# STUDY OF BEAM INSTABILITIES WITH A HIGHER-HARMONIC CAVITY FOR THE HALS<sup>\*</sup>

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# Abstract

Hefei Advanced Light Source (HALS), a diffractionlimited storage ring is on the design. In HALS project, a passive higher-harmonic cavity may be added in order to increase the beam lifetime of the storage ring. When the storage ring is operated with a small momentum compaction, instabilities limit the utility of the higher-harmonic cavity. In this paper, we run an algorithm (analytic modeling) to consider the Robinson instabilities for normal and superconducting cavity respectively. The Robinson instabilities are predicted with and without mode coupling. Coupled-bunch instability induced by resonant interaction with parasitic longitudinal mode is also considered. The analytic modeling may be used to give rf-cavity parameters that are more conducive to stability. The results show that the storage ring can operate at a higher beam current and the parasitic higher-order mode of the fundamental cavity has less impact on the beam by using superconducting harmonic cavity.

#### **INTRODUCTION**

The beam energy of Hefei Advanced Light Source (HALS) is about 2.5GeV, and the beam emittance is aimed at about 50 pm-rad. Because of small momentum compaction for the diffraction-limited storage ring lattice, the bunch lengths are usually short. The intense intra-beam scattering and Touschek effects will cause evident emittance growth and lifetime reduction. The low-emittance lattice also increase the susceptibility to Robinson instabilities. Here, a passive higher-harmonic cavity is used to increase Landau damping of synchrotron oscillation and lengthen the bunch, thereby suppressing Robinson instabilities and increase the Touschek lifetime.

In this paper, we operate an algorithm to analyze the Robinson instabilities and coupled-bunch instabilities for normal and superconducting cavity respectively in the HALS. The analysis results will help us to determine the rf-cavity parameters that are useful for beam stability in HALS.

## ANALYTIC MODELING

We use the parameters of HALS shown in Table 1 and Table 2 to consider Robinson instability for a given rf voltage  $V_{T1}$ , ring current *I*, and harmonic cavity tuning angle  $\phi_2$ . We also calculates whether resonant interaction with a real parasitic impedance R<sub>3</sub> at frequency  $\sim \omega_{CB}$  will excite a dipole coupled-bunch instability. Then we calculates

Table 1: The Machine Parameters for HALS

Parameter	Value
$V_{\rm T1}$ (fundamental cavity peak	500 kV
voltage)	
$\beta_1$ (fundamental cavity rf	2.6
coupling coefficient)	
Q1 (fundamental cavity quality	20000
factor)	
$R_1$ (fundamental cavity shunt	3.2 MΩ
impedance)	<del>.</del>
$\alpha$ (momentum compaction)	$3.42 \times 10^{-5}$
$T_0$ (recirculation time)	$2.24 \times 10^{-6}$ s
N (number of bunches)	224
E (beam energy)	2.5 GeV
$\sigma_E/E$ (natural electron energy	$5.39 \times 10^{-4}$
spread)	
$V_{\rm s}$ (synchronous voltage)	217.6 kV
$\tau_L$ (longitudinal radiation	0.0332 s
damping time)	
Q <sub>3</sub> (HOM quality factor)	700
R <sub>3</sub> (HOM impedance)	$6000\Omega$
$\omega_{CB}$ (HOM angular fre-	$7.536 \times 10^9 \text{ rad/s}$
quency)	

Table 2: The RF Parameters of	of Harmonic	Cavity
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Parameter	Normal	Superconducting
ν (har- monic cavity harmonic number)	3	3
$\beta_2$ (har- monic cavity rf coupling coefficient)	0	4
Q <sub>2</sub> (harmonic cavity quality factor)	11000	$1 \times 10^{9}$
R <sub>2</sub> (harmonic cavity shunt impedance)	1.65 MΩ	$1.25  imes 10^{11} \mathrm{M}\Omega$

whether the dipole coupled-bunch modes with longitudinal mode numbers of  $\pm 1$  will cause a dipole coupled-bunch instability. The Robinson instability theory can be seen in reference [1]. The calculation procedure refers to the reference [1, 2].

The analytical results are shown in Fig. 1.

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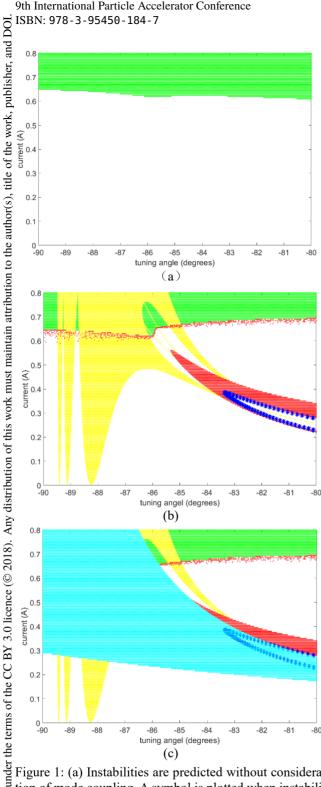


Figure 1: (a) Instabilities are predicted without consideration of mode coupling. A symbol is plotted when instabilities is predicted for a given ring current and Landau-cavity g tuning angle. Green point: zero-frequency instability. (b) ⇒Dipole-quadrupole mode coupling is include. Green point: Ξ zero-frequency coupled instability. Red point: fast modework coupling instability. Blue \*: coupled-quadrupole Robinson instability. Yellow point: instability for dipole coupledthis bunch oscillations with longitudinal mode numbers of  $\pm 1$ . rom (c) Included the resonant interaction with a longitudinal mode. Cyan point: parasitic mode coupled-bunch instabil-Content ity

