separated is Any by distance L, where βc is the beam velocity [2]. Time of flight can be represented as

$$t = N \frac{\lambda}{c} + \Delta t , \qquad (1)$$

Content from this work may be used under the terms of the CC BY 3.0 licence (© 2018). where N is the integer number of RF periods during time of flight, λ is the wavelength, and Δt is the fractional part of time of flight. The change of RF phase during time of flight is

$$\omega t = 2\pi N + \Delta \varphi \,, \tag{2}$$

where $\Delta \varphi = 2\pi c \Delta t / \lambda$ is an actual measured fractional part of phase change (see Fig. 2). Expressing beam velocity from Eqs. (1), (2):

$$\beta = \frac{L}{\lambda(N + \frac{\Delta\varphi}{2\pi})},\tag{3}$$

the beam energy is determined as

$$W = mc^{2}(\frac{1}{\sqrt{(1-\beta^{2})}} - 1), \qquad (4)$$



Figure 3: Time of flight measurement of absolute beam energy.



Figure 4: Measurement of time of flight of bunch train between two delta-t loops.

where mc^2 is the ion rest energy and N is the value of integer number of RF periods during time of flight. The Serror in determination of beam velocity is

$$\frac{d\beta}{\beta} = \sqrt{\left(\frac{dL}{L}\right)^2 + \left(\frac{\delta(\Delta\varphi/2\pi)}{N + \frac{\Delta\varphi}{2\pi}}\right)^2}$$
(5)
where dL/L is the error in measured distance between
pickup loops, and $\delta(\Delta\varphi)$ is the error in phase

 \overleftarrow{a} measurement, while error in determination of energy is

$$\frac{dW}{W} = \gamma(\gamma+1)\frac{d\beta}{\beta}.$$
 (6)

Separation in possible values of energy due to ambiguity in velocity determination is (N >> 1)

$$\frac{\Delta W}{W} \approx \frac{\gamma(\gamma+1)}{N} \,. \tag{7}$$

Determination of beam energy requires knowledge of Sinteger number of RF periods during time of flight. Direct observation of beam pulse trains (see Figs. 3, 4) allows measurement of absolute value of time of flight between two loops [3]. Beam velocity is determined as

$$\beta = \frac{L}{c[t - (\tau_{cable2} - \tau_{cable1})]},$$
(8)

where t is the observable time of flight at the scope, and au_{cable1} , au_{cable2} are cable lengths measured in terms of time required for signal propagation from pickup loop to the scope. Error in determination of velocity is

$$\frac{d\beta}{\beta} = \sqrt{\left(\frac{dL}{L}\right)^2 + \left(\frac{dt}{t - \Delta\tau_{cable}}\right)^2 + \left[\frac{d(\Delta\tau_{cable})}{t - \Delta\tau_{cable}}\right]^2} , \quad (9)$$

where cable length $\Delta \tau_{cable} = \tau_{cable2} - \tau_{cable1}$ the is difference.

RESULTS OF MEASUREMENTS

For beam energy measurement, we used existing Δt loops distributed along the linac. Both methods, presented above, included preliminary calibration of cable lengths, which was done using beam with known energy. Final beam energy of the linac 795.46 MeV is determined through operation of LANSCE Proton Storage Ring, which has circumference of 90.26 m and operational frequency of 2.792424 MHz. Final pair of Δt loops after linac (loops 48 - SY) is separated by the distance of $L_{48-SY} = 20.029$ m. Measured time of flight $t_{48} = 146.6$ nsec between that Δt loops with known beam velocity β_{48} = 0.84073 provides beam-based measured difference in cable length

$$\tau_{SY} - \tau_{48} = t_{48} - \frac{L_{48-SY}}{\beta_{48}c} = 67.1337 \text{ ns}$$
 (10)

For control of beam energy after every accelerating module, each module was subsequently delayed and TOF between loops 48-SY was determined. Measurements started from the last module 48 and performed until module 13, unless the signal from drifting beam was observed. Using Eqs. (4), (8), the particle velocity and beam energy were determined.

The similar technique was used to measure beam energy in the beginning of 805 MHz linac. Beam with known energy of 100 MeV after DTL was used to calibrate cable length difference of Δt loops after modules 11-12. Using distance between loops $L_{11-12} = 16.976$ m, velocity of 100 MeV particles $\beta_{100MeV} = 0.42799$, and measured TOF $t_{100 MeV} = 196.15$ ns, the beam-based measured difference in cable length of loops 11-12 is:

$$\tau_{12} - \tau_{11} = t_{100MeV} - \frac{L_{11-12}}{\beta_{100MeV}c} = 63.84528 \text{ ns}.$$
 (11)

Turning on subsequently modules 5-11, we measured TOF and beam energy after modules 5-11. The final determination of energy after module 12 was performed using measured beam energy after module 11, $E_{11} =$ 196.318 MeV, and Δt loops after modules 12-13.

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Typical error in determination of time of flight in absolute measurements is $dt = \Delta \tau_{cable} = 0.1 - 0.2$ ns. Distances between Δt loops are known with relative error of $dL/L = 5 \cdot 10^{-5}$. The error in absolute determination of energy is

$$\frac{dW}{W} \simeq (1.4 - 3.5) \cdot 10^{-3} \gamma (\gamma + 1) \,. \tag{12}$$

Absolute energy measurements were used as a reference for more precise energy determination using determination of RF phase difference in Δt loops.

