

Coherent Spontaneous-Superradiance and Stimulated-Superradiance of Bunched Electron Beams

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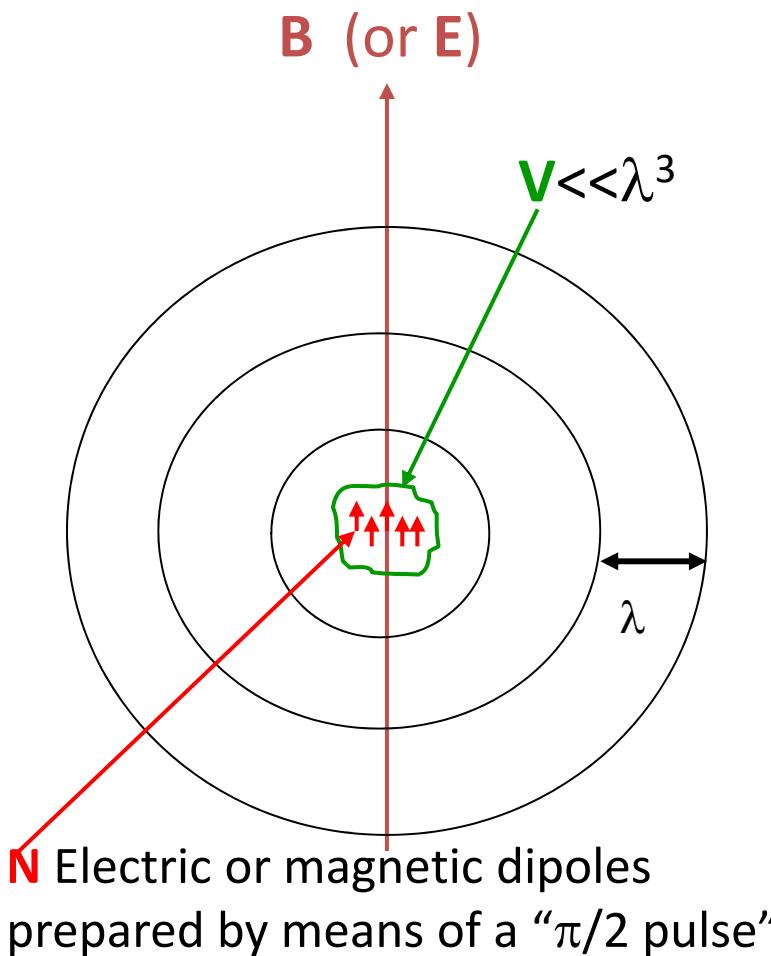
Rev. Mod. Phys. 91, 035003 – Published 19 August 2019

Tutorial - FEL International conference Hamburg, August 2019

CONCEPTS TO BE CONSIDERED

- Fundamental processes of coherent radiation emission from bunched beam current:
 - Spontaneous superradiant emission (SP-SR)
 - Stimulated Superradiant emission (ST-SR).
- Nonlinear SR, ST-SR dynamics of a trapped bunched beam.
- Tapering Enhanced Superradiant (TES) and Stimulated Superradiance (TESSA)
- Bunched-beam/radiation Self-interaction.
- Review of applications.

DICKE'S SUPERRADIANCE



$$I = (r+m)(r-m+1)I_0$$

$$r = 1/2, 1, 3/2, \dots, N/2$$

$$m = -r, \dots, 0, \dots, r$$

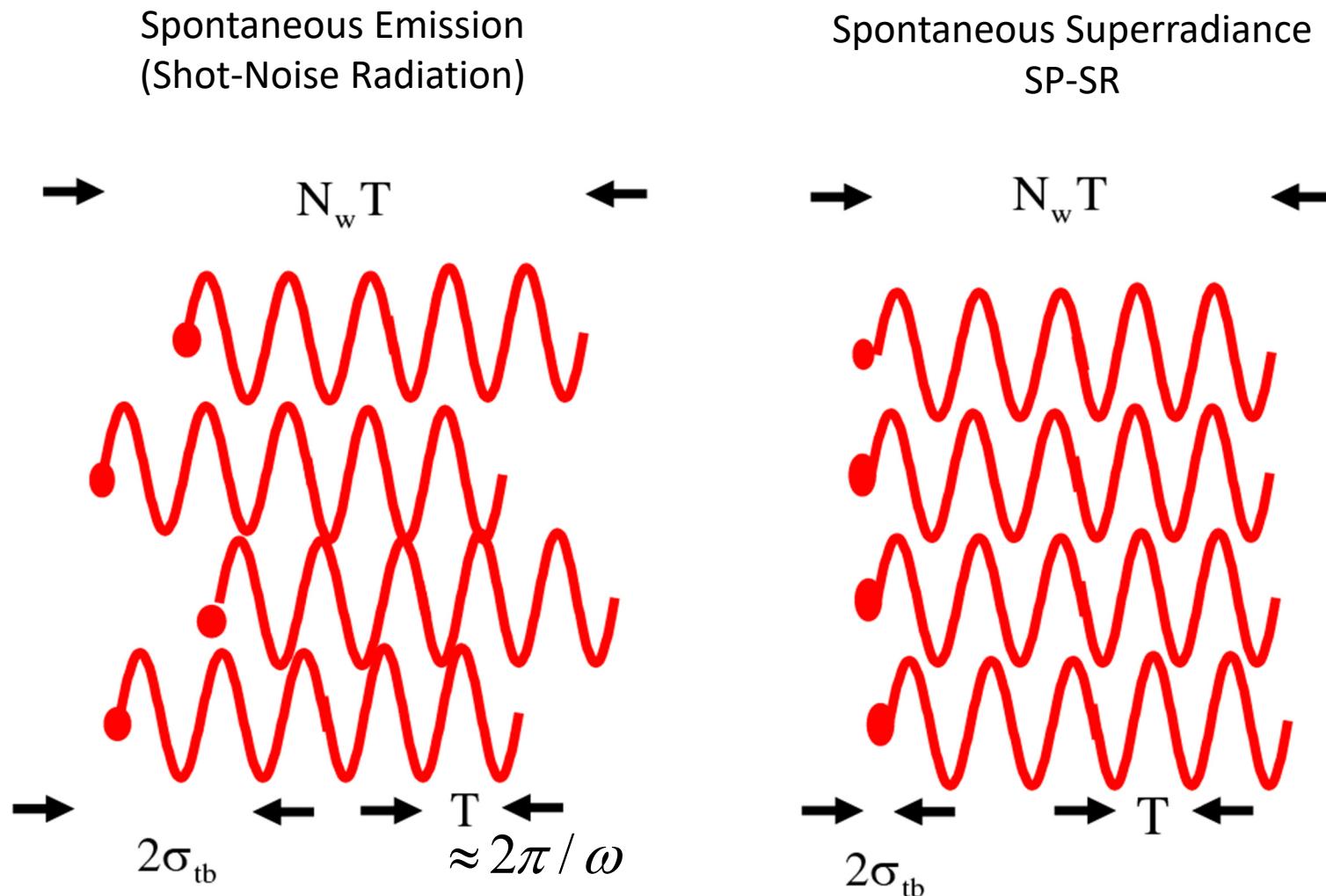
Spontaneous emission (quantum)

$$N=1, \quad r = m = 1/2 \quad \Rightarrow \quad I = I_0$$

Superradiant emission (classical)

$$r = N/2 \gg 1, \quad m = 0 \quad \Rightarrow \quad I = \frac{1}{4} N^2 I_0$$

COHERENT vs RANDOM SUPERPOSITION OF RADIATION WAVEPACKETS



Formulation of Radiation mode Expansion

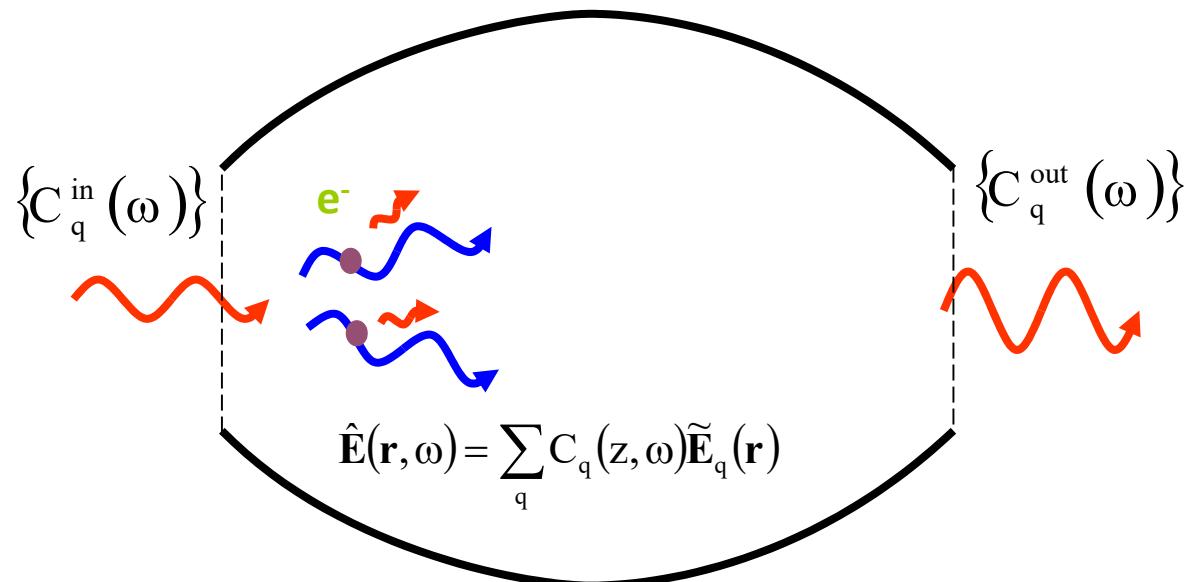
Spectral (Fourier)
Frequency domain

$$\check{\mathbf{E}}(\mathbf{r}, \omega) = \sum_{\pm q} \check{C}_q(z, \omega) \tilde{\mathbf{E}}_q(\mathbf{r})$$

$$\check{\mathbf{H}}(\mathbf{r}, \omega) = \sum_{\pm q} \check{C}_q(z, \omega) \tilde{\mathbf{H}}_q(\mathbf{r})$$

$$\frac{d\check{C}_q(z, \omega)}{dz} = -\frac{1}{4\mathcal{P}_q} \int \check{\mathbf{J}}(\mathbf{r}, \omega) \cdot \tilde{\mathbf{E}}_q^*(\mathbf{r}) dA$$

$$\check{C}_q^{out}(\omega) - \check{C}_q^{in}(\omega) = -\frac{1}{4\mathcal{P}_q} \int \check{\mathbf{J}}(\mathbf{r}, \omega) \cdot \tilde{\mathbf{E}}_q^*(\mathbf{r}) dV$$



WAVEPACKETS EMISSION BY PARTICULATE CHARGES

For particulate current: $\mathbf{J}(\mathbf{r}, t) = \sum_{j=1}^N -e\mathbf{v}_j(t)\delta(\mathbf{r} - \mathbf{r}_j(t)) \quad \Rightarrow$

Radiation wavepackets $C_q^{out}(\omega) = C_q^{in}(\omega) - \frac{1}{4\mathcal{P}_q} \sum_{j=1}^N \Delta\tilde{\mathcal{W}}_{qj}$

$$\Delta\tilde{\mathcal{W}}_{qj} = -e \int_{-\infty}^{\infty} \tilde{\mathbf{v}}_j \cdot \tilde{\mathbf{E}}_q^*(\mathbf{r}_j(t)) e^{i\omega t} dt$$

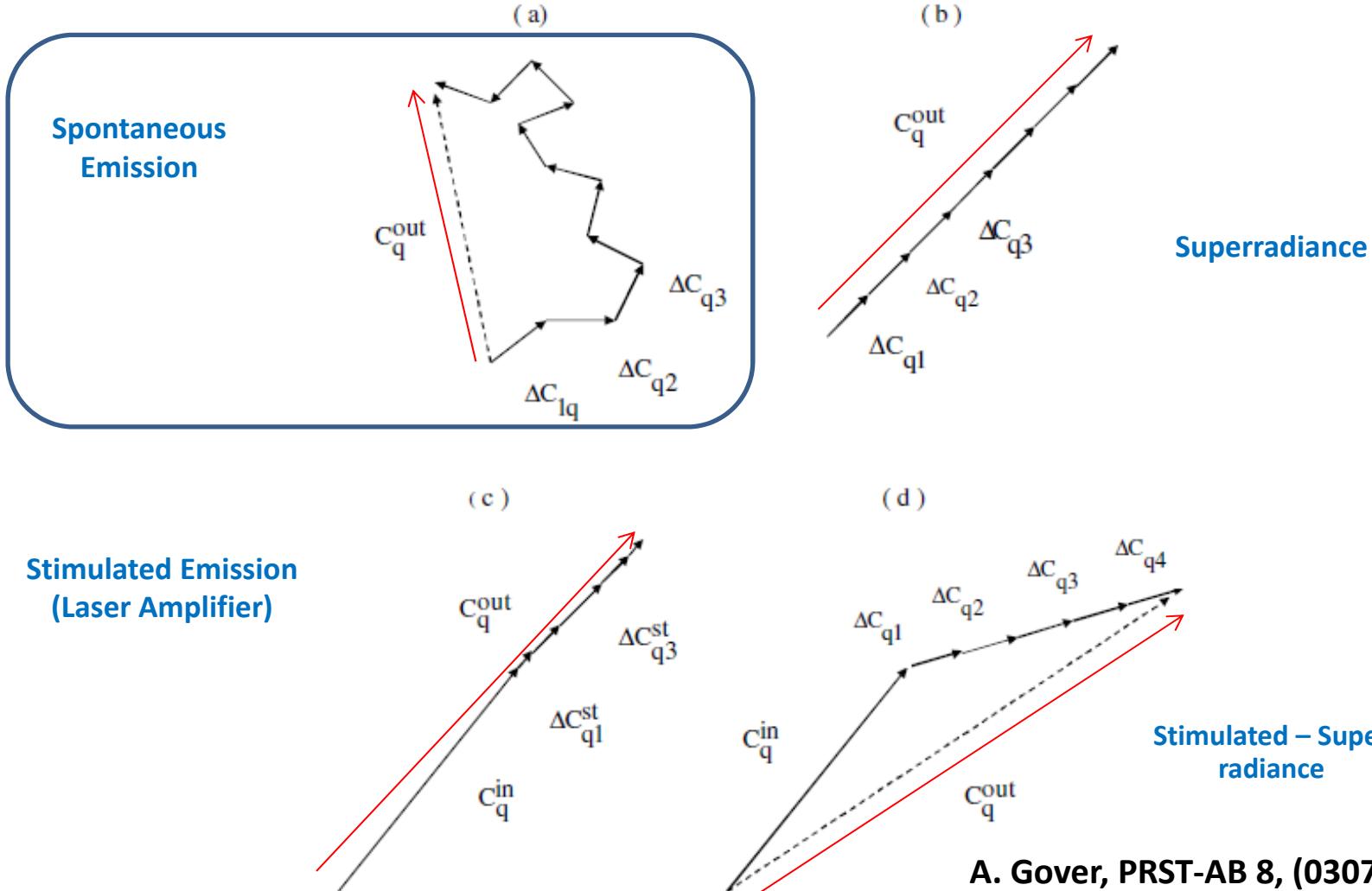
Zero order: $\Delta\tilde{\mathcal{W}}_{qj}^{(0)} = \Delta\tilde{\mathcal{W}}_{qe}^{(0)} e^{i\omega t_{j0}}$

Spectral radiant energy:

$$\begin{aligned} \frac{dW_q}{d\omega} &= \frac{2}{\pi} P_q |C_q^{out}(\omega)|^2 = \\ &= \left(\frac{dW_q}{d\omega}\right)_{in} + \left(\frac{dW_q}{d\omega}\right)_{SP/SR} + \left(\frac{dW_q}{d\omega}\right)_{ST-SR} \end{aligned}$$