In Fig. 1(a), the zero-frequency instability is predicted. The instability occurs when the currents exceeding 610 mA. In Fig. 1(b), the threshold currents of zero-frequency coupled instability is slightly increased in the larger tuning angle side compare with the zero-frequency instability. While another instability (suspect fast mode-coupling instability) occur when the currents is less than the threshold currents of zero-frequency coupled instability. This instability is plotted because the algorithm used to compute the coupleddipole and coupled-quadrupole Robinson frequencies does not converge, which may indicate a fast instability. When currents is below 560 mA, a fast mode-coupling instability is predicted. When current is below 410 mA, the coupledquadrupole instability is predicted. The instability for dipole coupled-bunch oscillations with longitudinal mode numbers of  $\pm 1$  occurs because the harmonic cavity has a low value of Q or is detuned far from the frequency  $\nu\omega_{g}$ , where  $\omega_g$  is the rf generator frequency. The parasitic mode coupled-bunch instability is shown in Fig. 3(c). The values of  $R_3$  and  $\omega_{CB}$  represent a typical higher-order mode of the fundamental rf cavity. Passive operation of the harmonic cavity is stable in the blank area in Fig. 1(c).

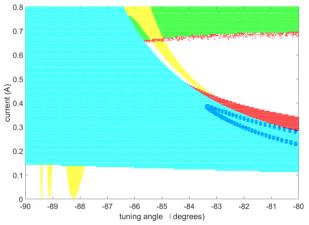


Figure 2: Instability are predicted with a HOM parasitic impedance  $R_3=12000 \Omega$ .

When the higher-order mode is not damped sufficiently, it will has a more serious impact on the beam. For  $R_3=12000\Omega$ , the analytical results are shown in Fig. 2. The parasitic mode coupled-instability is predicted in a larger area. The results manifest that the HOM must be damped strongly in order to reduce parasitic mode coupled-bunch instability.

We also analyze a superconducting harmonic cavity used the rf parameter in Table 2. The results are shown in Fig. 3. In Fig. 3, the plot range of the tuning angle is only from  $-90^{\circ}$  to  $-89.9^{\circ}$ . In Fig. 3(a), the zero-frequency instability is predicted in the range of I > 740 mA and tuning angle  $< -89.96^{\circ}$ . The quadrupole Robinson instability is happened at low current  $I \sim 20$  mA from tuning angle  $-89.926^{\circ}$ to  $-89.914^{\circ}$ . In Fig. 3(b), the zero-frequency coupled instability area become smaller compare with the uncoupled zero frequency instability. The plotted fast instability include the fast mode-coupling instability and the suspect fast mode-coupling instability. The parasitic mode coupled-bunch instability is given in Fig. 3(c). This parasitic

02 Photon Sources and Electron Accelerators

higher-order mode has a small influence on beam instability.

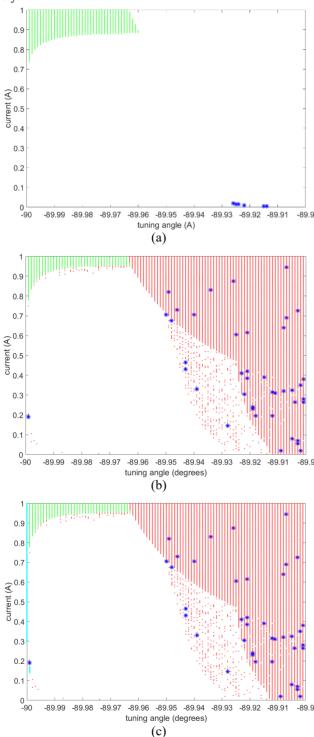


Figure 3: Instabilities are predicted for the HALS lattice with a passive superconducting harmonic cavity. The plotted symbols has the same meaning as in the Fig. 1. In Fig. 3(a) Blue \*: quadrupole Robinson instability.

Passive operation of the superconducting harmonic cavity is stable probably in tuning angle  $< -89.98^{\circ}$  and at ring current < 900 mA. Compare with normal harmonic cavity, superconducting harmonic cavity is less affected by the longitudinal cavity mode and can stable operate in a higher beam current.

For  $R_3=12000\Omega$ , the analytical results are shown in Fig. 4. The results show that the impact of the parasitic higher-order mode on the beam is slight.

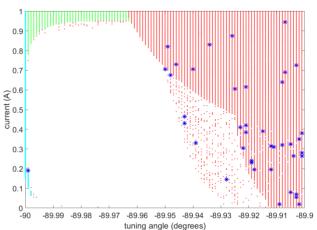


Figure 4: Instability are predicted with a HOM parasitic impedance  $R_3=12000\Omega$ .

#### DISCUSSION

We have studied instabilities when a normal and superconducting harmonic cavity is utilized respectively with the HALS lattice, whose momentum compaction is small. The results show that by using superconducting harmonic cavity, the storage ring can operate at higher beam current and the parasitic higher-order mode of the fundamental cavity has less impact on the beam. The stable operation area was obtain. We can change the rf parameters of the rf cavities to get a larger stable operation area. The analytic modeling can help us to get rf parameters, which are better for stable operation at higher beam current. We will study instabilities when the superconducting fundamental cavity with the higher harmonic cavity are utilized in the future.

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