A newly developed BPPM data acquisition system to control the 3D position of the beam centroid (x, y, phase) was used. As in previous method, we started measurements from the end of the linac using 48-SY Δt loops. Measured beam phase difference between loops $\varphi_{\rm sy} - \varphi_{\rm 48} = 192^{\circ}$ together with expected difference in beam phase

$$\Delta \varphi_{\text{expected}_{48}_{SY}} = 2\pi \left[\frac{L_{48-SY}}{\beta_{48}\lambda} - INT(\frac{L_{48-SY}}{\beta_{48}\lambda})\right] = 357.3^{\circ}, (13)$$

provides phase correction due to difference in cable lengths $\Delta \varphi_{corr} = \varphi_{expected 48 \text{ SY}} - (\varphi_{SY} - \varphi_{48}) = 165.3^{\circ}$. Using the RF phase difference measurement in loops 48-SY, and known value of integer RF periods for each value of energy from absolute measurements, N, the beam energy was determined in modules 35-48.

In order to perform beam energy measurement in modules 5-34, the calibration of each pair of Δt loops was done using measured beam energy from previous module while tested module was off:

$$\Delta \varphi_{corr} = 2\pi \left[\frac{L}{\beta_{in}\lambda} - INT(\frac{L}{\beta_{in}\lambda})\right] - \left(\varphi_{loop2}^{OFF} - \varphi_{loop1}^{OFF}\right), \quad (14)$$

where β_{in} is the measured beam velocity from previous module. Then, turning module on, the beam velocity after tested module was determined though measured phase difference $\varphi_{loop2}^{ON} - \varphi_{loop1}^{ON}$:

$$\beta = \frac{L}{\lambda(N + \frac{\varphi_{loop2}^{ON} - \varphi_{loop1}^{ON} + \Delta\varphi_{corr}}{2\pi})},$$
(15)

where value of N was known from absolute TOF measurement.

Error in determination of phase in considered method is $\delta(\Delta \varphi) \approx 1^{\circ}$. This method appears to be more accurate than the absolute energy measurement method. The estimated error is one order of magnitude smaller than that in absolute method:

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Table 1: Results of Beam Energy Measurements (MeV)	
of 805 MHz Linac	

module 5 6					
5	Design	Absolute Method	Error	RF Phase Method	Error
6	113.02200	112.36600	0.30020	112.63126	0.10571
	125.92400	126.31650	0.36160	125.64310	0.11520
7	139.39500	139.43210	0.42350	139.46410	0.13594
8	153.41299	152.26030	0.48800	152.67914	0.15129
9	167.95799	168.71800	0.57600	168.56892	0.18032
10	182.10600	181.66569	0.64971	181.73358	0.19748
11	196.68500	196.31799	0.73781	196.83067	0.21872
12	211.67900	211.75160	0.93160	211.42828	0.27409
13	226.37399	226.10460	1.58080	226.59654	0.30898
14	240.91701	240.16010	1.74820	241.85416	0.16909
15	255.77499	255.73830	1.94260	256.77493	0.17984
15	270.93399	270.06381	2.12999	270.95114	0.19682
17	286.00000	285.79910	2.34470	287.01511	0.21618
18	302.11700	301.34119	2.56681	301.75278	0.23536
19	317.56900	316.40970	2.79100	317.99795	0.25181
20	333.26401	330.68899	3.01201	333.71249	0.27355
21	349.19299	348.42291	3.29769	349.39008	0.29520
22	365.34698	362.83719	3.53921	365.59720	0.31939
23	381.71701	378.32379	3.80851	380.81226	0.34025
24	397.70901	395.00610	4.10950	398.26715	0.36809
25	413.88901	416.17419	4.50821	417.15771	0.39108
26	430.24899	432.56021	4.82989	433.03842	0.41787
27	446.78400	442.95270	5.03970	450.17416	0.44641
28	463.48700	461.29840	5.42160	465.09152	0.47326
29	480.35199	472.97269	5.67211	485.35034	0.51056
30	496.75000	493.65680	6.13070	502.63019	0.54392
31	513.28802	506.87000	6.43350	516.21478	0.57169
32	529.96301	525.51880	6.87430	535.40192	0.61105
33	546.76001	545.46118	7.36282	552.31592	0.64856
34	563.69897	561.35168	7.76522	570.46838	0.68841
35	580.75299	578.11871	8.20219	580.39380	0.57266
36	597.27399	601.96979	8.84691	596.78180	0.60292
37	613.90002	608.22082	9.02018	612.32794	0.63239
38	630.63000	621.09338	9.38302	625.14526	0.65725
39	647.45801	634.48639	9.76931	628.55579	0.66394
40	664.38300	655.62323	10.39597	662.19958	0.73214
41	681.40100	670.46771	10.84929	676.72748	0.76276
42	698.50897	693.96619	11.58931	690.29291	0.79196
43	715.70502	710.52197	12.12753	711.34650	0.83844
44	732.98499	719.08539	12.41091	725.06293	0.86955
45	749.66400	746.00488	13.32762	741.95288	0.90875
46	766.41699	755.41467	13.65643	761.34784	0.95493
47	783.24103	774.93658	14.35402	779.78540	1.00009
	795.46002	795,46002	15 10951	705 46002	1 02027

$$\frac{dW}{W} \simeq 1.5 \cdot 10^{-4} \gamma(\gamma + 1) \tag{16}$$

Separation in energy due to ambiguity, Eq. (7), is two orders of magnitude larger than error, Eq. (16). Table 1 8 contains results of measurements performed by both the methods. Obtained results create basis for more careful tuning of linear accelerator.

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