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What should Start to End Simulations for Light Source ERLs be sure to include?

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Content:

effects

- simulation codes
- start to end simulations





Storage Rings, LINACs and ERLs: variation in parameters

	Storage Rings	LINACs	ERLs
Mean path length / m	10 ¹³ (τ =10h)	10 ¹ - 10 ⁴	10 ¹ – 10 ²
Mean travel time / s	10 ⁴	<10 ⁻⁴	<10 ⁻⁵
Norm. emittance / mm mrad	20 – 50 (hor.)	1 (0.1)	
Current / mA	10 ²	<1	10 ²
Bunch length (rms) / ps	10 ¹	10 ⁻²	
Peak current / A	10 ¹	10 ³	
Bunch charge / nC	0.1 -1		
Rel. energy spread (sliced)	10 ⁻³	10 ⁻⁴	

ERLs compared to Storage Rings:

- large fraction of beam path goes through cavities
- single pass devices: no (less BBU) resonant behaviour → tracking procedures do not need to be as highly sophisticated as for Storage Rings



With these parameters most relevant effects for ERLs can be identified already:

• **Space Charge:** high charge density due to both small transverse (low emittance) and small longitudinal (bunch length) dimensions → energy spread & emittance growth

special care for the injector: no accumulation like in Storage Rings -> full bunch charge in one single gun shot <u>at lowest energies</u>

• <u>Coherent Synchrotron Radiation CSR</u>: short pulses emit powerful coherent radiation at wavelength comparable to the bunch length; interaction of bunch with CSR-field → energy spread & emittance growth

• <u>Cavity Wakes:</u> bunches passing the cavities deposite energy in monopole or higher order modes (HOM)
→ energy change, energy spread & emittance growth, in resonance even total beam loss (BBU)

• <u>Resistive and Geometric Wakes:</u> most relevant for small vacuum chambers (undulator sections) → energy spread & emittance growth

• **Ion Trapping:** the strong potential of the electron current attracts positively charged ions (like in Storage Rings); ion potential causes additional focusing, scattering on ions → energy spread & emittance growth



- Incoherent Synchrotron Radiation: only a minor effect at low energy ERLs, might lead to emittance and energy spread growth in high energy ERLs
- Intra Beam Scattering: build up times usually too long compared to beam travel time, can be relevant at extreme bunch densities

ERLs are used for various purposes: which parameters are most important for which device?

- ERL-FELs: high peak current, ultra low sliced energy spread and emittance some tolerance in projected values
- ERL-Light Sources: like FEL but no hard limits
- ERL-Cooler: cooling most efficient at high electron current, low energy spread and emittance, sensitive to projected values
- ERL-Collider: high current and large emittance

Focus on FEL- and Light Source ERLs: observe <u>sliced and projected energy spread and</u> <u>emittance</u> of compressed bunches



Next slides:

short introduction to the relevant effects that can be considered in single bunch Startto-End (S2E) simulations.

Note that the following effects are not considered:

• collective, multi bunch effects like BBU/Ion trapping: analytic formulas and simulation codes (e.g. ERLBBU, BI) to estimate current thresholds are available, but not implemented in the existing multi particle tracking codes

• FEL process: beam quality after FEL interaction must be good enough to allow for deceleration without particle losses – ask experts from WG 3 for this topic



Coherent Synchrotron Radiation CSR

- Synchrotron radiation is usually emitted incoherently
- if $\sigma_1 \le \lambda \rightarrow$ Coherent emission of radiation: $P_{CSR}(f) = N P_0(f) + N^2 P_0(f) F(f)$
 - P₀: P(single electron), F(f): form factor = FFT of longitudinal bunch charge density



• ERLs: CSR can be emitted in dipole magnets of bunch compressors, 180 degree arcs, doglegs,



Coherent Synchrotron Radiation CSR

• CSR fields from the tail of the bunch can catch up with the head







Fig.: 1D CSR wake for a Gaussian bunch at varying positions inside a dipole magnet



Coherent Synchrotron Radiation CSR



- CSR \rightarrow energy modulation over the bunch
- energy modulation \rightarrow growth of energy spread (correlated & sliced)
- energy modulation in dispersive sections \rightarrow emittance growth (correlated & sliced)



Space Charge

• electron bunch: cloud of $10^9 - 10^{10}$ charged particles with repulsive forces

$$\vec{F}(\vec{r}) = \frac{e^2}{4\pi\varepsilon_0} \sum_{l=1}^N \frac{\vec{r} - \vec{r}_l}{\left\|\vec{r} - \vec{r}_l\right\|^3} \qquad \vec{r} \neq \vec{r}_l \qquad \text{(static)}$$