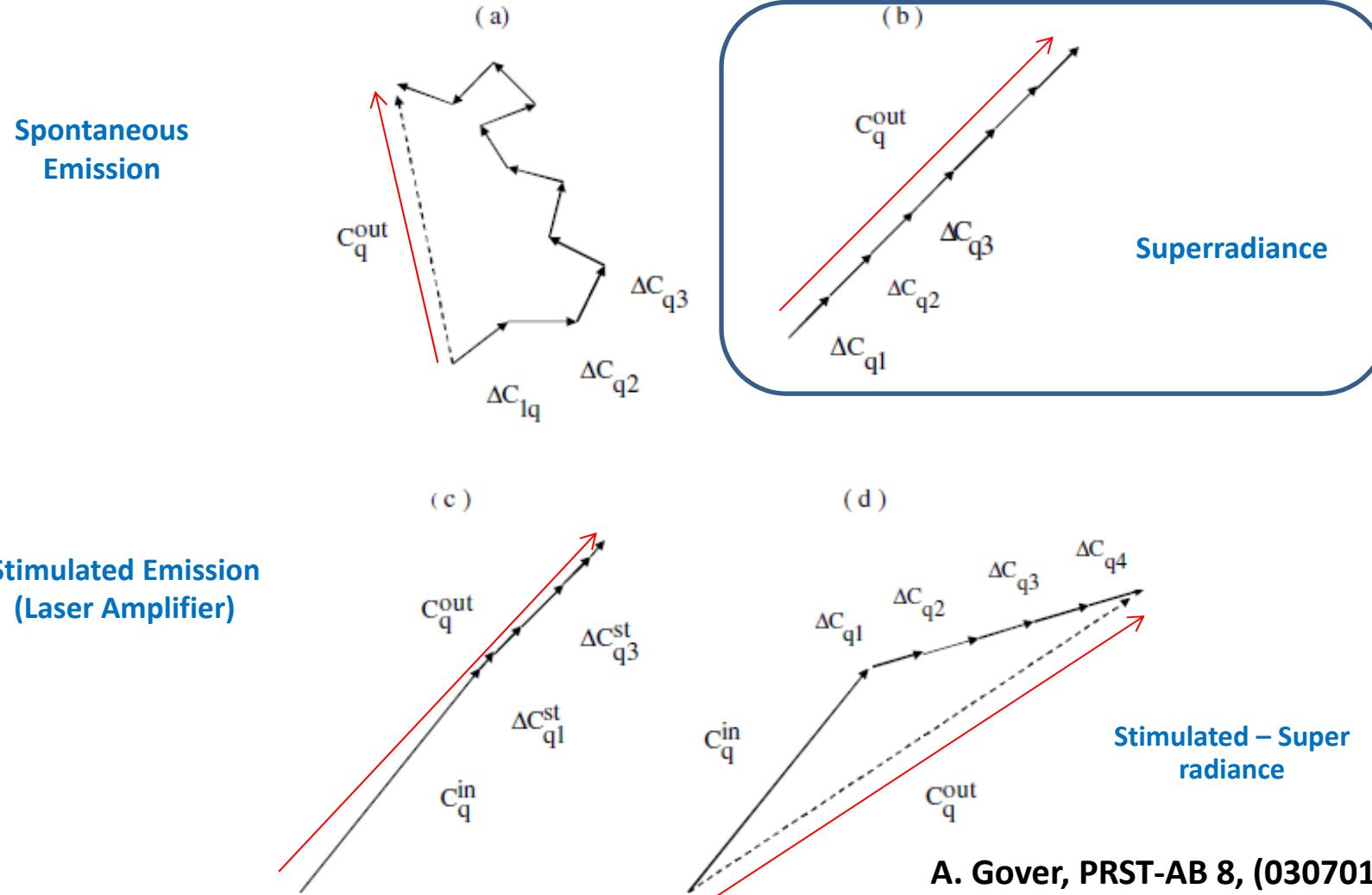
COMPLEX PLANE REPRESENTATION OF SINGLE MODE EXCITATION AMPLITUDES

$$\check{C}_q^{out} = \check{C}_q^{in} + \sum_j \Delta \check{C}_{qj}$$



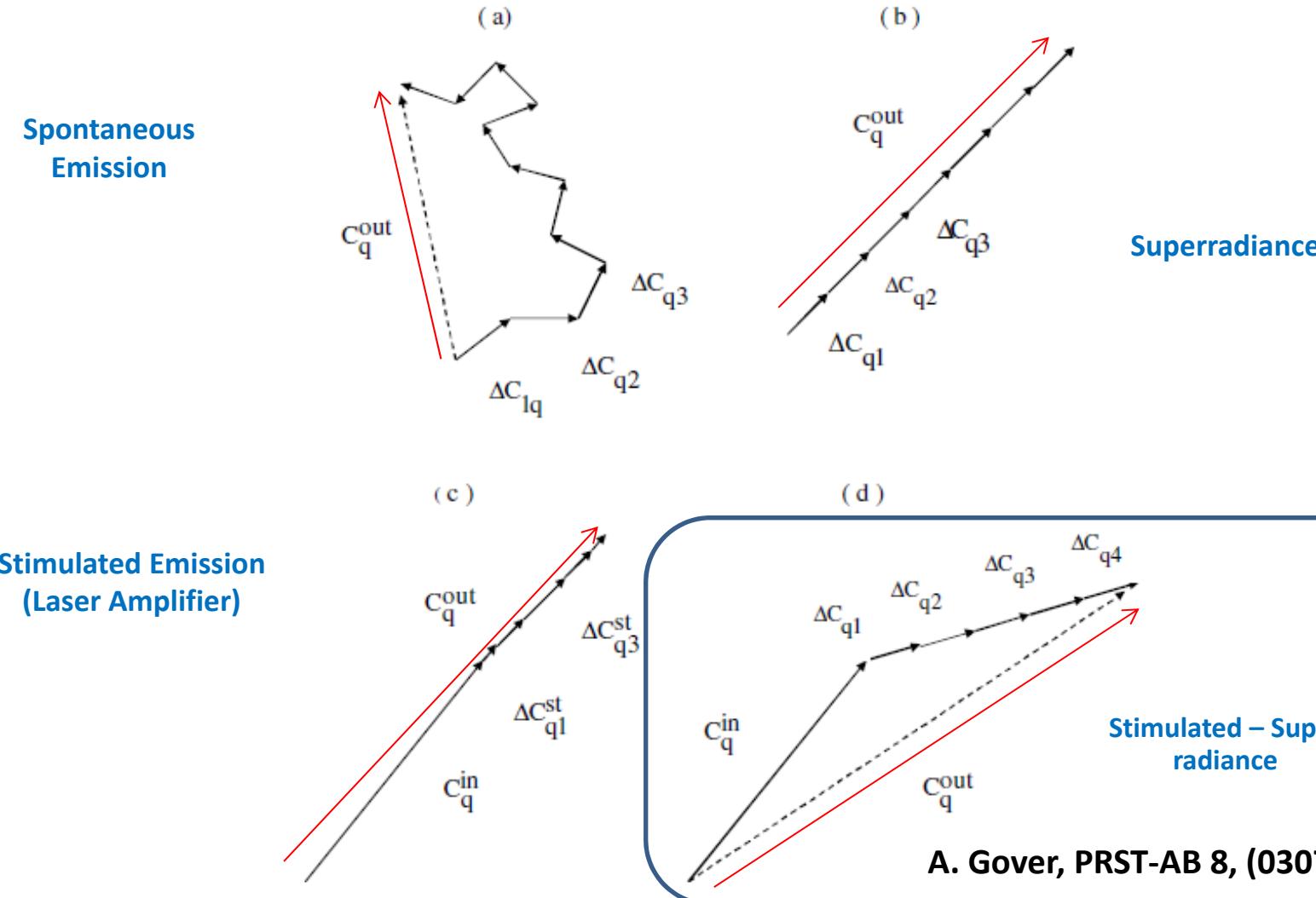
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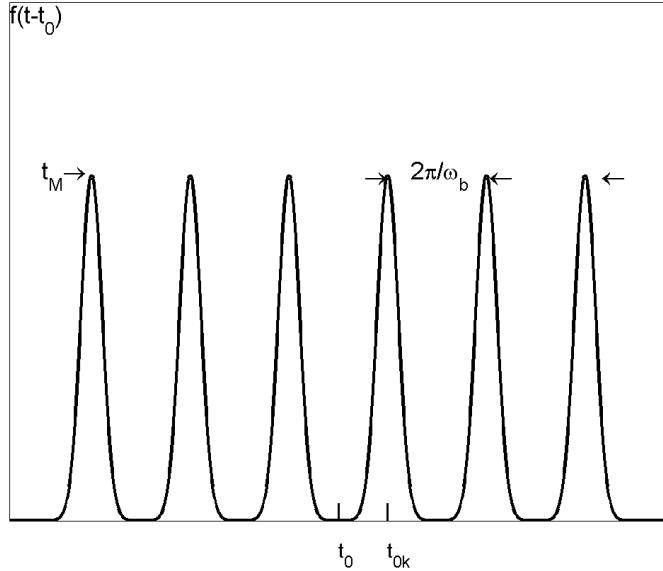


COMPLEX PLANE REPRESENTATION OF SINGLE MODE EXCITATION AMPLITUDES

$$\check{C}_q^{out} = \check{C}_q^{in} + \sum_j \Delta \check{C}_{qj}$$



A TRAIN OF PULSES (MACRO-PULSE)



$$\sum_{j=1}^N e^{i\omega t_{0j}} = \sum_{k=1}^{N_M} \sum_{j=1}^{N_{bk}} e^{i\omega t_{0j}} =$$

$$= N \cdot M_b(\omega) \cdot M_M(\omega) \cdot e^{i\omega t_0}$$

Bunch form factor:

$$M_b(\omega) = \frac{1}{N_M} \sum_{j=1}^{N_M} e^{i\omega t_{0j}} = \int_{-\pi/\omega_b}^{\pi/\omega_b} f(t'_0) e^{i\omega t'_0} dt'_0 \approx \mathcal{F}\{f(t'_0)\}$$

Macro-pulse form factor:

$$M_M(\omega) = \frac{1}{N_M} \sum_{k=1}^{N_M} e^{i\omega t_{0k}} = \frac{\sin(N_M \pi \omega / \omega_b)}{N_M \sin(\pi \omega / \omega_b)}$$

SR and ST-SR of UNDULATOR RADIATION

For a finite train of periodic bunches:

$$\left(\frac{dW_q}{d\omega} \right)_{SP-SR} = \frac{N^2 e^2 Z_q}{16\pi} \left(\frac{\bar{a}_w}{\beta_z \gamma} \right)^2 \frac{L^2}{A_{em}} |M_b(\omega)|^2 |M_M(\omega)|^2 \text{sinc}^2(\theta L/2)$$

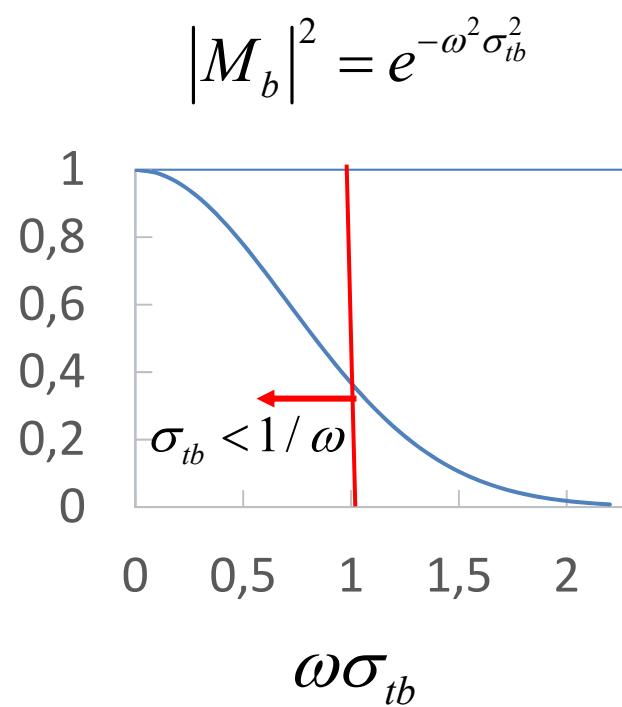
$$\left(\frac{dW_q}{d\omega} \right)_{ST-SR} = |\check{C}_q^{in}(\omega)| \frac{Ne}{2\pi} \left(\frac{\bar{a}_w}{\beta_z \gamma} \right) \sqrt{\frac{2Z_q \mathcal{P}_q}{A_{emq}}} L |M_b(\omega)| |M_M(\omega)| \text{sinc}(\theta L/2) \cos(\varphi - \theta L/2)$$

Detuning parameter: $\theta(\omega)L = \left(\frac{\omega}{v_z} - k_{zq}(\omega) - k_w \right) L \cong 2\pi \frac{\omega - \omega_0}{\Delta\omega}$

$$\Delta\omega = \omega_0/N_w, \quad \omega_0 = 2\gamma_{z0}^2 \lambda_w \quad \gamma_z = \gamma / \sqrt{1 + \bar{a}_w^2}, \quad \bar{a}_w = \bar{B}_w / k_w mc$$

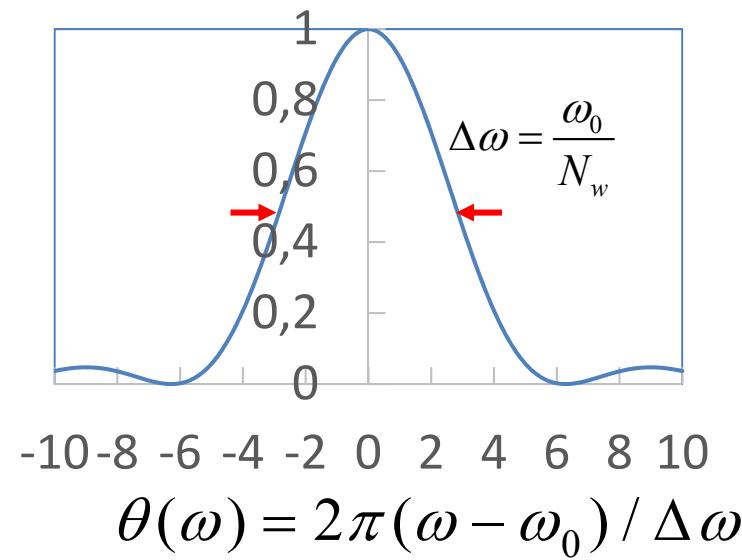
Bunching coefficient:

(for $f(t_0) = e^{-t_0^2/2\sigma_{tb}^2} / \sqrt{2\pi}\sigma_{tb}$)



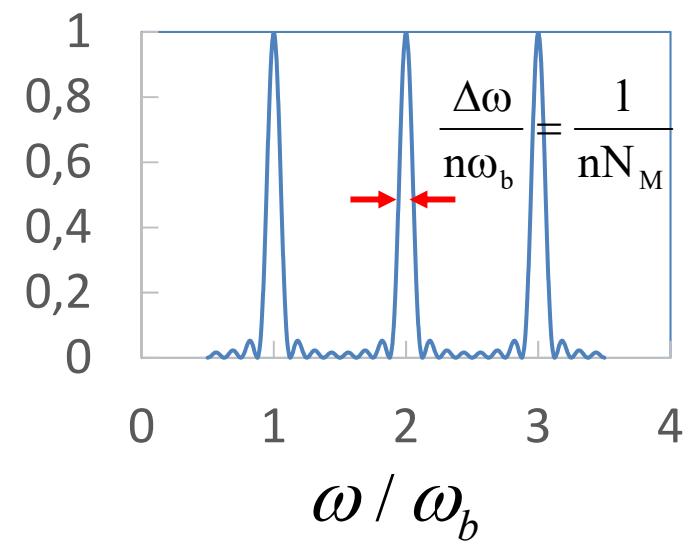
Finite interaction length frequency line-shape (N_w =# of wiggler periods)

$$\sin c^2(\theta(\omega)L/2)$$

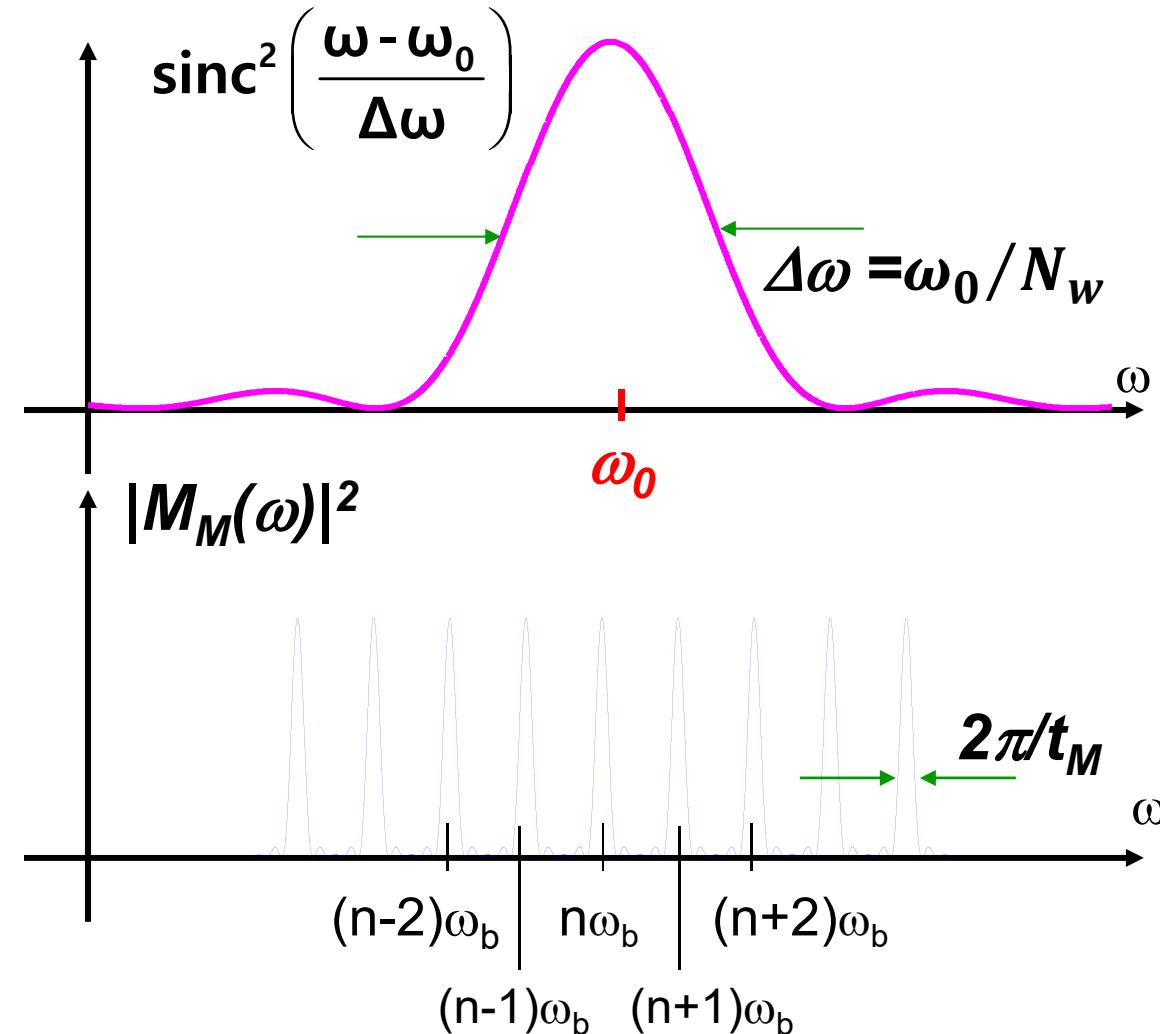


Pulse train form-factor: (N_M =# of bunches in macropulse)

$$|M_M(\omega)|^2 = \frac{\sin^2(N_M\pi\omega/\omega_b)}{N_M^2 \sin^2(\pi\omega/\omega_b)}$$



Superradiant Emission from a Pulse Composed of a Train of Bunches

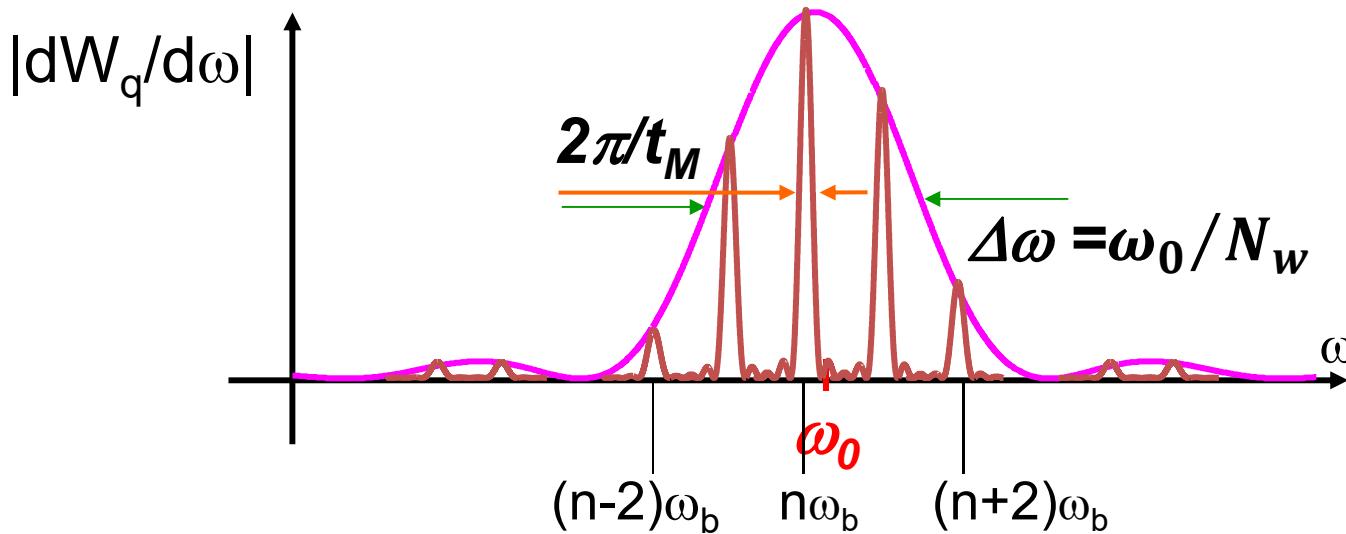


Single harmonic:

$$\omega_0 \approx n\omega_b$$

$$\begin{aligned} \omega_0 / N_w &= n\omega_b / N_w < \omega_b \Rightarrow \\ n &< N_w \end{aligned}$$

