

# A 1d Time-Dependent Theoretical Model of X-Ray Free-Electron Laser Oscillator

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  - Gain calculation
  - Cavity model
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# Background

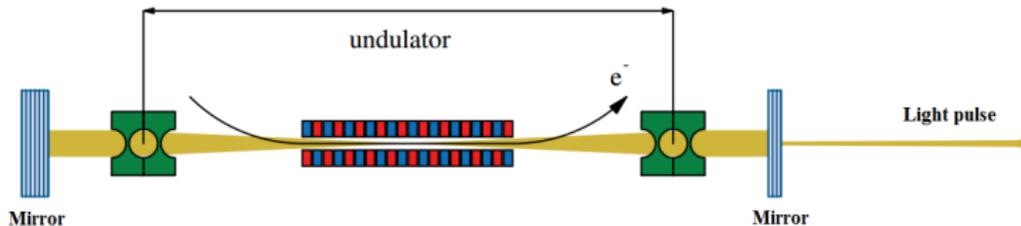
# Background

Two main ways to generate X-ray pulses by FEL: self-amplified spontaneous emission(SASE) and X-ray free-electron laser oscillator(XFEL).

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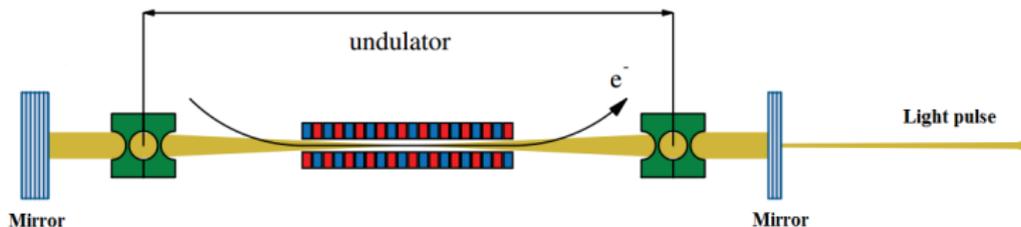
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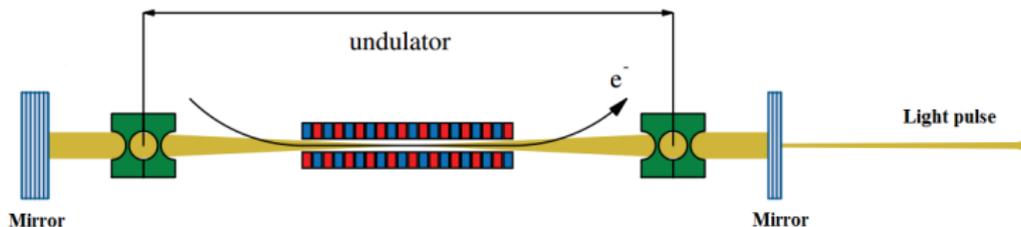


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- 2008 Kwang-Je Kim made a proposal for XFELO.
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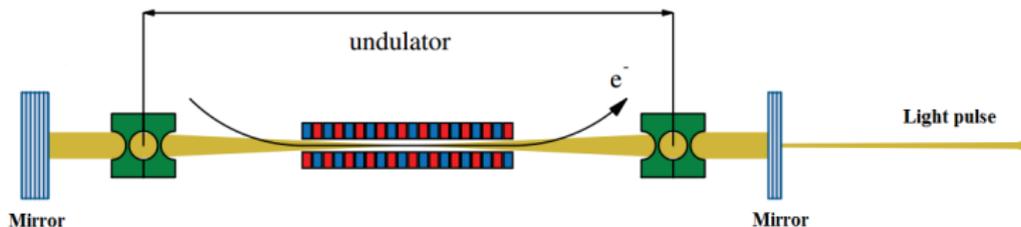
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Challenges:

- High repetition electron injector.
- Heat loading of the Bragg reflection crystal mirror.
- Crystal mirrors alignment.
- Time-consuming numerical simulation.