Transverse Space Charge:

• simple example: round beam with radius a:





• space charge in a drift acts like a quadrupole magnet, defocusing in both planes:

elliptic beam: $\frac{1}{F_x} = \frac{4I(s)}{(\gamma\beta)^3 I_A} \frac{L}{a(a+b)} = f(s) \qquad \begin{array}{l} a^2 = \varepsilon_x \beta_x \\ b^2 = \varepsilon_y \beta_y \end{array}$ L: considered path length $b^2 = \varepsilon_y \beta_y \qquad I_A$: Alfven current growth of projected emittance: $\varepsilon_x = \varepsilon_{x_0} \sqrt{1 + \beta_{x_0}^2 \left(\left\langle \frac{1}{F_x^2} \right\rangle - \left\langle \frac{1}{F_x} \right\rangle^2\right)}$

Longitudinal Space Charge (LSC)

- repulsive forces in direction of motion: particles in the bunch head gain energy, those in the tail lose energy $\vec{E}(r, z, t) \rightarrow$ induce momentum modulation \rightarrow growth of energy spread
- at low energies: space charge oscillations

$$\mathcal{D}_{LSC}^{2}(k) = c^{2} \frac{kI}{\gamma^{3}I_{A}} \frac{4\pi \left| Z_{LSC}(k) \right|}{Z_{0}}$$

 $Z_{LSC}(k) = \frac{iZ_0}{\pi k r_b^2} \left| 1 - \frac{kr_b}{\gamma} K_1\left(\frac{kr_b}{\gamma}\right) \right| \approx \frac{iZ_0 k}{4\pi\gamma^2} \left(1 + 2\ln\frac{\gamma}{kr_b} \right)$

on axis LSC impedance: (round, uniform beam)



Fig: LSC-induced energy modulation amplitude in a drift section E=120 MeV, I=120 A, \pm 5% initial density modulation at 50/100 μ m.

 $\frac{kr_b}{<<1}$

Courtesy of Z. Huang



Microbunching Instability

• LSC in combination with CSR: microbunching instability

small charge density modulation from the injector (e.g. due to modulation of photo-cathode laser power)

 \rightarrow LSC induces momentum variation

→ momentum modulation transforms to density modulation (by LSC oscillations at low energy or in dispersive sections)

 \rightarrow density modulation is source of strong CSR fields in compressors etc.





Microbunching Gain:

Final bunching factor $b_f(k_f)$ Initial bunching factor $b_0(k_0)$

$$b(k) = \frac{1}{Nec} \int I(z) e^{-ikz} dz$$

Fig.: CSR microbunching gain for the CSR2002 workshop benchmark case (chicane @ 5 GeV) as a function of the initial density modulation wavelength.

Courtesy of Z. Huang



Microbunching Instability



Fig.: simulated LCLS longitudinal phase space at the end of DL2 with 1% initial density modulation at 15 μm without a laser heater.

Fig.: like before (figure on the left) but this time with the laser heater.

Courtesy of Z. Huang, M. Borland



Microbunching Instability



65.2 65.1 65 64.9 64.8-2 -1 0 1 2

Fig.: momentum distribution after bunch compressor measured at the DUV-FEL (hor. axis corresponds to momentum).

Courtesy of T. Shaftan

Fig.: simulated longitudinal phase space with projections of the DUV-FEL (after bunch compressor)



Cavity Wake Fields

- bunches deposit power in the cavities → cavity wake fields (beam loading both to fundamental and to HOMs)
- effects are considered in different time scales: single bunch and multi bunch effects (BBU)

Single Bunch Effects:

- wake field varies over the bunch length \rightarrow bunch energy is modulated (long. wakes) and emittance can be increased (transv. wakes)
- wake functions: simulations give empirical analytical expression, e.g. TESLA module (8 cavities)





Resistive & Geometric Wakes

• **<u>Resistive Wall Wakes</u>**: longitudinal wake for a round pipe of radius a, (dc) conductivity σ

$$W(s) = \frac{4Z_0c}{\pi a^2} \left(\frac{e^{-s/s_0}}{3} \cos \frac{\sqrt{3}s}{s_0} - \frac{\sqrt{2}}{\pi} \int_0^\infty \frac{dx \, x^2 e^{-x^2 s/s_0}}{x^6 + 8} \right) \qquad s_0 = \left(\frac{2a^2}{Z_0\sigma}\right)^{\frac{1}{3}}$$