Superradiant Emission from a Pulse Composed of a Train of Bunches

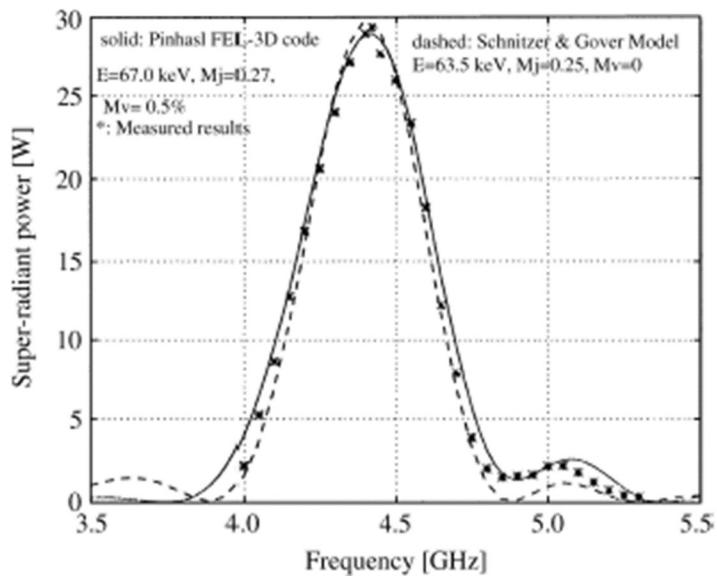


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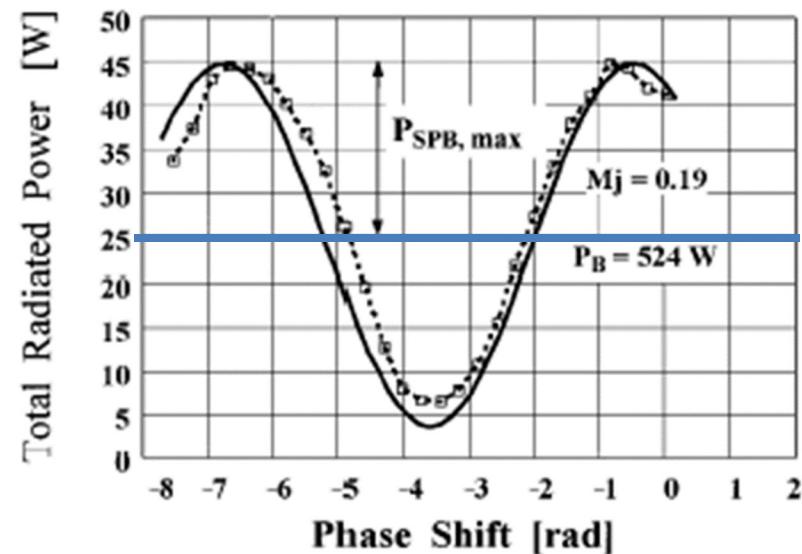
$$\omega_0 / N_w = n\omega_b / N_w < \omega_b \Rightarrow \\ n < N_w$$

PREBUNCHED – FEM SUPERRADIANCE MEASUREMENT



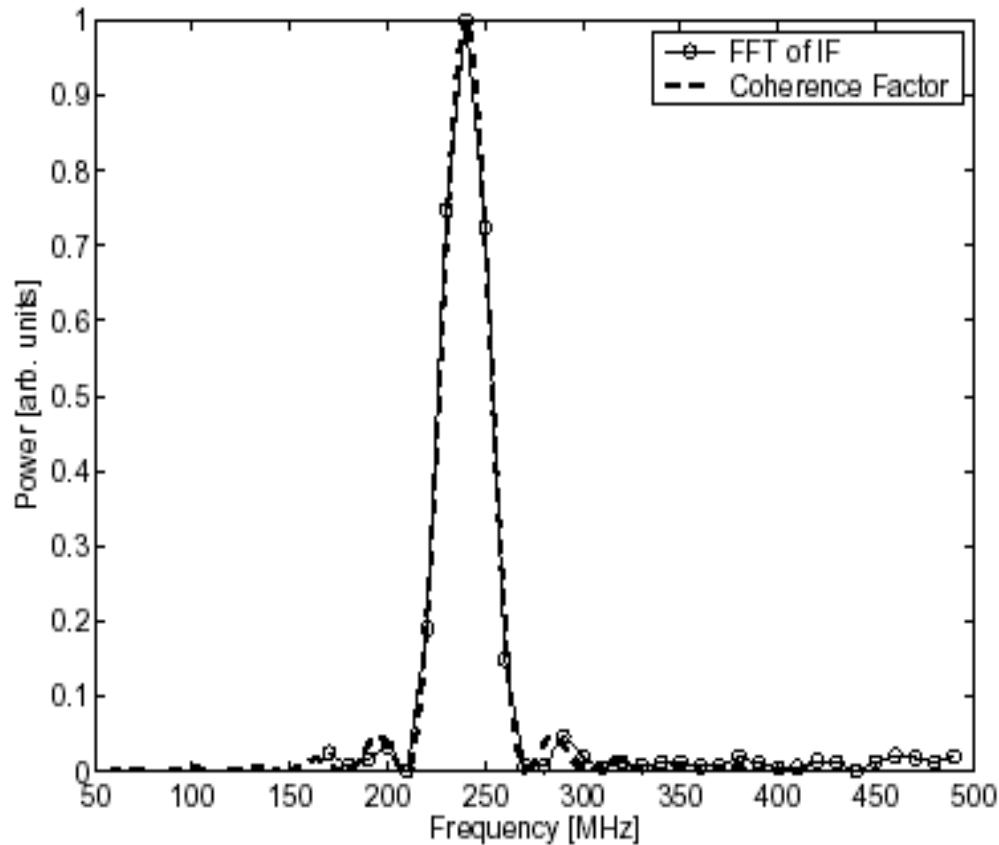
A. Cohen et al, PRL, **74**, 3812 (1995) ;
M. Arbel et al, NIM **A445** (2000)

PREBUNCHED-FEM STIMULATED SUPERRADIANCE MEASUREMENT



M. Arbel et al, PRL, **86**, 256 (2001)

Mesured multi-bunch coherent Smith-Purcell linewidth (MIT - S.E. Korbley et al PRL 2005)



$$f_n = n f_b = 240 \text{ GHz}$$

14 1.7 GHz

$$\Delta f / f_n = 1 / n N_M = 1.3 \cdot 10^{-4}$$

14 550

SP-SR HARMONIC POWER EMISSION of A LONG PERIODICALLY BUNCHED e-BEAM ($N_M \gg N_w$)

Discrete harmonics: $\omega_n = n\omega_b = 2\gamma_{z0}^2 \lambda_w$

Spectral radiant energy:

$$\left(\frac{dW_q}{d\omega} \right)_{SR} = \frac{N^2 e^2 Z_q}{16\pi} \left(\frac{\bar{a}_w}{\beta_z \gamma} \right)^2 \frac{L^2}{A_{em}} |M_b(\omega)|^2 |M_M(\omega)|^2 \text{sinc}^2(\theta L/2)$$

Integrate over frequencies \rightarrow Power of harmonic n:

$$P_{q,n} = \frac{dW_{qn}(\omega = \omega_n)}{d\omega} \left(\frac{\omega_b}{N_M} \right) / t_M$$

$$P_{q,n} = \frac{I_0^2}{8} \sqrt{\frac{\mu_0}{\epsilon_0}} \left(\frac{\bar{a}_w}{\beta_z \gamma} \right)^2 |b_n|^2 \frac{L_w^2}{A_{em,q}(\omega_n)} \text{sinc}^2(\theta(\omega_n)/2)$$

[Bunching parameter :

$$b_n = M_b(\omega = \omega_n) = e^{-\omega_n^2 \sigma_{tb}^2 / 2}$$

ELECTRON BEAM BUNCHING SCHEMES

- Photo-cathode emission (sub-pSec)
- Klystron (RF cavity) (Superradiance)
- Optical klystron oscillator (Stimulated-Superradiance)(Vinokurov and Skrinsky in 1977)
- Ultrafast laser bunching:

HGHG (Li Hua Yu)

EEHG (Stupakov)

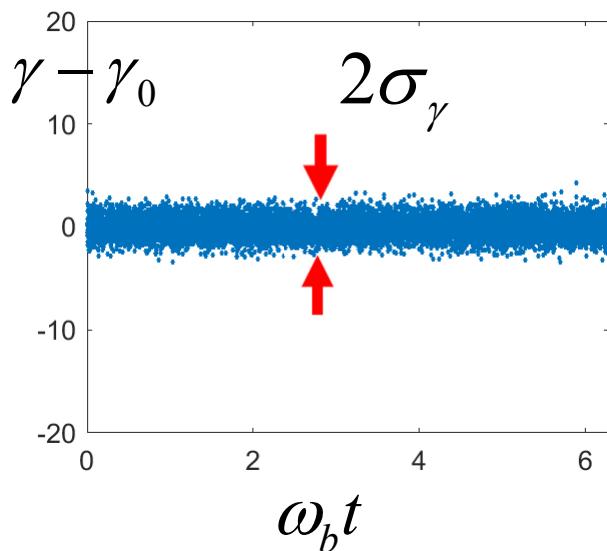
E-SASE (Zholent)

Extensive reference list: Gover et al Rev. of Mod. Phys. Vol/EID: 91/035003 (August 2019)

BUNCHING BY LASER MODULATION AND DISPERSIVE SECTION

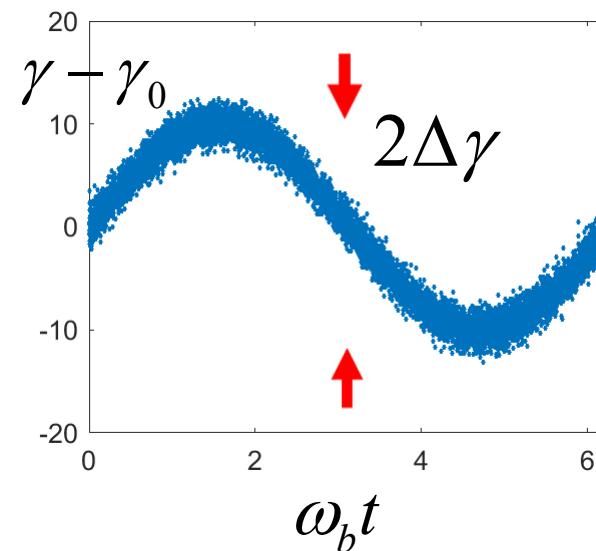
Energy distribution:

$$f(\gamma_j - \gamma_0) = \frac{1}{\sqrt{2\pi}\sigma_{\gamma_0}} e^{-(\gamma_j - \gamma_0)^2 / 2\sigma_{\gamma_0}^2}$$



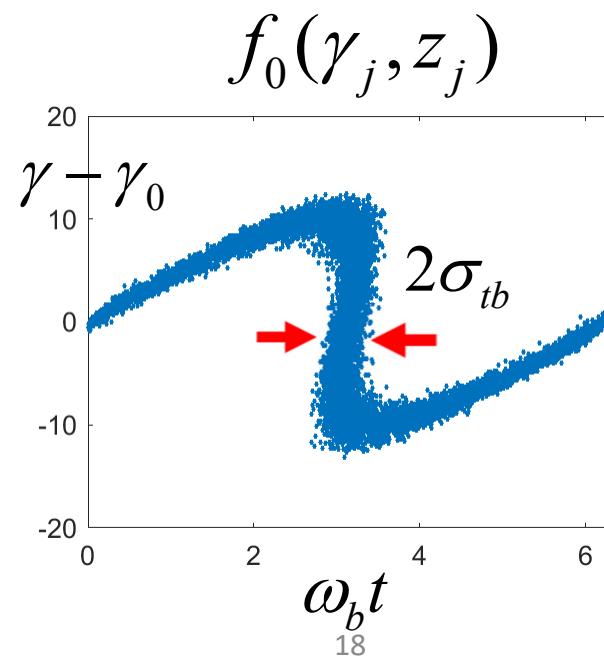
Energy modulation:

$$\gamma_i = \gamma_0 + \Delta\gamma_{\text{mod}} \sin \omega_b t_j$$



Dispersive section:

$$t_j' = t_j + \frac{R_{56}}{c} (\gamma_j - \gamma_0)$$



BUNCHING OPTIMIZATION

Normalized parameters *:

$$p_j = \frac{\gamma_j - \gamma_0}{\sigma_{\gamma 0}}, \quad A = \frac{\Delta \gamma_{\text{mod}}}{\sigma_{\gamma 0}}, \quad B = \omega_b \frac{R_{56}}{c} \frac{\sigma_{\gamma 0}}{\gamma_0}$$

→ $f_0(p_j, t_j) = \frac{1}{\sqrt{2\pi}} e^{-[p_j - A \sin \omega_b t_j - B p_j]^2 / 2}$

$$b_n = \left\langle e^{i \omega_n \Delta t_j'(t_j, \gamma_j)} \right\rangle = \int_{-T_b/2}^{T_b/2} dt_j \int_{-\infty}^{\infty} dp_j f_0(t_j, p_j)$$

$$\boxed{b_n = J_n(nAB) e^{-\omega_n^2 \sigma_{tb}^2 / 2}} \quad \left(\sigma_{tb} = \frac{R_{56}}{c} \frac{\sigma_{\gamma 0}}{\gamma_0} \right)$$

Maximal bunching for $A B \approx 1$:

$$b_n \approx \frac{0.67}{n^{1/3}} e^{-\omega_n^2 \sigma_{tb}^2 / 2}$$

*Hemsing E., Stupakov G., Xiang D., & Zholents A., Rev. Mod. Phys., 86(3), 897 (2014)

Phase-merging enhanced harmonic generation

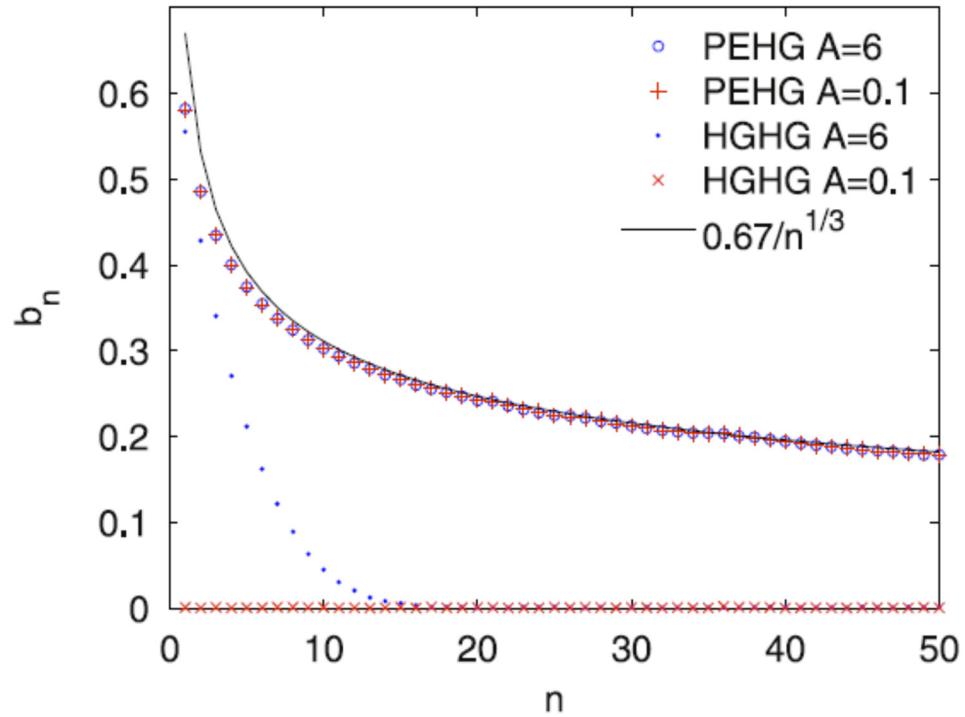
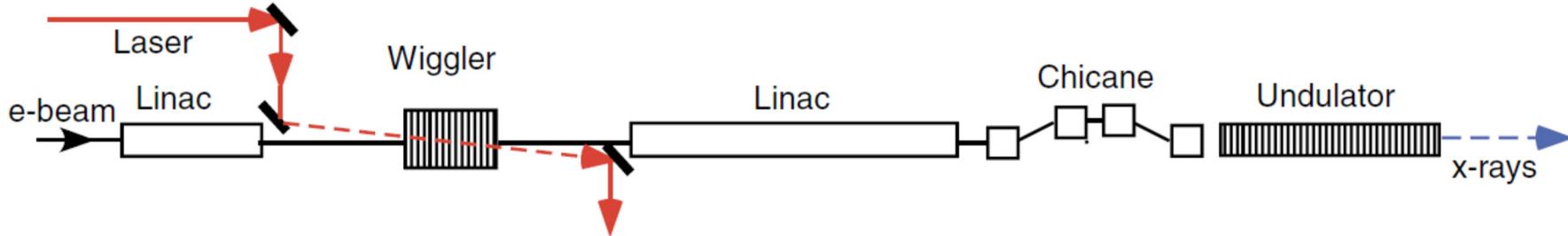


Figure 3. Comparison of the bunching factor of PEHG and standard HGHG with different energy modulation amplitudes. The black line is the theoretical prediction of the maximal bunching factor of PEHG.

Feng, C., Deng, H., Wang, D., & Zhao, Z. (2014). Phase-merging enhanced harmonic generation free-electron laser. *New Journal of Physics*, 16(4), 043021.

