Recent Progress on the MaryLie/IMPACT Beam Dynamics Code

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in collaboration with:

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What is MaryLie/IMPACT

A 3D parallel PIC code that

- combines 5th order MaryLie beam optics with
- IMPACT space charge & RF cavity modeling
- and augments these with new capabilities

Multiple capabilities in a single environment

 Map generation and analysis, particle tracking, envelope tracking, fitting/optimization

Combines the functionality of high order optics code (MaryLie, others) and space-charge code (IMPACT, PARMILA) and rms envelope code (e.g. TRACE3D) in a single unified environment

Other codes combine high order optics w/ space charge

• Synergia (talk by P. Spentzouris on Thursday)



- Note that the rapidly varying s-dependence of external fields is decoupled from slowly varying space charge fields
- Leads to extremely efficient particle advance:
 - Do not take tiny steps to push ~100M particles
 - Do take tiny steps to compute maps; then push particles w/ maps

Particle-In-Cell (PIC)



What to do after a map is created?

- Concatenate with previous maps
- Use it to track particles
- "Sandwich" with the previous map
- Use it to track envelopes

Some ML/I features

- 5th order optics
- 3D space charge
- realistic RF fields
- acceleration
- ref traj computed "on the fly"
- 5th order rf cavity model (D. Abell)
- multi-reftraj for rf cavities
- envelope tracking
- soft-edged magnets
- methodical treatment of units
- MAD-style input compatibility
- "automatic" commands

The second	s Code with Space Charge
A l'araner Deam Dynamic	s code with space charge
based on the MaryLie Lie Alg	gebraic Beam Transport Code
and the IMPACT Paralle	el Particle-In-Cell Code *
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Treatment of Units in

IVIL/I ML/I allows arbitrary specification of 1, ω , δ , where **6-vector = (x/1, p_x/\delta, y/1, p_y/\delta, \omegat, p_t/\omega1\delta)**

Common choices are:

Magnetostatic systems: δ=p₀, 1 =1, ω 1 /c=1
 Systems w/ acceleration: δ=m₀c, ω=ω_{bunch}, 1 =c/ω

New UNITS command:

myunits, units: l=1.0, p=0.8, w=2.856e9
myunits, units: type=magnetic, l=1.0
myunits, units: type=dynamic, l=c/w
Etc.

New Commands (selected)

autoslice: automatic slicing of thick elements

SLICES = # of slices, L = distance between slices, CONTROL = local/global/none

autoapply: automatic application of a commands

- NAME = name of menu element or line
- autotrack: automatic tracking of particles
 - TYPE=<taylorN, symplecticN>, envelope=true/false

autoconcat : automatic concatenation of maps

poisson: select/set parameters of Poisson solver

• NX=,NY=,NZ=,ngridpoints= *fixed/variable*, boundingbox=*fixed/variable*...

units: specification of units

• TYPE=<static, dynamic>, L = scale length, P = scale momentum, F = scale freq, W = scale angular freq, T = scale time

ML/I Data Handling and I/O

Uses headers in particle data files (developed under SciDAC w/ FNAL) to ease macroparticle input from codes with different units

- Descriptive text, scale length, scale momentum, scale time
- Frees the user from the headache of unit conversion

H5Part (implementation in ML/I under development)

- Parallel I/O crucial in very large scale simulations
- See presentation by A. Adelmann, PSI
- Working w/ Tech-X (D. Abell, K. Paul) to develop standard for map I/O
 - Scale length, momentum, time
 - Initial reference trajectory
 - Final reference trajectory
 - Map coefficients

RMS Envelope Calculations

based on observation that envelope Hamiltonian is of the form H=H_{ext}+H_{sc+emit}

envelope tracking can be done using the same split-operator functionality in ML/I as is used for particle tracking

odotrack, autotrack: envelope=true

Subroutine to update the envelopes due to H_{ext} is just a matrix-vector multiply

Subroutine to compute "envelope kick" from H_{sc+emit} involves calling math library routine to compute integrals in 3D envelope equations

 easily incorporated into other split-operator based spacecharge simulation codes

Code for computing "envelope kick"

```
uu=env(1)**2
vv=env(3)**2
ww=(env(5)*gamma*beta*clite/(w*xl))**2
```

```
call scdrd(uu,vv,ww,g311,g131,g113)
```

```
qtot=bcurr/bfreq
xi0=xmc2*clite/(q*fpei)
xlam3=1.d0/(5.d0*sqrt(5.d0))
qcon=1.5d0*qtot*xlam3*clite/(xi0*(beta*gamma*xl)**2)*p0/xp
tcon=qcon*(gamma*beta*clite/(w*xl))**2
```

```
r11= qcon*g311*env(1) + emxn2*(xp/(p0*xl))/env(1)**3
r33= qcon*g131*env(3) + emyn2*(xp/(p0*xl))/env(3)**3
r55= tcon*g113*env(5) + emtn2/env(5)**3*(xl*xp*xw*2*xmp0**2/p0)
```

```
env(2) = env(2) + r11*tau
env(4) = env(4) + r33*tau
env(6) = env(6) + r55*tau
```

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(developed in collaboration w/ A. Adelmann, J.