# Motivation

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- The single-pass gain is calculated theoretically.
- Matlab codes for the cavity model.
- It takes a few minutes for a fully simulation.
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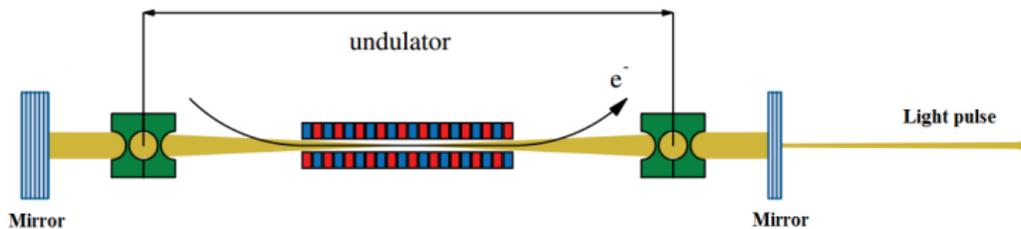
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The new approach takes few minutes which makes the theoretical analysis of single-pass gain, power growth, time-dependent laser profile evolution and cavity desynchronization become more efficiently.

# Theoretical model of FEL0

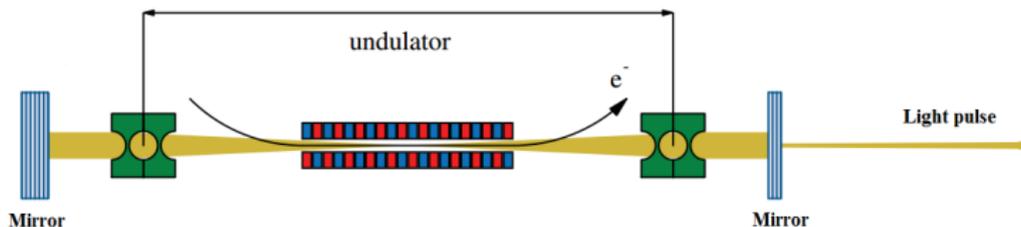
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Three main steps for FEL0 new model:

- Calculating the single-pass gain theoretically.
- Initializing the start-up radiation field.
- Simulating the laser power and profile transformation.

We take the advantages of electron distribution density function to get the single-pass gain. The motion of single electron in the phase space  $(\theta, \eta)$  is described by the pendulum equation:

$$\begin{aligned}\frac{d\theta}{dz} &= 2k_u\eta \\ \frac{d\eta}{dz} &= -\frac{\epsilon}{2k_u L_u^2} \sin\theta\end{aligned}$$

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Combining the equations above yields the following partial differential equation

$$\frac{\partial\rho}{\partial z'} + \eta' \frac{\partial\rho}{\partial\theta} + \sin\theta \frac{\partial\rho}{\partial\eta'} = 0$$

# Gain calculation

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<sup>1</sup>Boscolo I, et al,. IEEE Journal of Quantum Electronics, 1982, 18(11):1957-1961.

Assuming the initial condition of electron beam with a Gaussian distribution, the solution can be found by the method of characteristics<sup>1</sup>

$$\rho = \frac{1}{2\pi} \frac{1}{\sqrt{2\pi}\sigma_{\eta'}} \times \exp \left\{ -\frac{1}{2\sigma_{\eta'}^2} \left[ \frac{\eta' \operatorname{cn}(z'; C) - \sin\theta \operatorname{sn}(z'; C) \operatorname{dn}(z'; C)}{1 - \cos^2 \frac{\theta}{2} \operatorname{sn}(z'; C)} - \eta'_0 \right]^2 \right\}$$

where  $C^2 = \frac{\eta'^2}{4} + \cos^2 \frac{\theta}{2}$ .

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Using the law of conservation of energy, the single-pass power gain is

$$G = \sqrt{m_e c^2 K [JJ] k_u^{-1}} \frac{I}{c\beta} \frac{1}{2\pi \sum^2} \frac{1}{\epsilon_0 E_0^{3/2}} \langle \Delta\eta' \rangle$$

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Taking account of three dimensional effects, we use the equivalent formula

$$\frac{\sigma'_E}{E_0} = \sqrt{\left( \frac{\sigma_E}{E_0} \right)^2 + \left( \frac{\varepsilon \lambda_u}{4\lambda\beta} \right)^2}$$

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# Cavity model

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In this way, the complex reflectivity is simplified as

$$r(y) = \begin{cases} y - \sqrt{y^2 - 1} & \text{if } y > 1 \\ y - i\sqrt{1 - y^2} & \text{if } |y| \leq 1 \\ y + \sqrt{y^2 - 1} & \text{if } y < -1 \end{cases}$$