ULTRA-COMPACT X-RAY FEL BASED ON e-SASE*

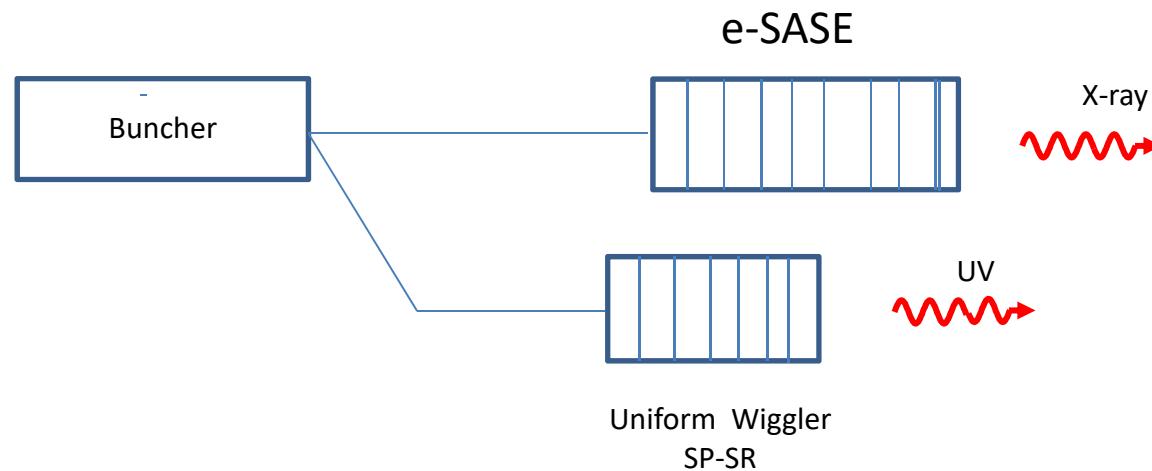
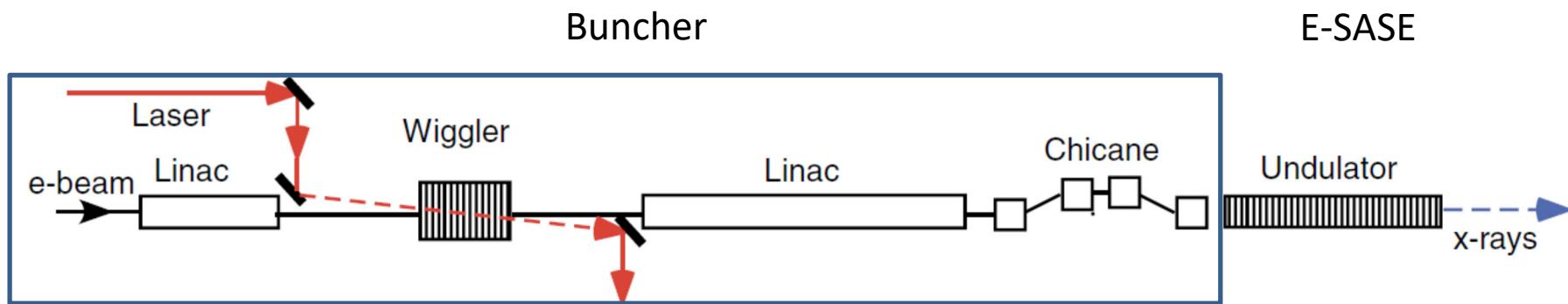


Towards Ultra-Compact X-ray FEL, J. Rosenzweig, UCLA Moore Foundation Workshop 1/22-25/2019

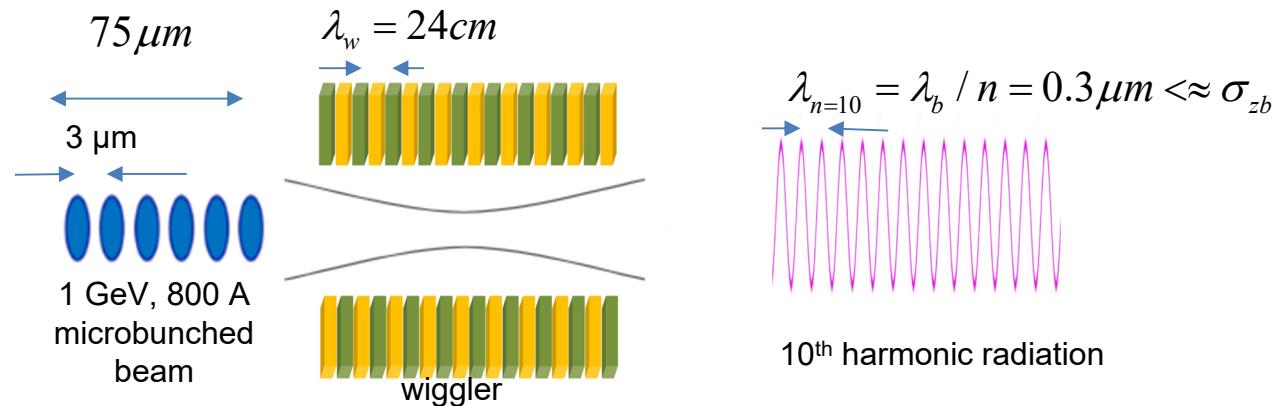
$$E=1GeV \ (\gamma = 2000) \quad \lambda_b = 3.2\mu \quad Q_M = 200pC \quad I_M = 800A$$

$$\text{with } \lambda_w = 1.2mm \Rightarrow \lambda = 1.57\text{\AA}$$

EXERSIZE # 1: HIGH HARMONIC UV SUPERRADIANT SOURCE



A Model Problem with a SAMURAI beam: 10th harmonic SR



$$\bar{a}_w = 10, \quad N_w = 10, \quad L_w = 2.4 m \quad \rightarrow \quad \lambda_w = 2 \frac{\gamma^2}{1 + \bar{a}_w^2} \lambda_{n=10} = 24 cm$$

$$L_w = 2z_R(\lambda_{n=10}) \quad \rightarrow \quad A_{em,q}(\omega_n) = \frac{z_R \lambda_n}{2} = \frac{L_w \lambda_{10}}{4}$$

$$\rightarrow P_{q=TE_{0,0},n=10} = \frac{I_0^2}{8} \sqrt{\frac{\mu_0}{\epsilon_0}} \frac{L_w^2}{A_{em}} = 2.4 GW$$

*RIVER ROBLES - STUDENT UCLA PHYSICS DEPT.

Steady-state (single frequency) phasor formulation for periodic pre- bunching

SINGLE-MODE PHASOR MODEL for SP-SR /ST-SR ($\omega_0 = n\omega_b$)

$$\tilde{C}_q(z) = \tilde{C}_q(0) - \frac{1}{4P_q} \iiint_0^z \tilde{\mathbf{J}}_\perp(\mathbf{r}, \omega_0) \cdot \boldsymbol{\varepsilon}_{q\perp}^*(\mathbf{r}_\perp) e^{-ik_{zq}z} d^2\mathbf{r}_\perp dz$$

$$P_q(z) = P_q \left| \tilde{C}_q(z) \right|^2 \quad P_q(z) = P_q(0) + P_{SR}(z) + P_{ST-SR}(z)$$

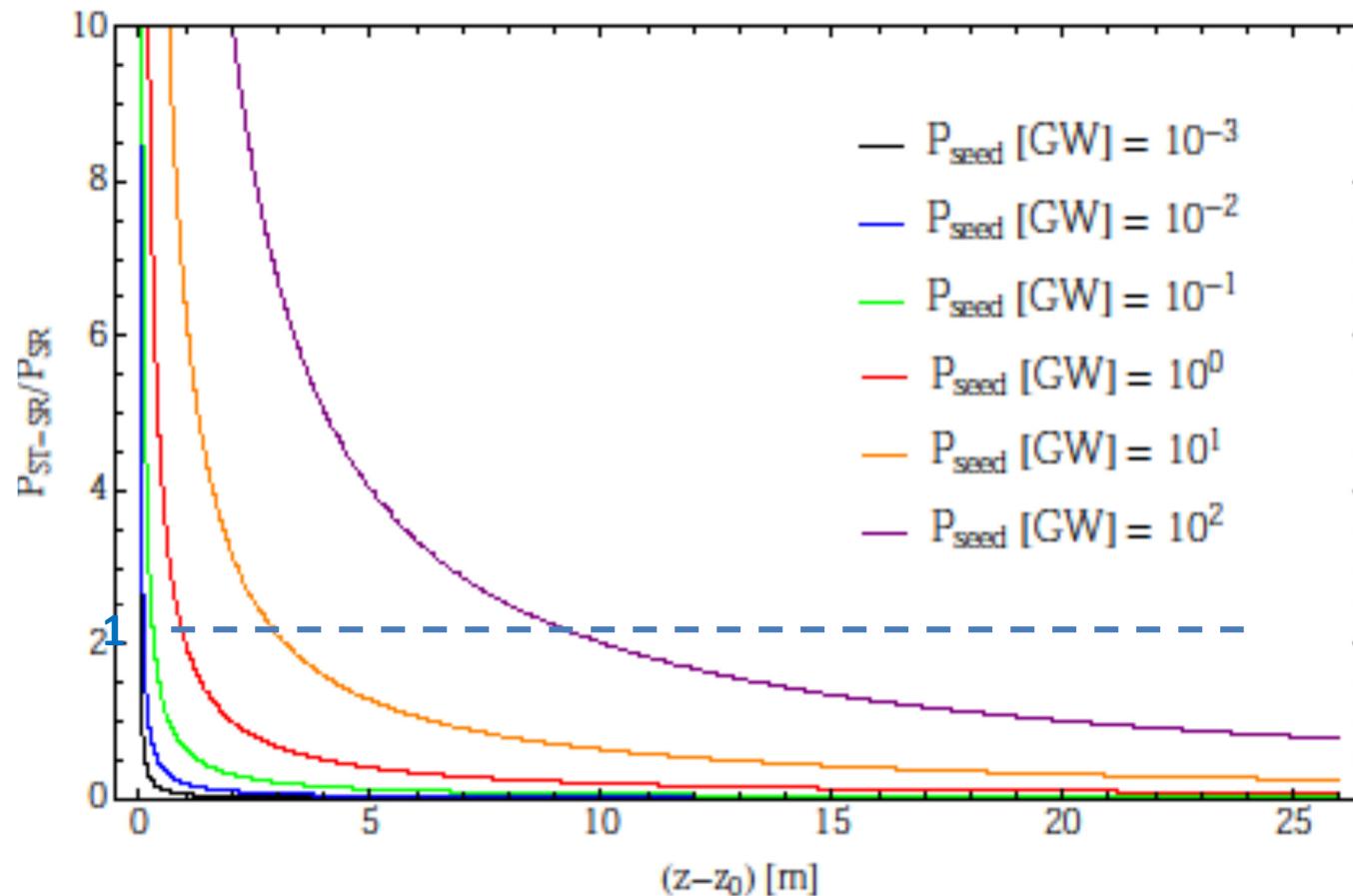
$$P_{SP-SR} = \frac{1}{32} Z_q |\bar{I}_m|^2 F^2 \frac{z^2}{A_{em}} \sin c^2 \left(\frac{\theta_0 z}{2} \right)$$

$$P_{ST-SR} = \frac{1}{4} |\bar{I}_m| |E_\perp(0)| F z \cos \left(\varphi_0^r - \varphi_0^b - \theta_0 z/2 \right) \sin c \left(\frac{\theta_0 z}{2} \right)$$

Relative weight of ST-SR and SP-SR Power

$$\frac{P_{ST-SR}}{P_{SP-SR}} = \frac{\frac{1}{4} \bar{I}_m \sqrt{\frac{2Z_q}{A_{em}}} \sqrt{P_{in}} F_z}{\frac{1}{32} Z_q |\bar{I}_m|^2 F^2 \frac{z^2}{A_{em}}} = 8 \frac{A_{em}}{Z_q \bar{I}_m F_z} \sqrt{\frac{27q}{A_{em}}} \sqrt{P_{in}} = 8 \frac{A_{em}}{Z_q \bar{I}_m F_z} E_{in}(0)$$

Ratio of 0-order ST-SR to SP-SR (calculated for the parameters of the tapered section of LCLS X-FEL)



SR and ST-SR IN THE NONLINEAR REGIME

PERIODIC TIGHT BUNCHING MODEL

$$\mathbf{J}(\mathbf{r}, t) = Q_b \mathbf{v}_e(t) f(\mathbf{r}_\perp) \sum_{n=-\infty}^{\infty} \delta[z - z_e(t - nT_b - t_0)]$$

➡ $\mathbf{J}(z, t) = \mathbf{J}_0 + \sum_{n=1}^{\infty} 2Re \left[\tilde{\mathbf{J}}_n e^{-in\omega_b t} \right]$

$$\frac{d\tilde{C}_q(z)}{dz} = \frac{1}{4P_q} \int \tilde{\mathbf{J}}(\mathbf{r}) \cdot \mathbf{E}_q^*(\mathbf{r}) d^2 r_\perp \quad P_{q,n} = P_q |\tilde{C}_{q,n}|^2$$

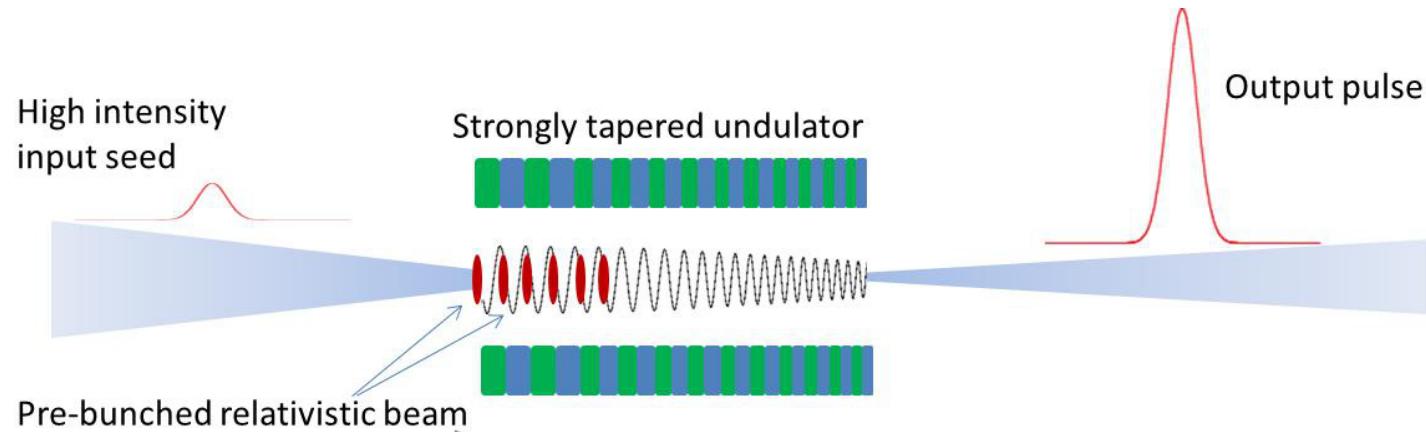
$$N_b m c^2 \frac{d\gamma}{dz} = \frac{Q_b}{\beta_z} \boldsymbol{\beta} \cdot \mathbf{E}(\mathbf{r}, t_e(z)) \quad P_e = N_b m c^2 (\gamma - 1) / T_b$$

Self-consistent nonlinear model formulation for an infinite pulse (or finite pulse with zero-slippage):

$$\frac{dP_{q,n}}{dz} + \frac{dP_e}{dz} = 0$$

NONLINEAR SR AND ST-SR INTERACTION OF A PREBUNCHED BEAM IN AA WIGGLER

TAPERING-ENHANCED STIMULATED SUPERRADIANT AMPLIFICATION - TESSA



N.M. Kroll, P.L. Morton, M.N. Rosenbluth, IEEE J. Quant. Electron., QE-17, 1981

N. Sudar, P. Musumeci et al "High Efficiency Energy Extraction ...Tapered Undulator"
PRL 117, 174801 (2016)

A. Gover, R. Ianconescu, A. Friedman, C. Emma, N. Sudar, P. Musumeci, C. Pellegrini,
"Superradiant and stimulated-superradiant emission of bunched electron beams"
Rev. Mod. Phys. 91, 035003 – Published 19 August 2019

SYNCHRONIZM CONDITION OF A TRAPPED BUNCH

$$\theta(z) = \frac{\omega_0}{v_z(z)} - \overbrace{k_w(z) - k_{zq}(z)}^{pm-wave} \quad (\omega_0 = n\omega_b)$$

$$\theta(z) = -2k_w(z) \frac{\delta\gamma}{\gamma_r(z)}$$

Controlled trap dynamics
 $\gamma = \gamma_r(z) + \delta\gamma$

Resonant energy of a trapped bunch:

$$\theta(z) = 0 \implies \gamma_r^2(z) = \frac{1 + \bar{a}_w^2(z)}{2} \frac{k_0}{k_w(z)} \quad (\bar{a}_w(z) = \frac{e\bar{B}_w(z)}{k_w(z)mc})$$

Phase of ponderomotive (pm) wave relative to bunches:

$$\Psi(z) = - \int_0^z \theta(z') dz' + \Psi(0) + \underbrace{[\varphi_q(z) - \varphi_q(0)]}_{\approx 0} \quad \left(\tilde{C}_q(z) = |\tilde{C}_q(z)| e^{i\varphi_q(z)} \right)$$

RADIATION EMISSION AND BUNCH DYNAMICS – UNIFORM WIGGLER

Single mode: $\tilde{E}(z) = \tilde{C}_q(z) |\tilde{\varepsilon}_{q\perp}(0)|$ $\gamma_r = const$

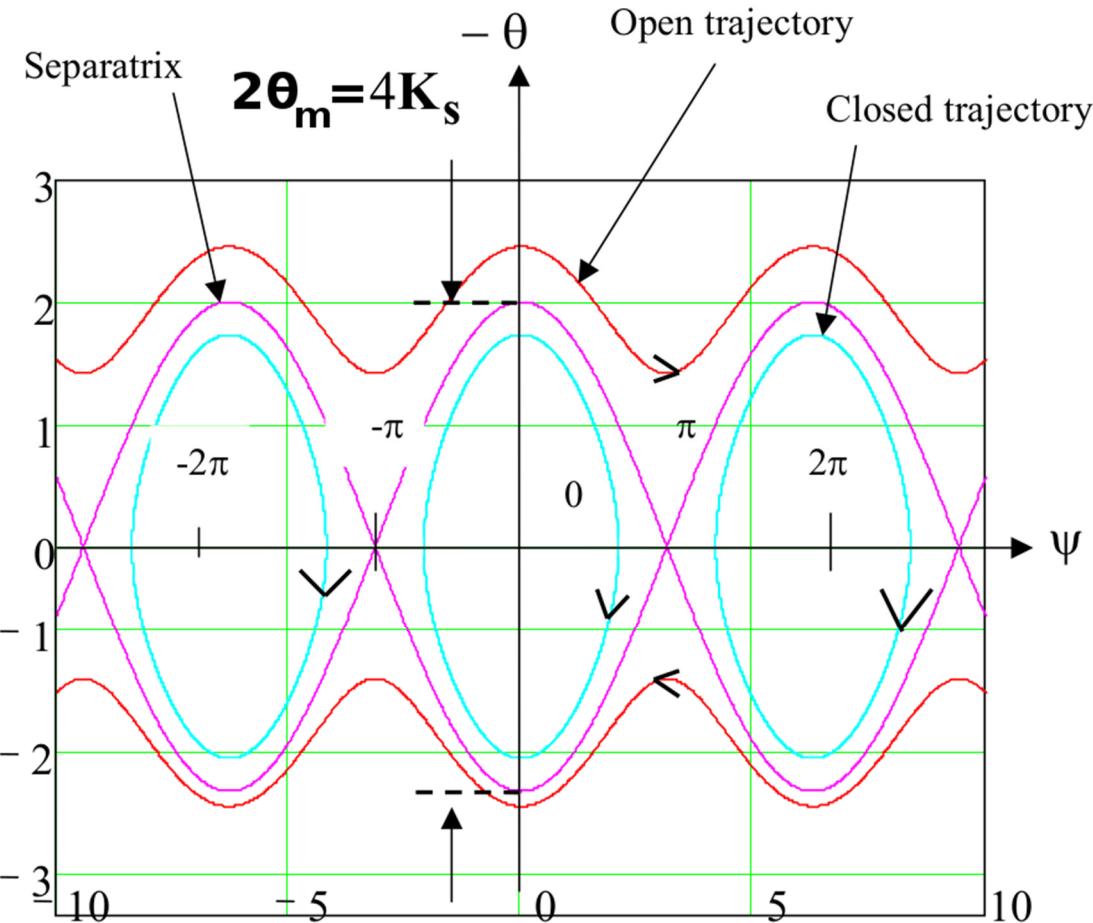
$$\frac{d|\tilde{E}|}{dz} = b \sin \Psi$$