- Amundson, P. Spentzouris)
 - KV beam in a FODO channel
 - Free expansion of a cold, uniform density bunch
 - Cold beam in a FODO channel with RF cavities
 - Thermal beam in a constant focusing channel
 - Bi-thermal beam in a constant focusing channel

Channel with RF



Plot of xrms, yrms, and 0.4*trms vs. distance in 1 period of a FODO channel with rf cavities. ML/I results (symbols) are on top of curves obtained from the rms equations

#comments

Transport in a quad channel w/ rf cavities and space charge

#menu beam: beam,particle=proton,ekinetic=.250,bfreq=7.d8,bcurr=0.1d0 units: units, type=dynamic, l=1.0d0, w=2.d0*pi*700.d6

dr: drift, l=0.10 slices=4
fquad: quadrupole, l=0.15 g1=6. lfrn=0. tfrn=0. slices=6
dquad: quadrupole, l=0.30 g1=-6. lfrn=0. tfrn=0. slices=12

gapa1: rfgap,freq=7.d8,phasedeg=45.,file=f1,steps=100, slices=5
gapb1: rfgap,freq=7.d8,phasedeg=-1.,file=f2,steps=100, slices=5

pois: poisson, nx=64,ny=64,nz=128 !solver parameters

raysin: raytrace, type=readonly file1=partcl.data

dump: particledump,file=adump,sequencelength=100,precision=6

slice: autoslice, control=local

```
dotrack:autotrack, type=taylor1
post: autoapply, name=prntall
prntmoms: moments, precision=9, nunits=1
prntref: reftraj, precision=9, nunits=1
prntmax: maxsize, precision=9, nunits=1
```

```
cell,line=(fquad dr gapa1 dr dquad dr gapb1 dr fquad dump)
linac, line= 20*cell
prntall, line=(prntmoms,prntref,prntmax)
```

#labor	
slice	!slice elements
pois	!set poisson solver parameters
raysin	!read in some rays
prntall	<pre>!print moments and reference trajectory</pre>
dotrack	!tell the code to do autotracking
post	lafter every slice, print moments
linac	
end	

Recent advances in RF cavity maps

5th order from field data on cylinder (D. Abell, Tech-X)

Multiple reference particles in longitudinal phase space



Algorithmic Advances

Integrated Green function

 Solves the long-standing problem that has plagued certain codes (e.g. PARMELA, original IMPACT) when grid aspect ratios become large

Shifted Green function

- Originally developed for long-range beam-beam
- Same algorithm works for cathode image effects
- Wavelet-based methods
- Multi-level gridding, AMR

addresses a Critical Issue: high aspect

ratios

Some Poisson solvers used in static electric and gravitational particle-in-cell simulations lose accuracy when the grid aspect ratio >> 1

Some important problems involve extreme aspect ratios:

- Long beams in rf accelerators; pancake beams
- Beams in induction linacs: L~ 10s of meters; R ~ cm
- Galaxies
- Standard grid-based approaches involve using a very large # of grid points in the long dimension, leading to prohibitively long run times
 - As a result, it is *extremely* difficult model high aspect ratio systems accurately using standard grid-based approaches

101 approach recognizes mai veriam

physical quantities may vary on vastly

different scales
 The Green function, G, and source density, p, may change over different scales

G is known apriori; ρ is not

We should use all the information available regarding G so that the numerical solution is only limited by our approximate knowledge of ρ

Example: 2D Poisson equation in free space

$$\phi(x,y) = \int G(x-x',y-y')\rho(x',y')dx'dy'$$

(Hockney and
Eastwood)
$$\phi_{i,j} = \sum G_{i-i',j-j'} \rho_{i',j'}$$

- This approach is equivalent to using the trapezoidal rule (modulo treatment of boundary terms) to approximate the convolution integral
- This approach makes use of only partial knowledge of G
- The error depends on how rapidly the integrand, ρG, varies over an elemental volume
 - If ρ changes slowly we might try to use a large grid spacing; but this can introduce huge errors due to the change in G over a grid length