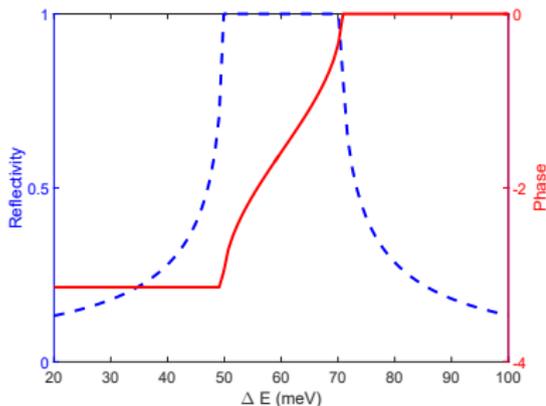
where  $y = \frac{1}{|\chi_H|} \left[ \frac{2(E - E_H)}{E_H} + \chi_0 \right]$ ,  $E_H$  is the Bragg energy and  $\chi_0$  and  $\chi_H$  are Fourier components of the dielectric susceptibility of the crystal.

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**Figure:** The complex reflectivity of Bragg crystal at various incident photon energy deviation from Bragg energy.

# XFEL0 parameters

A typical XFEL is studied using parameters shown in Table below:

Parameter	Value	Unit
Beam energy $E_0$	7	GeV
Energy spread $\sigma_E$	1.4	MeV
Normalized emittance $\varepsilon_n$	0.2	$\mu\text{m}\cdot\text{rad}$
Peak current $I$	10	A
Electron bunch length $\sigma_t$	1.0	ps
Undulator period $\lambda_u$	17.6	mm
Number of undulator $N_u$	3000	
Laser wavelength $\lambda$	0.1	nm
Cavity loss	5%	
Bragg mirror reflectivity $R$	94%	

# Electron distribution and gain function

## Electron distribution and gain function

The electron is trapped in the “bucket” and transform its energy to light like in the IR FEL case. However, the bucket which traps the electron becomes flatter, and the energy modulation is smaller due to the relative larger electron energy.

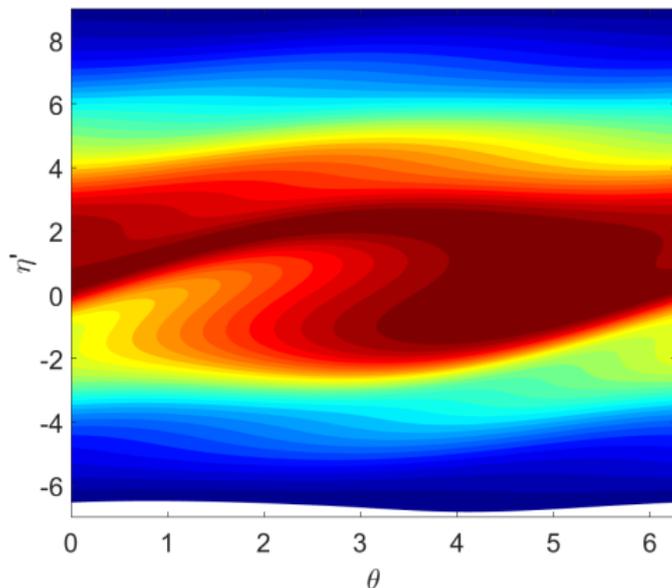
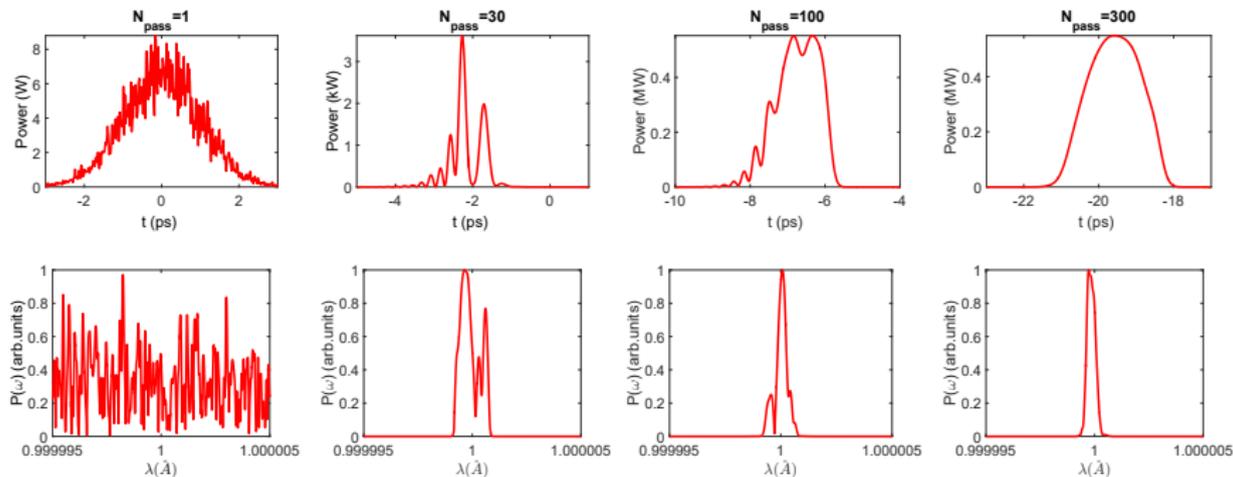