$\frac{d\theta}{dz} = K_s^2(z) \sin \Psi$	$\frac{d\delta\gamma}{dz} = \frac{\gamma_{zr}^2 \gamma_r}{k_0} K_s^2(z) \sin \Psi$
$\frac{d\Psi}{dz} = -\theta + \frac{b}{ \tilde{E} } \cos \Psi$	$\frac{d\Psi}{dz} = \frac{k_0}{\gamma_{zr}^2 \gamma_r} \delta\gamma + \frac{b}{ \tilde{E} } \cos \Psi$

$$K_s^2(z) = \frac{k_0 e}{2\gamma_{zr}^2 \gamma_r^2 m c^2} \bar{a}_w(z) |\tilde{E}(z)|$$

$$b = \frac{\bar{I}_b \bar{a}_w(z) Z_q}{2 A_{em,q} \gamma_r}$$

The Θ - Ψ phase-space trajectories of the pendulum equation



RADIATION EMISSION AND BUNCH DYNAMICS – TAPERED WIGGLER

$$\gamma = \gamma_r(z) + \delta\gamma(z)$$

$$\frac{d\theta}{dz} = \frac{k_0}{\gamma_{zr}^2 \gamma_r} \left(\frac{d\gamma}{dz} - \frac{d\gamma_r}{dz} \right)$$

$$\frac{d|\tilde{E}|}{dz} = b \sin \Psi$$

$$\frac{d\theta}{dz} = K_s^2(z) [\sin \Psi - \sin \Psi_r]$$

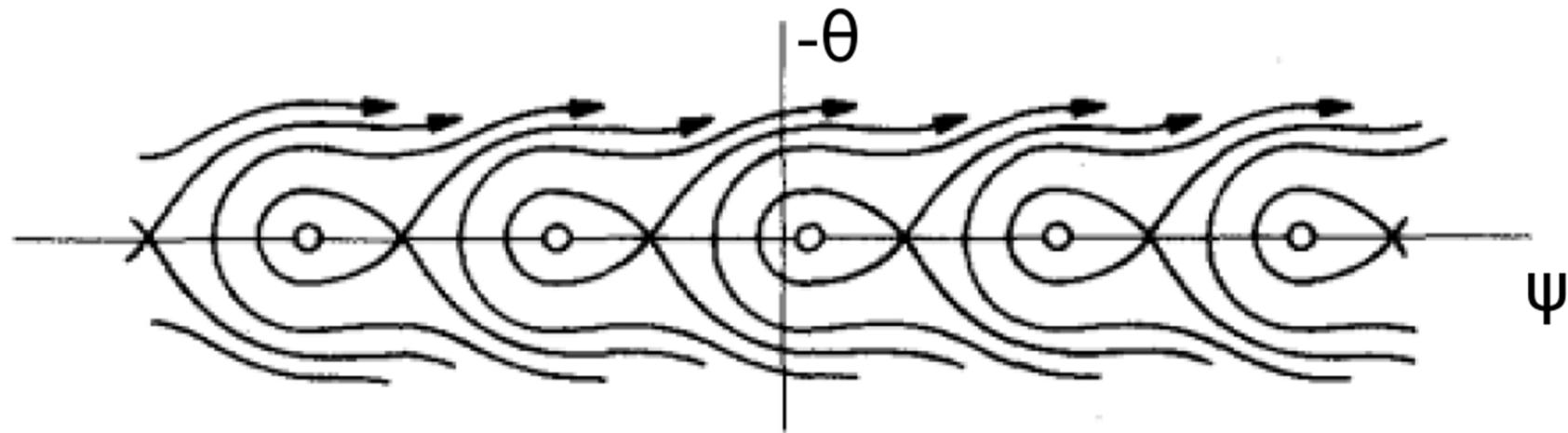
$$\frac{d\Psi}{dz} = -\theta + \frac{b}{|\tilde{E}|} \cos \Psi$$

$$\frac{d\delta\gamma}{dz} = \frac{\gamma_{zr}^2 \gamma_r}{k_0} K_s^2 K_s^2(z) [\sin \Psi - \sin \Psi_r]$$

$$\frac{d\Psi}{dz} = -\theta + \frac{b}{|\tilde{E}|} \cos \Psi$$

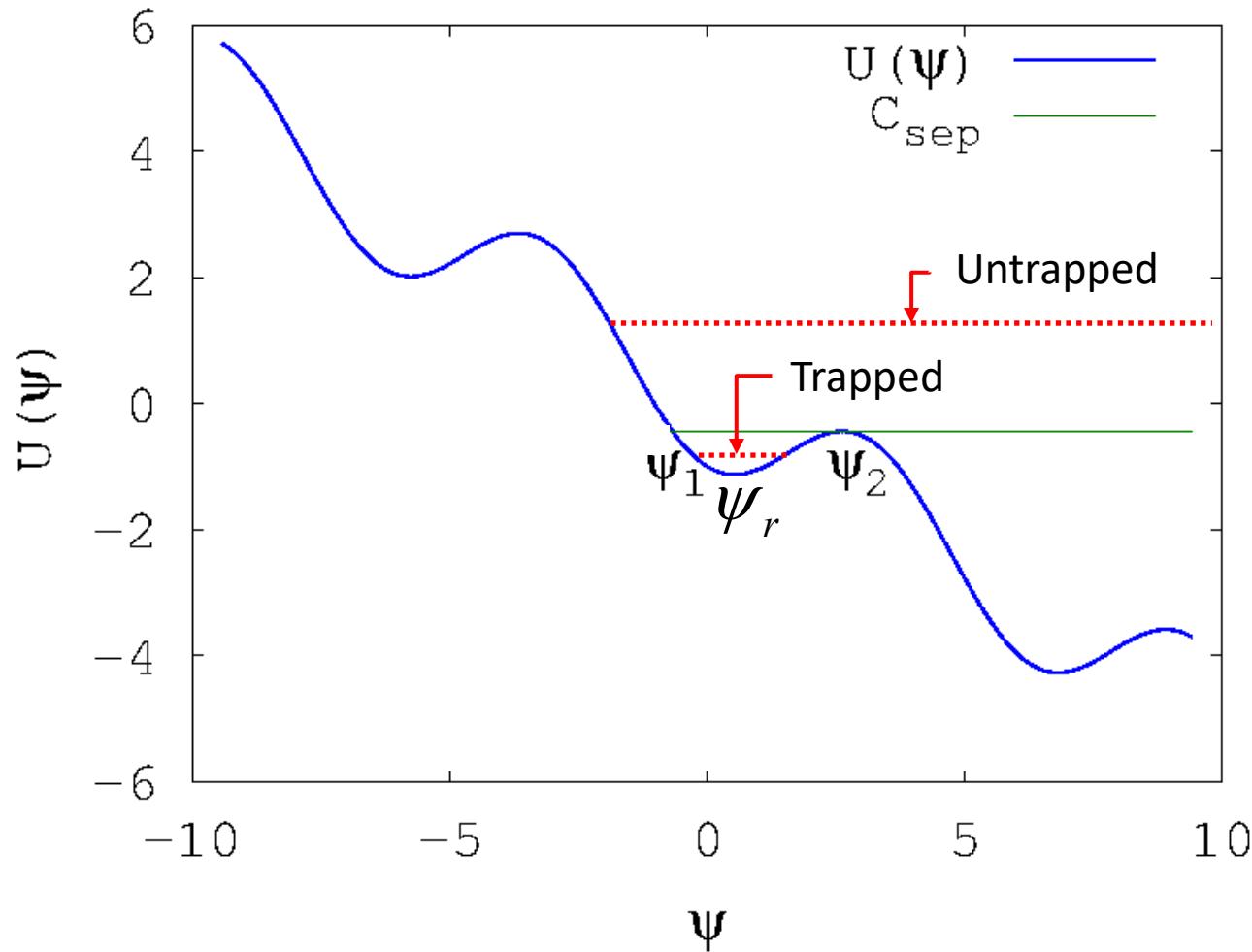
$$\sin \Psi_r = \frac{k_0}{\gamma_{zr}^2 \gamma_r K_s^2} \frac{d\gamma_r}{dz}$$

PHASE-SPACE TRAJECTORIES IN A TILTED PENDULUM

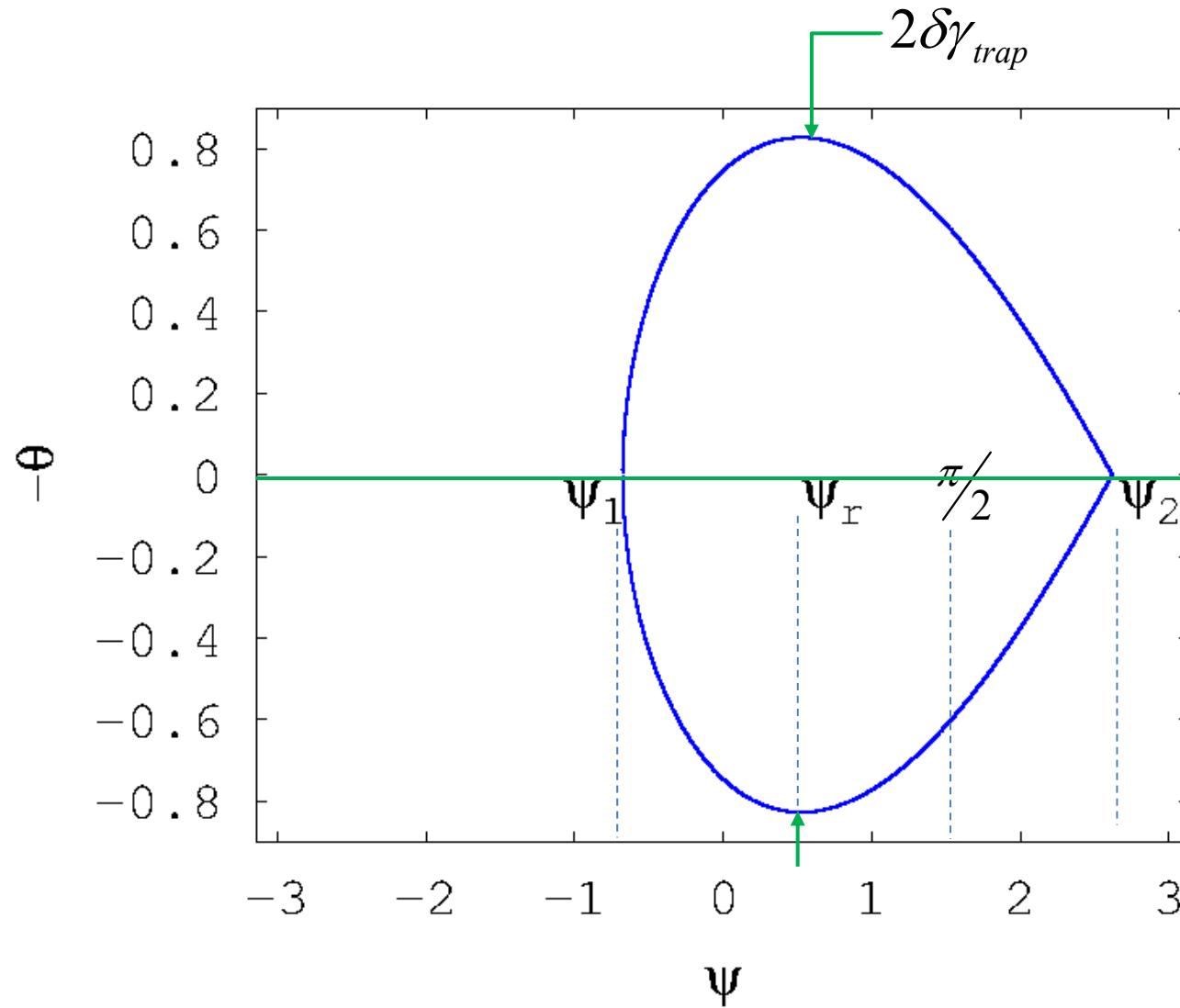


N.M. Kroll, P.L. Morton, M.N. Rosenbluth, "Free-Electron Lasers with Variable Parameter Wigglers", IEEE J. Quant. Electron., VOL. QE-17, NO. 8, AUGUST 1981

The potential energy for a tapered wiggler



SEPARATRIX OF TRAP IN A TAPERED WIGGLER



ST-SR IN THE NONLINEAR REGIME

2nd ORDER PERTURBATIVE SOLUTION

For short interaction length in normalized parameters ($u = z / L_w$):

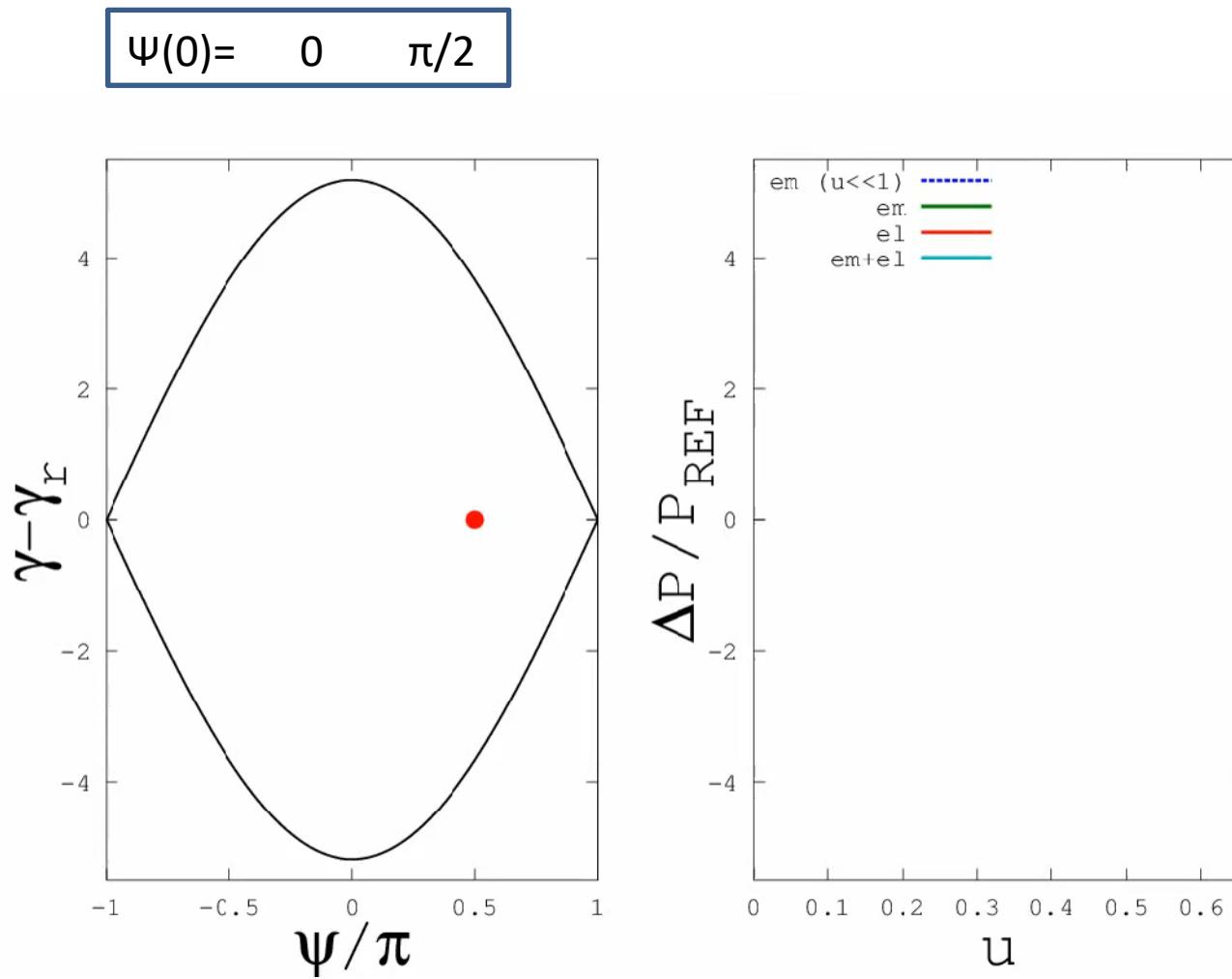
$$\begin{aligned}\Delta P_q / P_{REF} &= \bar{E}^2(u) - \bar{E}^2(0) \\ &= [2u\bar{E}(0)(\sin \psi(0) - \sin \psi_r) + u^2 \sin^2 \psi(0)] + 2u\bar{E}(0)\sin \psi_r\end{aligned}$$

ST-SR SR TAPERING

For $\psi(0) = \psi_r$:

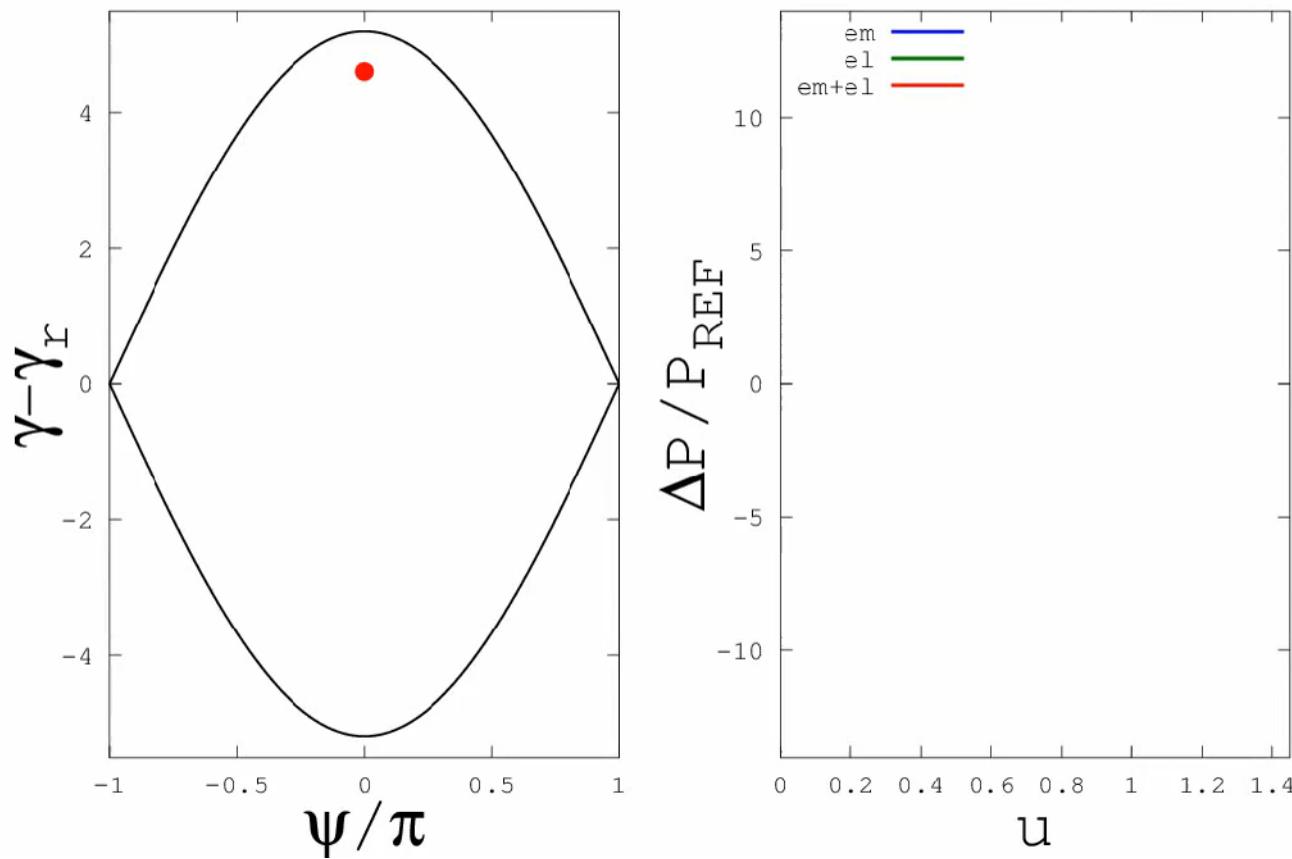
$$\Delta P_{em} / P_{REF} = 2u\bar{E}(0)\sin \psi_r + u^2 \sin^2 \psi_r$$