IGF Algorithm

- Assume the charge density, ρ , varies in a prescribed way in each cell
- Use the analytic form of the Green function to perform the convolution integral <u>exactly</u> in each cell, then sum over cells
- Example: linear basis functions to approximate ρ in a cell:

$$\phi(x_{i}, y_{j}) = \sum_{i', j'} \rho_{i, j} \int_{0}^{h} dx' \int_{0}^{h} dy'(h_{x} - x')(h_{y} - y')G(x_{i} - x_{i'} - x', y_{j} - y_{j'} - y') +$$

$$\sum_{i', j'} \rho_{i+1, j} \int_{0}^{h} dx' \int_{0}^{h} dy' x'(h_{y} - y')G(x_{i} - x_{i'} - x', y_{j} - y_{j'} - y') +$$

$$\sum_{i', j'} \rho_{i, j+1} \int_{0}^{h} dx' \int_{0}^{h} dy'(h_{x} - x')y'G(x_{i} - x_{i'} - x', y_{j} - y_{j'} - y') +$$

$$\sum_{i', j'} \rho_{i+1, j+1} \int_{0}^{h} dx' \int_{0}^{h} dy'x'y'G(x_{i} - x_{i'} - x', y_{j} - y_{j'} - y') +$$

Shifting the indices results in a single convolution involving an integrated effective Green function:

$$\phi_{i,j} = \sum G_{i-i',j-j'}^{eff} \rho_{i',j'}$$

Improvement over Standard

Applic Herental integrals can be done analytically; formulas are very lengthy

- Requires more FLOPS than simply using G_{ii} but...
- In situations where the grid is <u>fixed</u>, this only needs to be done <u>once</u> at the start of a run. Amortized over many time steps, this does not significantly impact run time.
- Accuracy: Method works as long as the elemental integrals are computed accurately and as long as the grid and # of macroparticles are sufficient to resolve variation in ρ

 IGF maintains accuracy even for extreme aspect ratios (>1000:1)

> As a result, IGF performs orders of magnitude better than the standard convolution algorithm for realistic problems involving large aspect ratios

Example: 2D gaussian ellipse

- Aspect ratio is 1:500 -- x_{max}=0.002, y_{max}=1
- Calculation of fields using (1) standard Hockney algorithm and (2) IGF approach
 - In both cases, performed convolutions for the fields directly (rather than calculating the potential and using finite differences to obtain fields)

Calculation performed using a mesh of size

- Hockney: 64x64, 64x128, 64x256,..., 64x16384
- IGF: 64x64

IGF field error





Simulation of a high-aspect ratio bunch using an integrated Green Function (IGF) and a conventional algorithm (Hockney). IGF on a 64x64 grid (purple) is more accurate than a standard calculation using 64x2048 (blue), 64x4096 (green), and 64x8192 (red).

IGF summary

- For the 2D Gaussian test problem, the standard Hockney algorithm would require ~500 times more computational effort to achieve the same worst-case accuracy as a simulation using the IGF approach.
- IGF works whether the aspect ratio is large, small, or near unity, i.e. it is generally applicable.
- 2D implemented
- 3D implemented w/ constant basis function
- **3D w/ linear basis functions leads to messy formulae**
 - Collaborating w/ D. Abell (Tech-X) to produce improved implementation

Extension of IGF to Beams in Pipes

- IGF is especially useful when applied to beams in pipes, since the Green function falls off exponentially in z, though ρ(z) may change slowly over meters
- Due to shielding in beampipe, sum can be truncated in the "long" direction:

$$\phi_{i,j} = \sum_{i'=1}^{N_x} \sum_{j'=j}^{j \pm j_{cutoff}} G_{i-i',j-j'}^{eff} \rho_{i',j'}$$

If grid length in z is >> pipe radius, can truncate at nearest neighbors:

$$\phi_{i,j} = \sum_{i'=1}^{N_x} (G_{i-i',j-1}^{eff} \rho_{i',j-1} + G_{i-i',j}^{eff} \rho_{i',j} + G_{i-i',j+1}^{eff} \rho_{i',j+1})$$

- For a rectangular pipe, can rewrite Green function as a sum of convolutions and correlations; can still use FFT-based approach to sum over elements
- Applicability to circular pipes is still an open problem

Conclusion

Under SciDAC we have developed a multiphysics 3D parallel framework

- Space charge
- Nonlinear optics
- Wakes (SciDAC/BNL, tested in Synergia)
- Improved algorithms
- Test suite
- Available for use through A. Dragt and R. Ryne
- Questions: rdryne@lbl.gov