Figure: The electron density distribution function in phase space of one slice.

# Light power and profile

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The radiation in the cavity start from shot noise which has significant fluctuations, becomes smooth as passing number increases, and finally reaches saturation and remain steady state.



**Figure:** Snapshots of output radiation pulse for a typical X-ray FEL at 1.0 Å. The top and the bottom row show the longitudinal pulse temporal profile and corresponding spectrum respectively.

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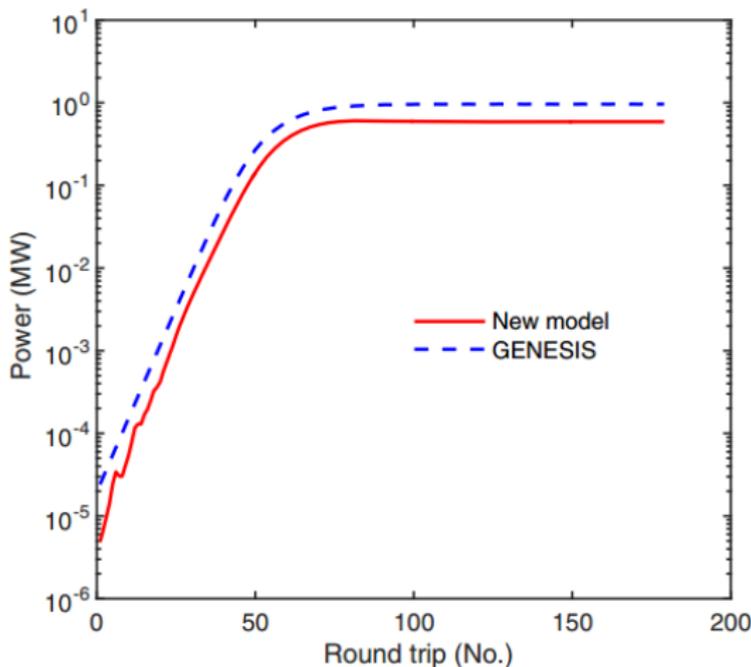
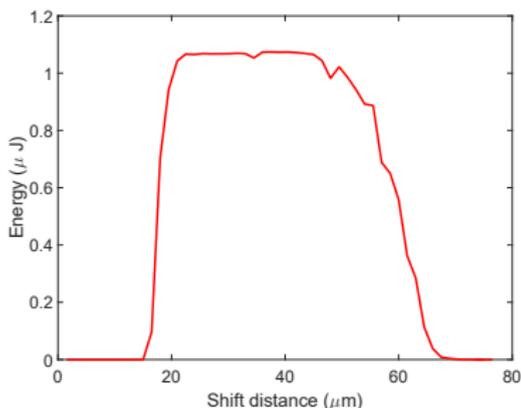


Figure: The enhancement of output laser peak power with various passes  $N_{pass}$ .

# Cavity desynchronization

The complex reflectivity of the crystal mirrors causes an extra phase shift of optical field and leads to the pulse slides backward. In the theoretical model, the electron beam is constantly delayed a distance to overlap with the optical field.