UNIFORM WIGGLER: HIGH GAIN ST-SR



Uniform wiggler: maximal extraction

$$\Psi(0) = 0 \quad \pi/2$$

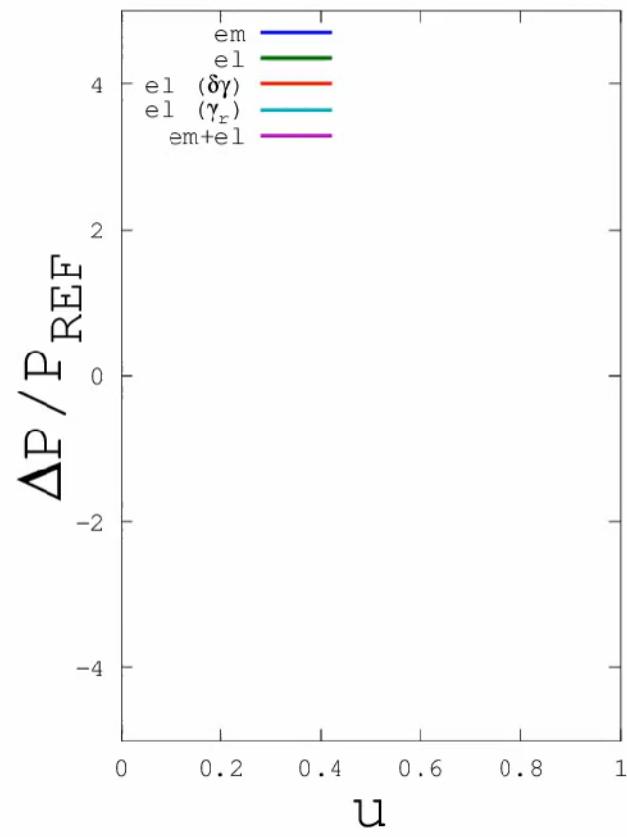
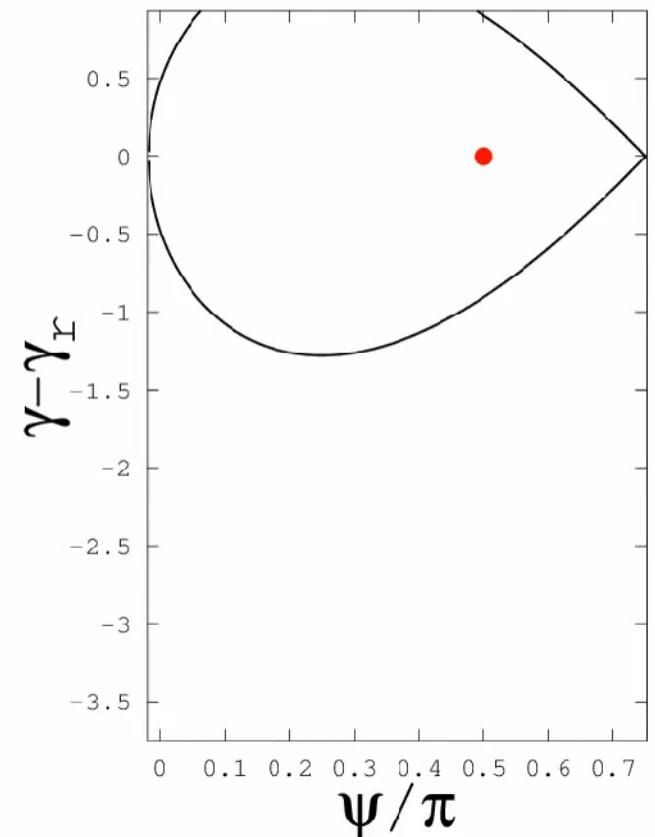


Tapered wiggler: maximal gain ST-SR

$$\Psi(0) = 0$$

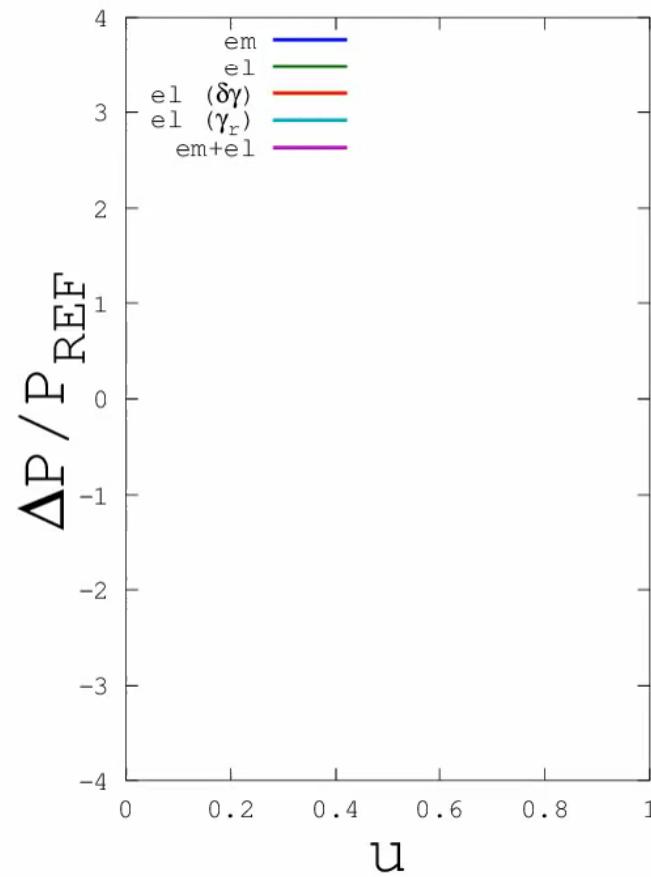
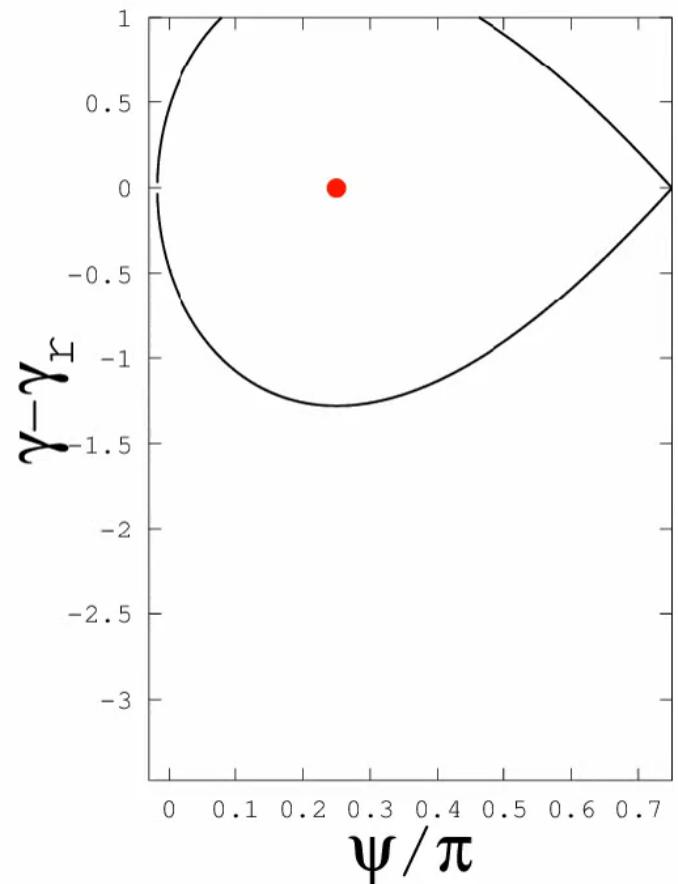
$$\Psi_r$$

$$\pi/2$$



Tapered wiggler: best trapped bunch

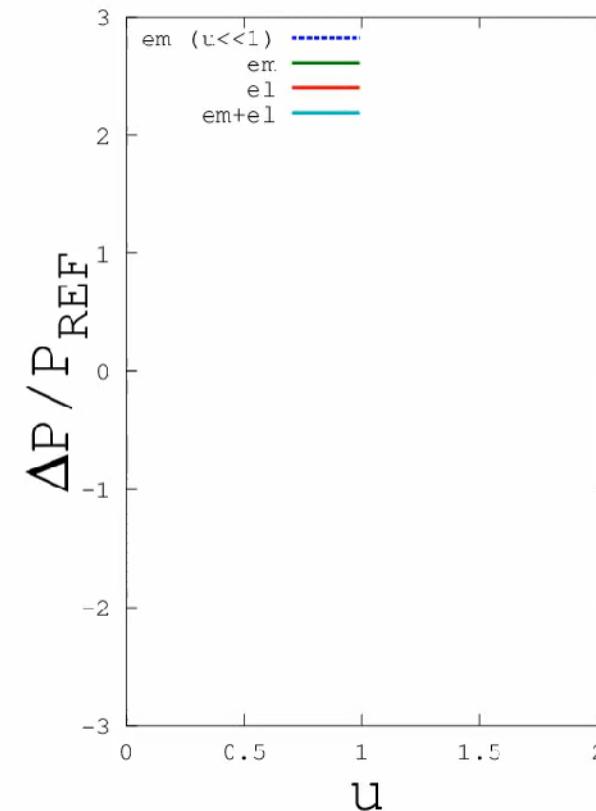
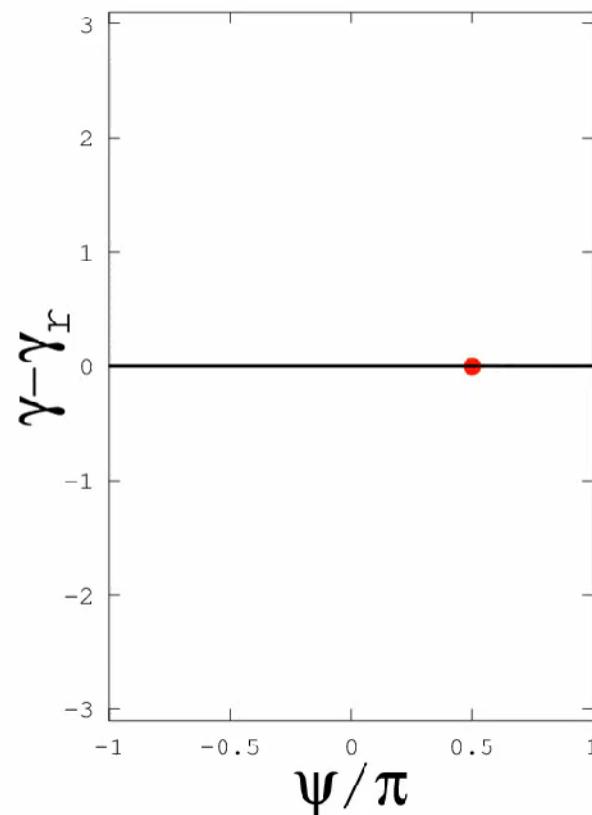
$$\Psi(0) = 0 \quad \Psi_r \quad \pi/2$$



Uniform wiggler: self-interaction

$$\Psi(0) = \begin{matrix} 0 \\ \pi/2 \end{matrix}$$

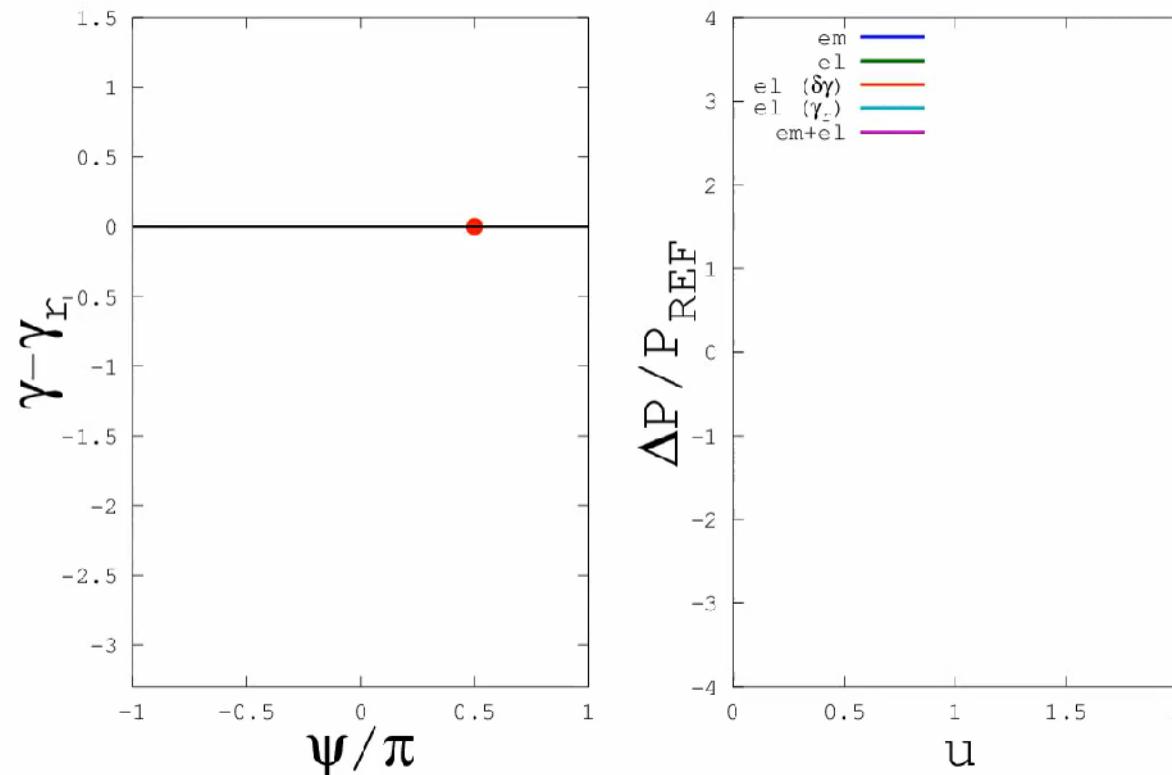
(NO RADIATION REACTION FORCE IS INVOKED)



Abraham-Lorentz Radiation Reaction, Dirac, P. A., 1938, Proc. R. Soc. A 167, 148.

SELF INJECTION: SR TO TESSA

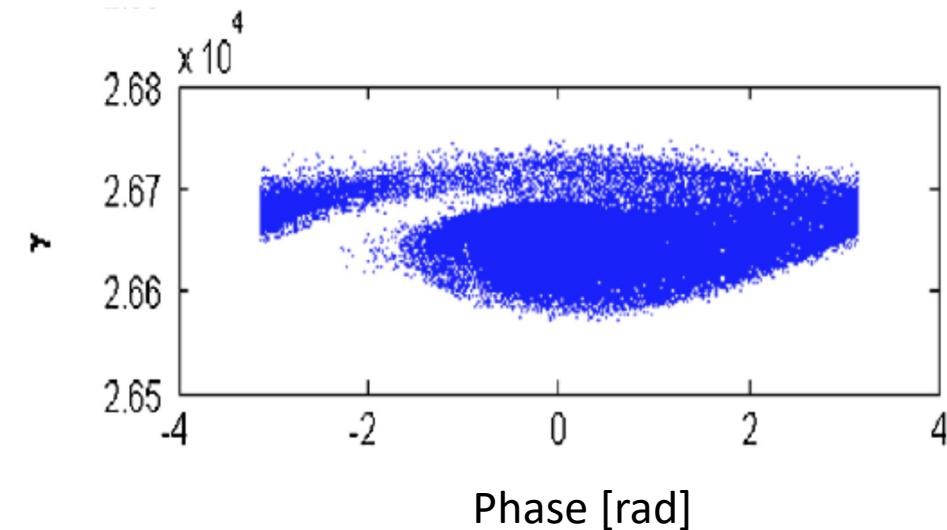
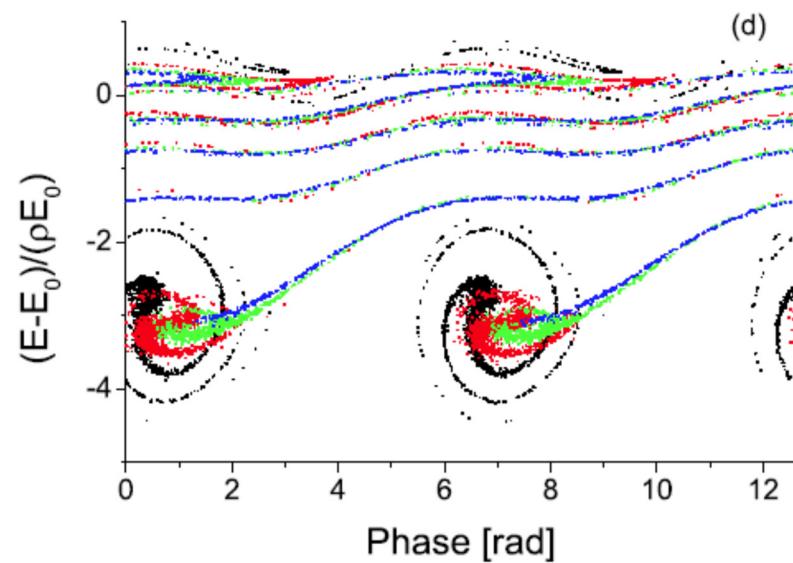
$$\Psi(0) = 0 \quad \pi/2$$



Snively, E. C., Xiong, J., Musumeci, P., & Gover, A. *Optics Express*, 27(15), 20221(2019).
Broadband THz amplification and superradiant spontaneous emission in a waveguide FEL.

EXTENTION TO DISTRIBUTED BUNCH

Phase – space trajectories of a realistic bunch in a tapered wiggler FEL



N. M. Kroll, P. L. Morton, M. N. Rosenbluth, IEEE J. Quant. Electron. **17**, 1436 (1981)

Y. Jiao et al, "Modeling and multidimensional optimization of a tapered free electron laser" PRST-AB **15**, (050704) (2012)

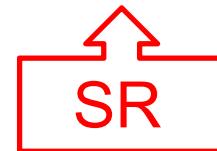
E. A. Schneidmiller, M. V. Yurkov, "Optimization of a high efficiency FEL amplifier" PRST-AB **18**, 03070 (2015)

TESSA: EXTENTION TO DISTRIBUTED BUNCH

Bunches get trapped in the tapered undulator in ponderomotive buckets.

f_t : fraction trapped.