• analog expression for the transverse resistive wake





Resistive & Geometric Wakes

• Surface Roughness Wakes: small cavity-like structures



$$W(s) = \frac{Z_0 c h^2 k_1^3}{4\pi a} f(k_1 s)$$

$$k_1 = 2\pi / p$$
 $f(x) = \frac{1}{2\sqrt{\pi}} \frac{\partial}{\partial x} \frac{\cos(x/2) + \sin(x/2)}{\sqrt{x}}$

$$W(s) = -\frac{Z_0 ch^2}{16\pi^{3/2} a} \left(\frac{k_1}{s}\right)^{3/2} \qquad \text{for} \quad \sigma_l \ll k_1$$

Fig. (left): Geometry for the "small corrugations model" of roughness impedance.

Fig.(bottom): wake of shallow, periodic corrugations for a Gaussian bunch: h << p << a, $\sigma_1 < k_1$







Simulation Codes: Space Charge

- C. Limborg et al: "Code Comparison for Simulations of Photo-Injectors", PAC03:
- various codes with various approaches (even more than compared here: GPT, BEAMPATH, ITACA, ...)

code	Algorithm	?D	SC calculated from	comput. time
HOMDYN	solving envelope equation (semi-analytical)	2D	one uniformly charged cylinder	short
PARMELA	Lorentz transformed space	2D	uniformly charged rings	long
ASTRA	charge forces			medium
PARMELA SPCH3D	fast Particle-in-Cell	3D	macro particles	long
TREDI	retarded potentials, Lienard- Wiechert	3D	macro particles	long

- Simulation results: good agreement inside gun some disagreement in drift (HOMDYN)
- Thermal emittance: initial distribution can be defined in all codes, no code includes physics behind
- Shottky effect: only considered by ASTRA
- Longitudinal profile: all codes but HOMDYN
- Transverse non-uniformity: 3D codes, ASTRA upgraded



Simulation Codes: CSR

- again, various codes with various approaches:
 - 1D: ELEGANT, CSR-Calc, CSR-Track (1D option): CSR effects calculated from long. charge density, fast
 - 3D: CSR-Track / TraFiC⁴, TREDI, slow

ELEGANT: • fully equipped 3D optics simulation tool with macro-particle tracking

- special dipole and drift elements added to include CSR effects in 1D appraoch
- transverse and longitudinal wake fields can be applied
- LSC elements included

CSR-Track / TraFiC4: - 1D option like in ELEGANT

- Green's function method for fast 3D simulations

- "direct method" for many particles combining Green's function method with meshed field method

TREDI: • Lienard-Wiechert potentials approach (also for space charge effects)

ICFA CSR workshop, DESY-Zeuthen, 2002 : reasonable agreement between codes (exception: TREDI)



Simulation Codes: Combine Codes for Start-to-End Simulations

How to combine codes? How can output(code1) be transformed to input(code2)?

• often parametrization of bunch is done. Potentially neglected: a) correlations between phase space dimensions e.g. momentum chirp b) parameter variations over the bunch length (sliced values)

• better: 6D macro particle phase space is easy to share between codes. Sometimes a problem: adjusting number of particles

→ select codes for Start-to-End simulations also with respect to their abilities to share the bunch data between the programs

• In the last years: S2E simulations for LCLS, TTF, X-FEL, BESSY-FEL, ...

the following codes have often been used:

PARMELA/ASTRA \rightarrow for space charge dominated machine parts (gun, low energy sections)

ELEGANT/CSR-Track/CSR-calc \rightarrow dispersive sections (bunch compressors, arcs, doglegs)

GENESIS/FAST/GINGER/MEDUSA → FEL process

- first attempts for master programs controlling all involved codes: GlueTrack (Zagorodnov, DESY) combining ASTRA, CSR-Track and GENESIS
- individual solutions: ASTRA → ELEGANT: script using tools from SDDSToolkit

ELEGANT → GENESIS: SDDSToolkit program "elegant2genesis"



Summary

- Quality of the electron bunches has to be conserved during beam transport, acceleration, compression and all other manipulations
- Most relevant (single bunch) effects with the potential to deteriorate bunch quality: space charge and CSR fields, cavity wakes, resistive wall and geometric wakes
- Simulation codes for each of these effects are available but no code that includes all
- For Start-to-End simulations output-input conversion has to be done very carefully: 6D macro particle phase space best suited. With parametrizations: correlations in phase space have to be included