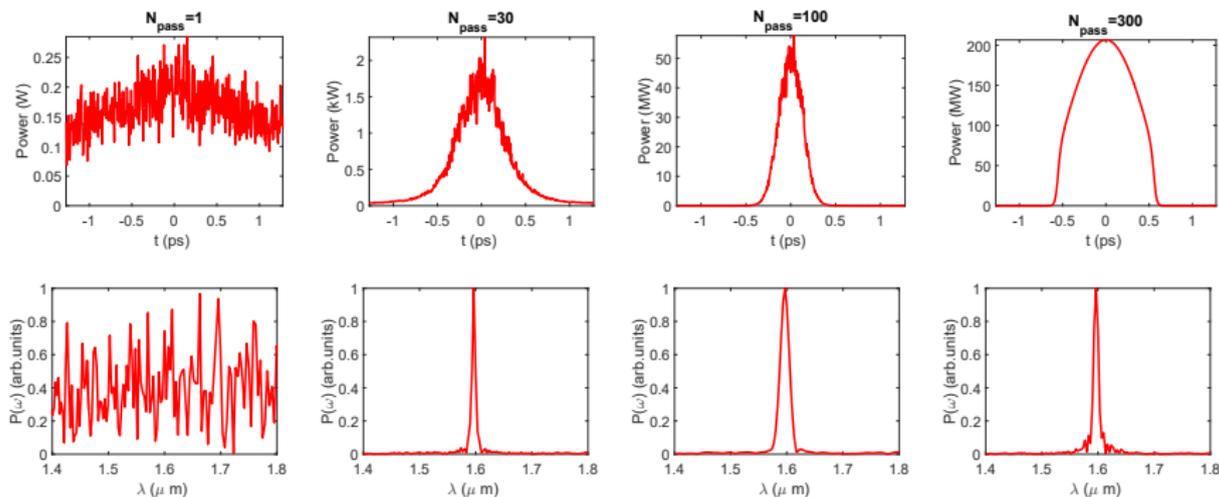
**Figure:** The output laser energy as a function of desynchronization.

**Figure:** The electron beam and light power profile.

# Infrared FEL0 Light power and profile

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The radiation in the cavity start from shot noise which has significant fluctuations, becomes smooth as passing number increases, and finally reaches saturation and remain steady state.



**Figure:** Snapshots of output radiation pulse for a typical infrared FEL0 at  $1.6\mu\text{m}$ . The top and the bottom row show the longitudinal pulse temporal profile and corresponding spectrum respectively.

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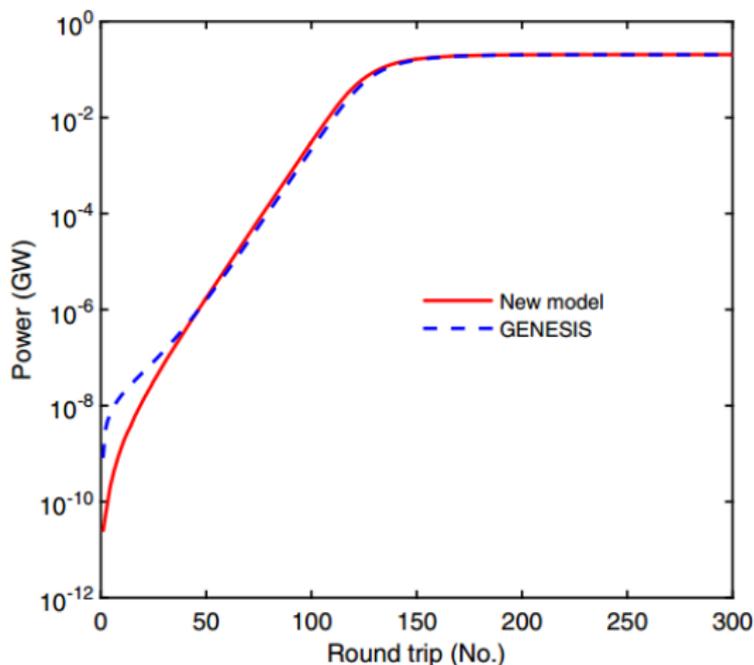


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# Primary ideas

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Solution<sup>2</sup>:

- Solving the electron density partial differential equation to get single-pass gain.

## Gain function

$$G = \sqrt{m_e c^2 K [JJ] k_u^{-1}} \frac{I}{c\beta} \frac{1}{2\pi\Sigma^2} \frac{1}{\epsilon_0 E_0^{3/2}} \langle \Delta\eta \rangle$$

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- Producing the initial electric field by sampling according to its probability distribution function.
- Simulating the evolution of the light power inside the cavity using the single-pass gain function.