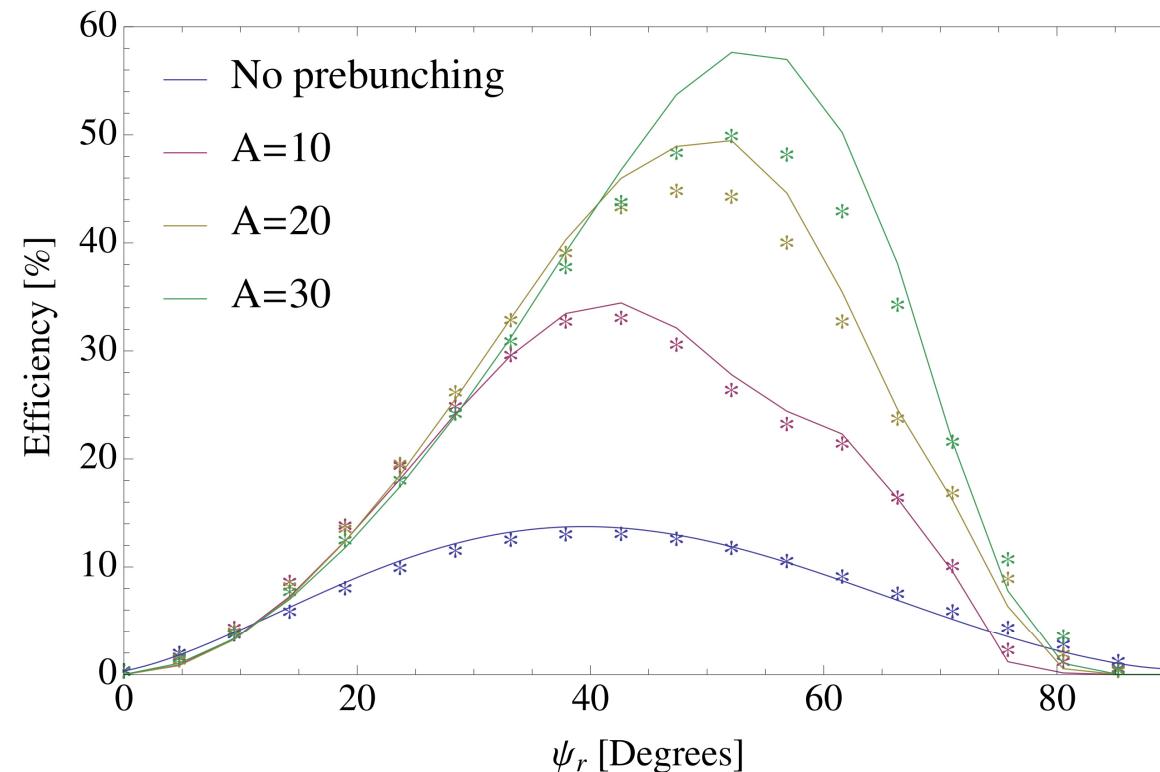
$$P_{rad}(z) = P_0 + E_0 \frac{\bar{a}_w(0)}{\gamma_0} f_t I_z \sin \psi_r + \frac{Z_0}{4A_{emq}} \left(\frac{\bar{a}_w(0)}{\gamma_0} \right)^2 (f_t I_z \sin \psi_r)^2$$



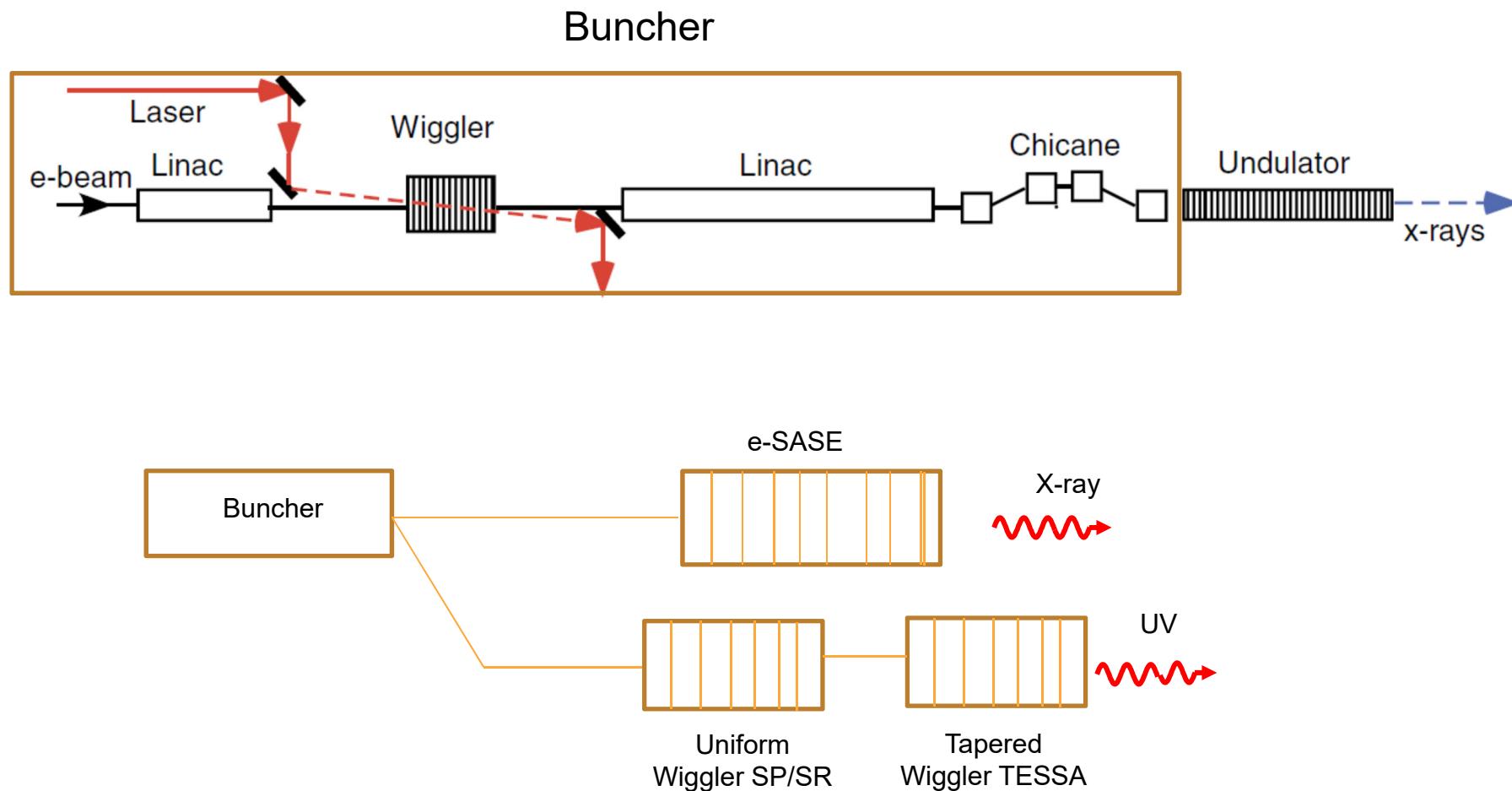
Need to optimize $f_t(\psi_r, \Delta\gamma_{mod}) \sin \psi_r$

TRAPPING EFFICIENCY

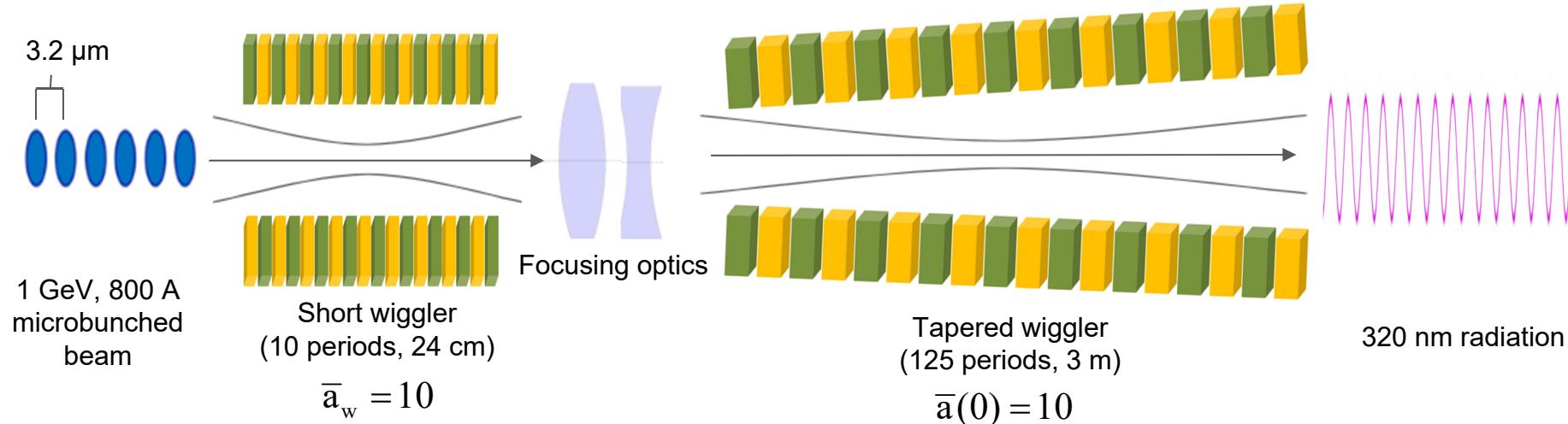
$$f_t = \int_{\Psi_1}^{\Psi_2} d\psi \int_{-\delta\gamma_{trap}/\sigma_{\gamma 0}}^{\delta\gamma_{trap}/\sigma_{\gamma 0}} dp f_0(p, \psi) \quad p = \frac{\gamma - \gamma_0}{\sigma_{\gamma 0}} \quad A = \frac{\Delta\gamma_{\text{mod}}}{\sigma_{\gamma 0}}$$



EXERCISE # 2 - HIGH HARMONIC UV RADIATION TESSA SOURCE



A Model Problem with a SAMURAI beam: 10th harmonic SR and TES*



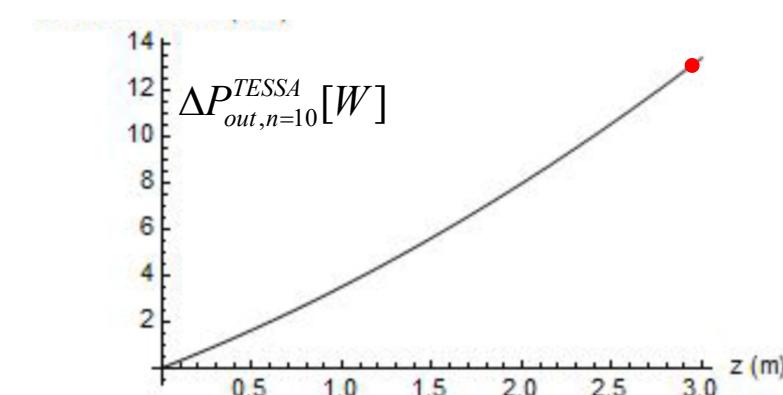
$$\Delta P_{rad}^{TESSA} = E_0 \frac{\bar{a}_w(0)}{\gamma_0} f_t \bar{L}_w \sin \psi_r + \frac{Z_0}{4A_{emq}} \left(\frac{\bar{a}_w(0)}{\gamma_0} \right)^2 (f_t \bar{L}_w \sin \psi_r)^2$$

Summary:

$$P_{in}^{TESSA} = P_{out}^{SR} = 2.4GW$$

$$\Delta P_{out}^{TESSA} = 13.4GW$$

$$P_{out} = P_{in}^{TESSA} + \Delta P_{out}^{TESSA} = 16GW$$



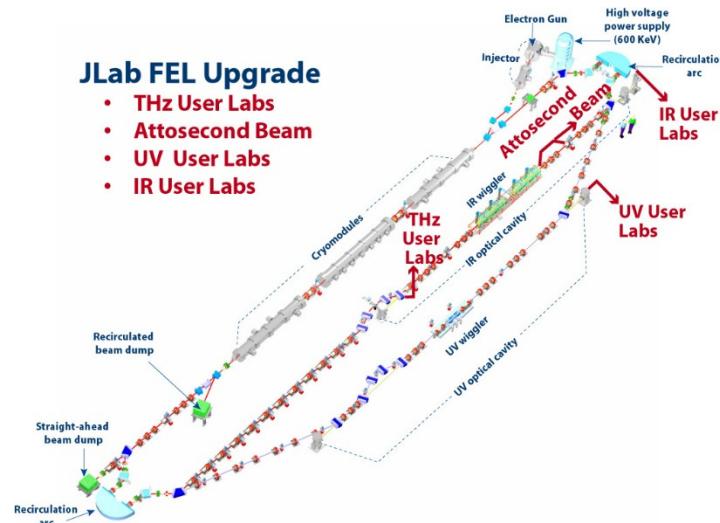
APPLICATION OF SUPERRADIANCE:

Extensive reference list: Gover et al Rev. of Mod. Phys. Vol/EID: 91/035003 (August 2019)

Energy Retrieval LINAC (ERL)

JLab FEL Upgrade

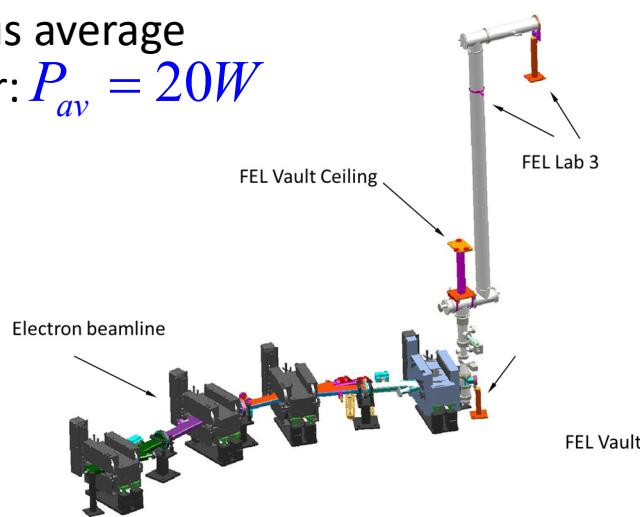
- THz User Labs
- Attosecond Beam
- UV User Labs
- IR User Labs



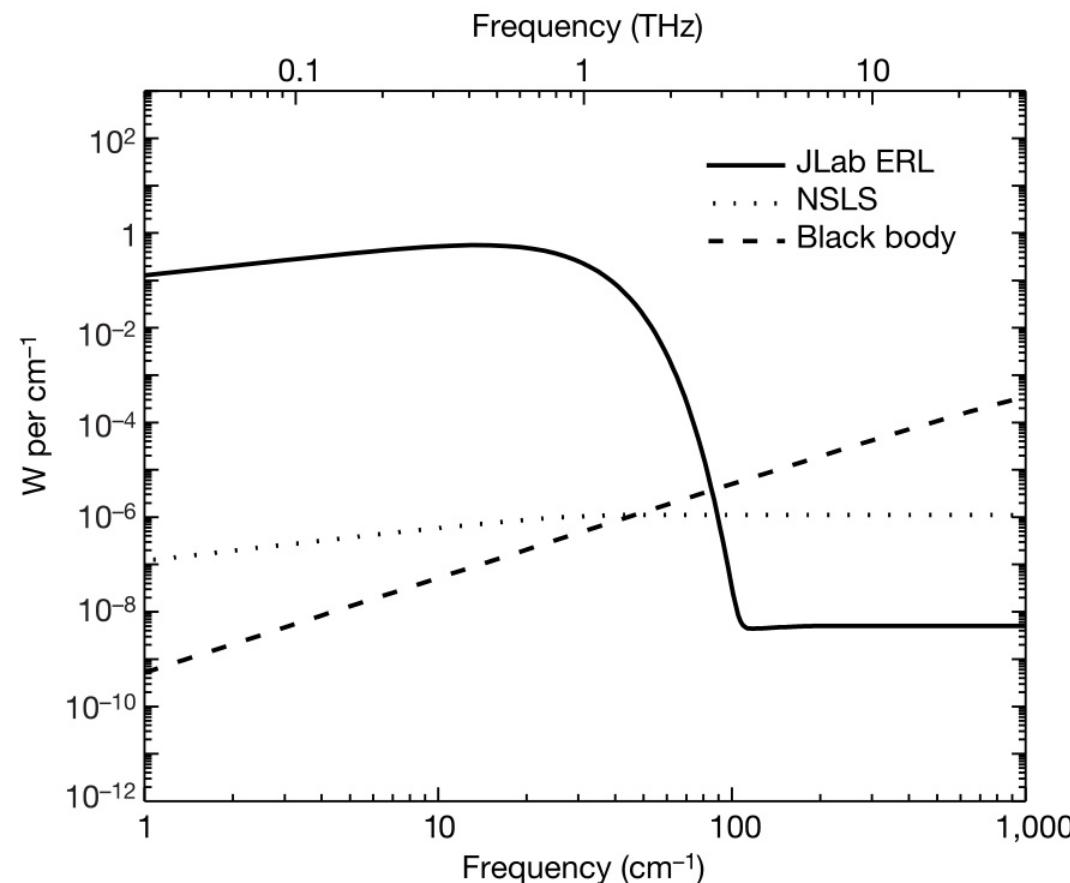
Terahertz beamline transports

Continuous average

THz power: $P_{av} = 20W$



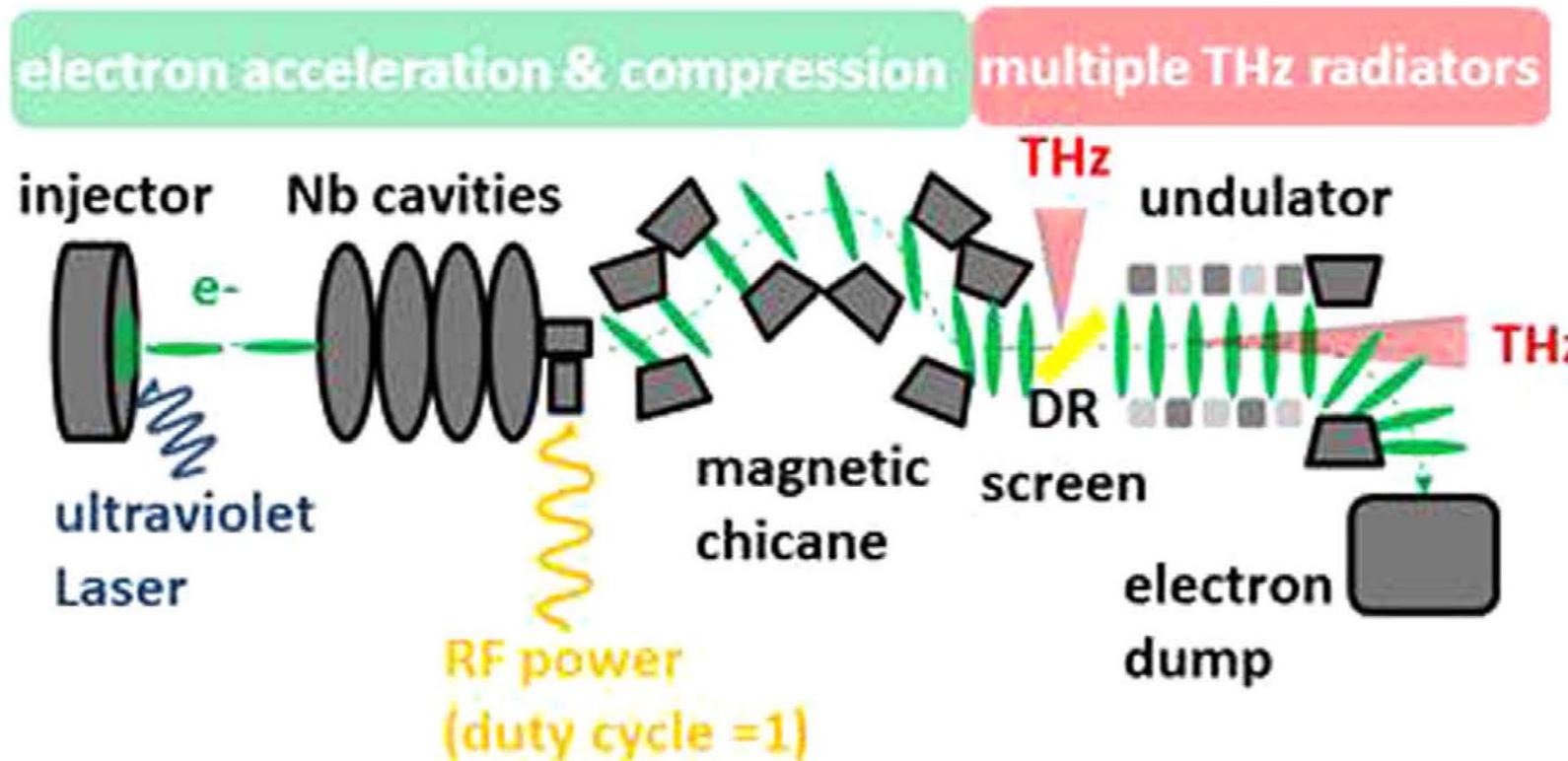
JEFERSON LAB FEL COHERENT SYNCHROTRON RADIATION



Carr, G. L. et al, 2002, Nature 420, 153.

