Compact low emittance lattices with longitudinal gradient bends ?

T.M.E. cell example: 12° dipole, 2.4 GeV $\rightarrow \epsilon = 1.67$ nm, $\mu_x = 284^\circ$



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Longitudinal variation of radiation integrands inside TME dipole center edge I2, 10000*I5 0.25 $I5 = \int \mathcal{H} h^3 ds$ 0,20 $I2 = \int h^2 ds$ 0.15 $\Phi = \int h ds$ 0.10 $h = B / (B\rho)$ 0.05 0.00 0,2 0.4 0.6 0.0 $\varepsilon J_x / \gamma^2 = C_q I_5 / I_2 \rightarrow min$ Problem: while Φ = const.

 \rightarrow Lower field towards edges....

Previous work

J.Guo & T.Raubenheimer LOW EMITTANCE E-/E+ STORAGE RING DESIGN USING BENDING MAGNETS WITH LONGITUDINAL GRADIENT, EPAC-2002 Factor 2..3 emittance reduction possible, application to NLC damping rings.

A.Wrulich & R.Nagaoka (1992) EMITTANCE MINIMIZATION WITH LONGITUDINAL DIPOLE FIELD VARIATION, Joint PSI-Synchrotron SOLEIL Accelerator Physics Report, 2004 ATTEMPT TO MINIMIZE THE EMITTANCE WITH LONGITUDINAL DIPOLE FIELD VARIATION, Small emittance light source workshop, Lund, 2004 Analytical and numerical treatment of exponential and polynomial field decays Zero emittance for infinite B-field in magnet center Factor 4 emittance reduction with ~ feasible magnets in DBA lattice (SOLEIL)

A.Ropert et al. PROSPECTS FOR A LONG TERM LATTICE UPGRADE AT THE ESRF, EPAC 2004 Application to existing ESRF geometry with 1.8T magnets: 0.5 nm minimum emittance



2-field dipole: $\mu = 0$ pure normal dipole $\mu = \mu_{max}$ pure s.c. dipole $\Rightarrow \mu \text{ for } \epsilon \rightarrow min ?$

Analytical solution for 2-field dipole

 $\partial^{2}\epsilon \left(\beta_{xo},\eta_{o},\mu\right) / \partial\beta_{xo} \partial\eta_{o} \rightarrow min \Rightarrow \epsilon_{min} \left(\mu\right)$



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sub-T.M.E. Cell: $\epsilon = 0.43$ nm (~ $\frac{1}{4}$ T.M.E.) at 2.4 GeV based on 2-field 12° dipole: 2.7°(0.6T) + 6.6° (7.2 T) + 2.7°(0.6T)



Trial: combined function bend

Longitudinal profile $B(s) = (B_{max} - B_{off}) / (1 + (\alpha s)^2)^{\frac{1}{2}} + B_{off}$ Numerical minimization of $\epsilon (\beta_{xo}, \eta_o, \alpha)$ for given L, Φ



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Longitudinal variation of radiation integrands

Optimum longitudinal profile and optics Result of ϵ ($\beta_{xo}, \eta_o, \alpha$) \rightarrow min

Constant field bend in TME mode



sub-T.M.E. Cell: $\epsilon = 0.46$ nm ($J_x = 1.3$) at 2.4 GeV based on 12° combined function bend (max. 1.6 T)



D.A. ~ \checkmark (ideal cell) $\mu_x = 311^\circ \rightarrow \text{long cell 8.4 m}$





10 x 36° arc, 362 m circumference, $\varepsilon = 0.47$ nm at 2.4 GeV straights: 10 x 6 m dispersions free, 30 x 3 m dispersive

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Relaxed cell: $\beta_{xo} \mathbf{f} \ \eta_o \mathbf{f} \Rightarrow$ less variation of $\mathcal{H}(s)$ \Rightarrow *less gain from longitudinal gradient*

Alternative:give up on long.grad.bend (\odot) $\epsilon \rightarrow \sim \epsilon \ge 4$ relax the cell $\epsilon \rightarrow \sim \epsilon \ge 3$ split bend (use $\mu_x > 180^\circ$ space !) $\epsilon \rightarrow \sim \epsilon / 8$ $\epsilon \rightarrow \sim 1.5 \ge \epsilon$





Conclusion

- Feasible bending magnets with longitudinal gradient provide an emittance which is about ¼ of the T.M.E.
- This works well only for ideal cells with exact matching, having serious disadvantages:
 - horiz. phase advance > 180° requires a 2^{nd} focus and a long cell
 - sensitive to errors and little tunability
 - large chromaticities and little flexibility on phases restrict D.A.
- In relaxed cells, the particular emittance reduction from the longitudinal gradient is partially lost.
- Longitudinal gradient bends offer a promising way to build compact low emittance lattices, but more studies are required to achieve at realistic designs.