## Light power profile evolution equation

$$E_{n+1}(t) = [E_n(t)g(t) + \delta E(t)] R_{total}$$

<sup>2</sup>Li K, et al., Physical Review Accelerators and Beams, 2017, 20(3): 030702. 

# Applications: cascaded XFELO

Utilizing FEL oscillator with multi-stage undulators enables gain cascading in a single-pass, making it possible to achieve shorter single pulse lengths, higher peak power, and even higher pulse energy than normal FEL oscillator.<sup>3</sup>

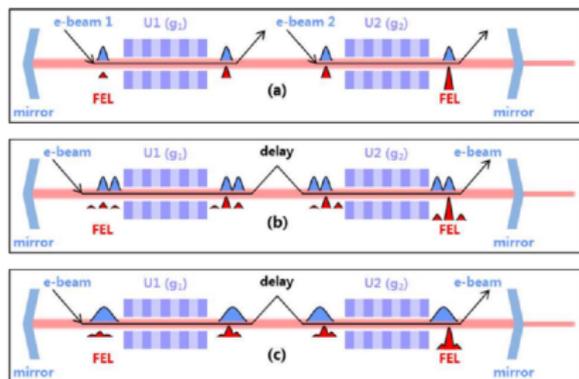


Figure: The schematic view of cascaded FELO.

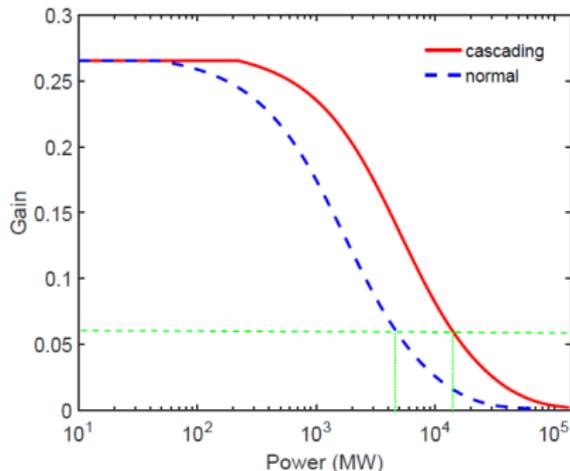


Figure: The gain of cascaded FELO calculated by new model.

<sup>3</sup>Li K, et al., Physical Review Accelerators and Beams, 2017, 20(11):110703.

# Applications: XFEL design at SCLF

The quasi-CW, FEL quality electron beams at Shanghai Coherent Light Facility (SCLF) is suitable to consider an X-ray free electron laser oscillator (XFEL) operation.

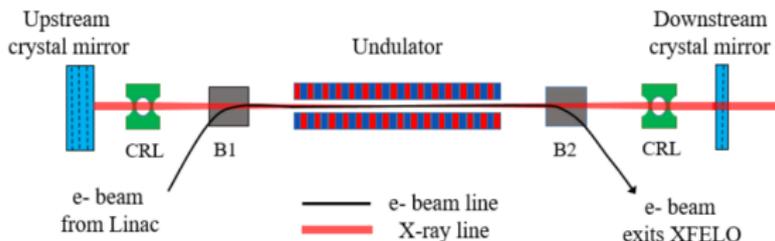


Figure: Schematic view of XFEL for SCLF.

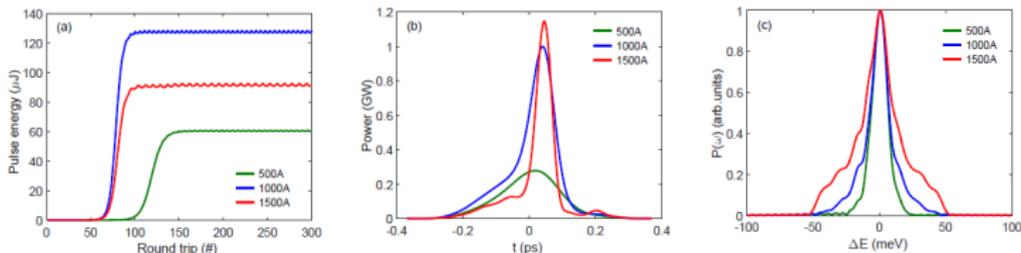


Figure: Performances of XFEL for SCLF.



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# Thanks

*Thank you!*

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