TERA-HERZ ELBE (TELBE) HELMHOLTZ ZENTRUM DRESDEN ROSENDORF

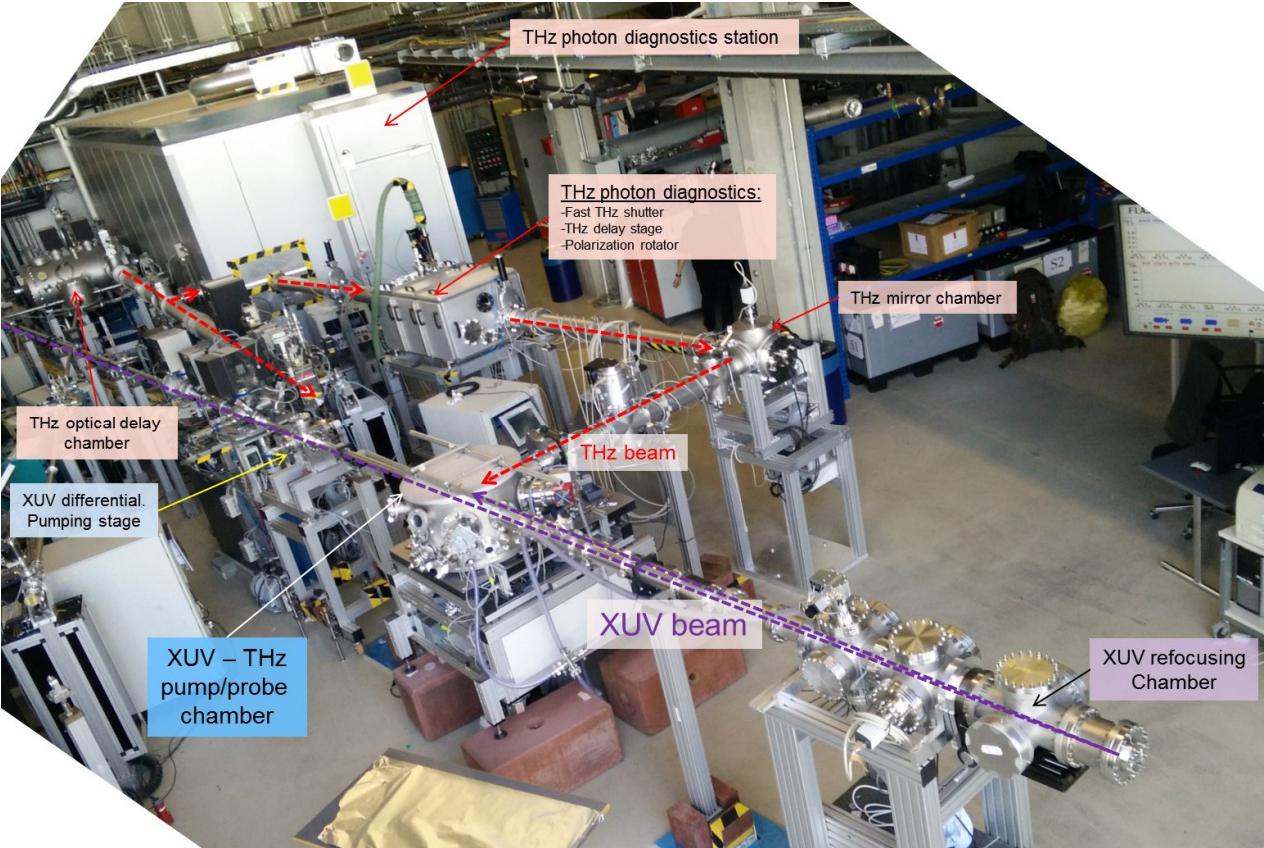
CONTINUOUS TRAIN OF SINGLE PULSES FROM SRF ACCELERATOR THZ SOURCES:
SR UNDULATOR RADIATION,
SR TRANSITION RADIATION



Beam energy: 24 MeV
Charge/pulse: 100 pC
Number of wiggles: 8
THz frequency: Up to 1.4THz
THz energy/pulse: 1 μ J

Green, B., et al., 2016, Sci. Rep. 6, 22256.

FLASH – DESY THz Beamline



tunable: 10 - 300 μm ;
up to 100 $\mu\text{J}/\text{pulse}$;
 \sim 10% bandwidth,

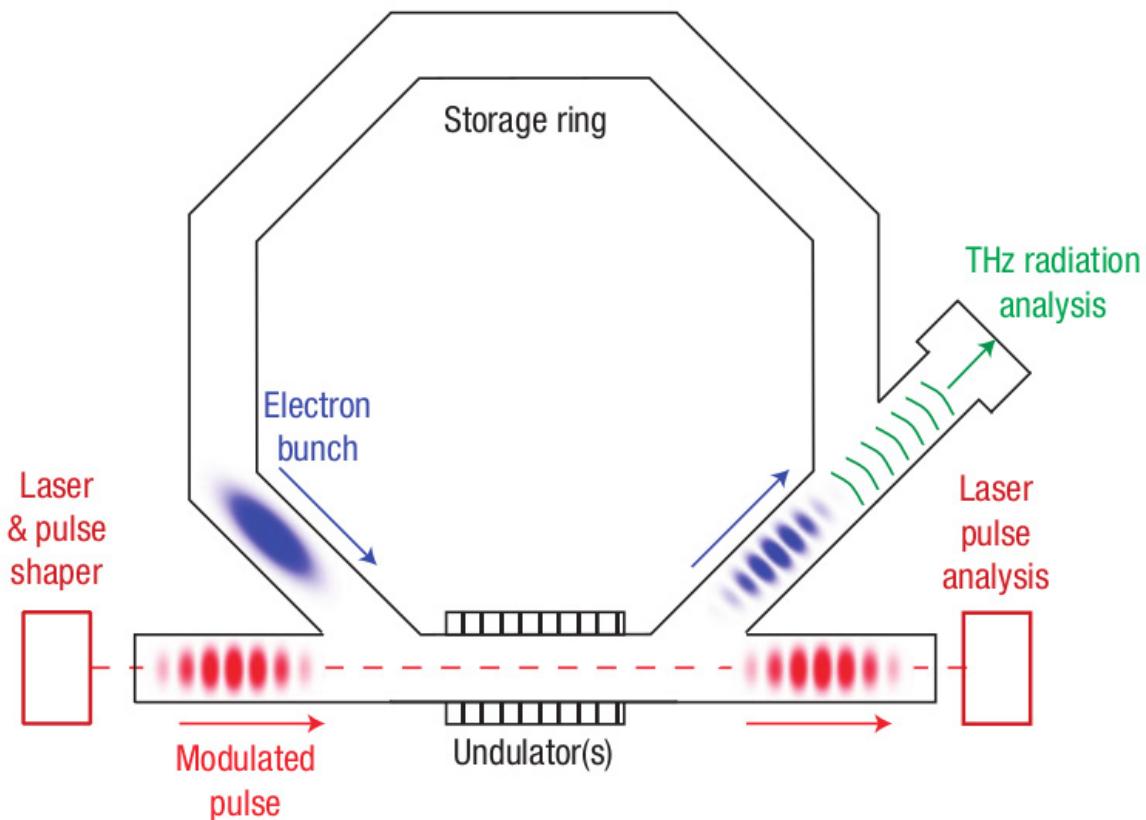
broadband at 200 μm ;
up to 10 $\mu\text{J}/\text{pulse}$;
 \sim 100% bandwidth

B. Faatz et al NIM A 475 (2001) 363

M. Gensch et al [Infrared Physics and Technology](#) 51, 423-425 (2008)

UVSOR OKAZAKI JAPAN

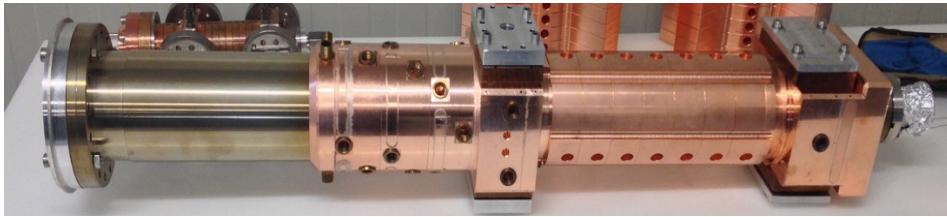
THz SR EMISSION IN A STORAGE RING



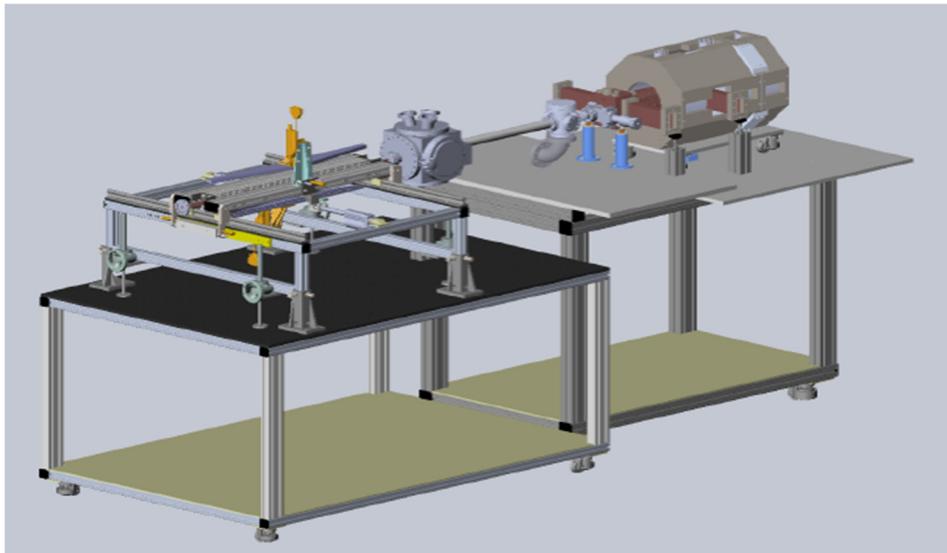
- $E=750\text{MeV}$
- The energy of storage beam pulses is modulated in an undulator by interaction with a THz modulated laser beam. They emit Synchrotron SR at the bend.

Bielawski, S., et al., 2008, Nat. Phys. 4, 390.

TeraHertz Superradiant FEL – ARIEL/TAU (ISRAEL)



Hybrid photocathode gun PBPL UCLA
Alesini et al, EPAC, Edinburgh 2006



Status of the FEL Facility in Israel
A. Friedman et al, ID: 2151 - THP077

Delhi Light Source (DLS): A Compact FEL-THZ facility - Subhendu Ghosh

RF-Linac cavity can work for:

Electron Energy at the exit: 3-6.5 MeV

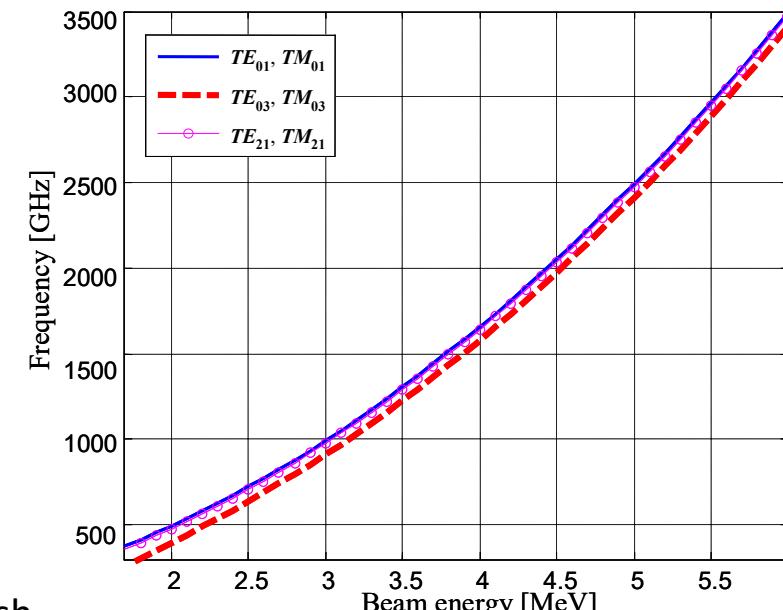
Macrobunch charge: 1 nC

Microbunch duration sub-pSec

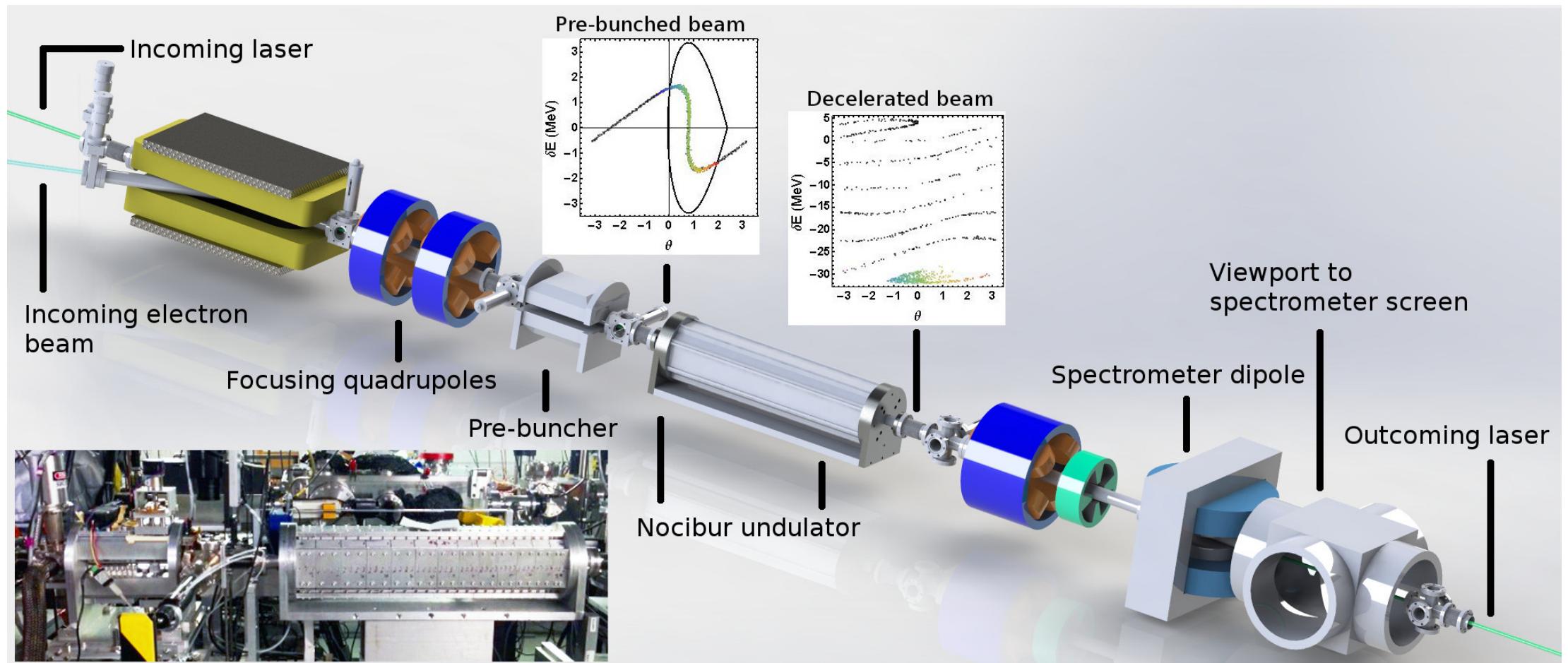
Microbunch Rep.Rate: single/3.5 THz

Macrobunch Duration: 12 ps

Macropulse Rep.Rate: 100 Hz

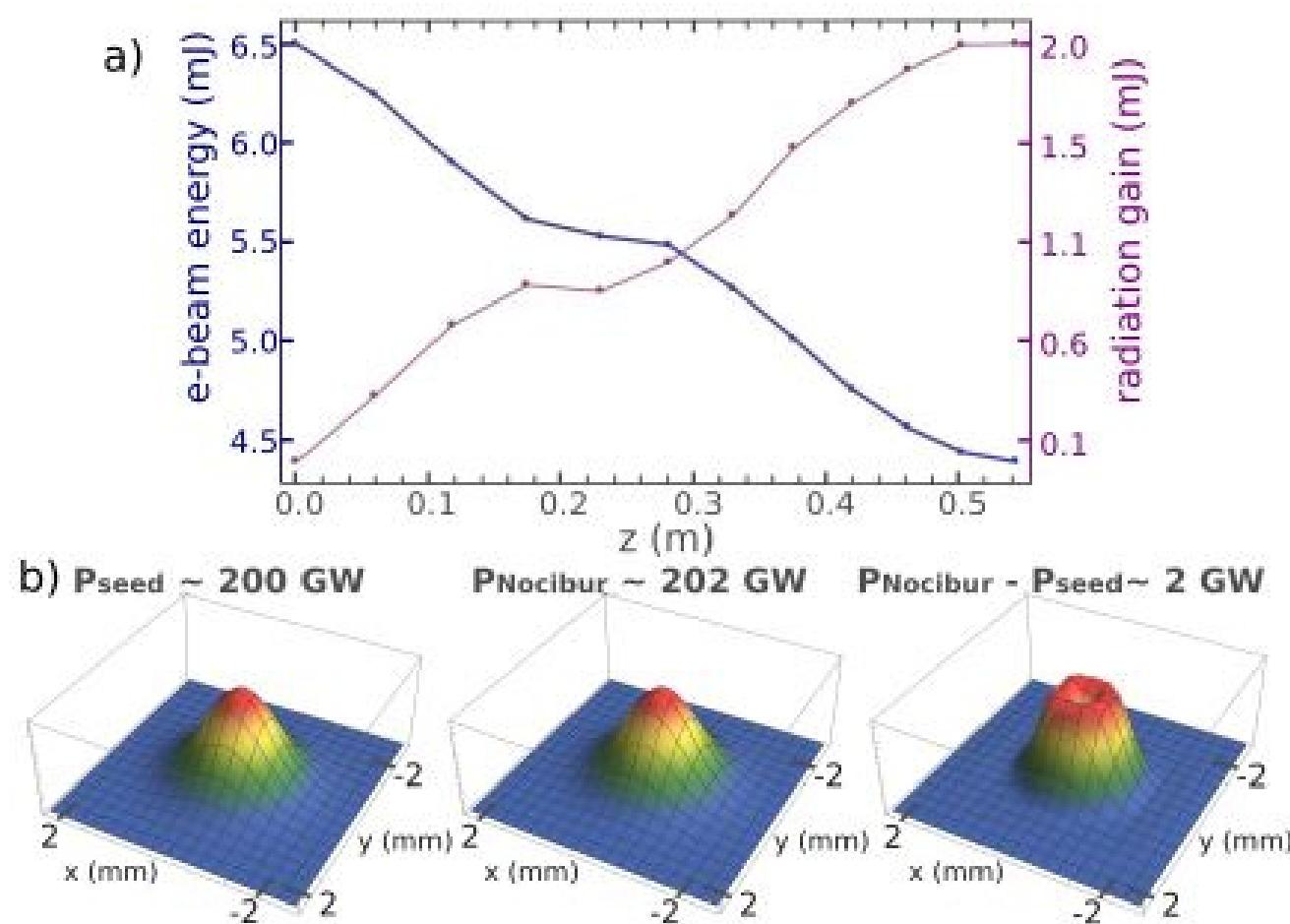


NOCIBUR TESSA EXPERIMENT UCLA-ATF

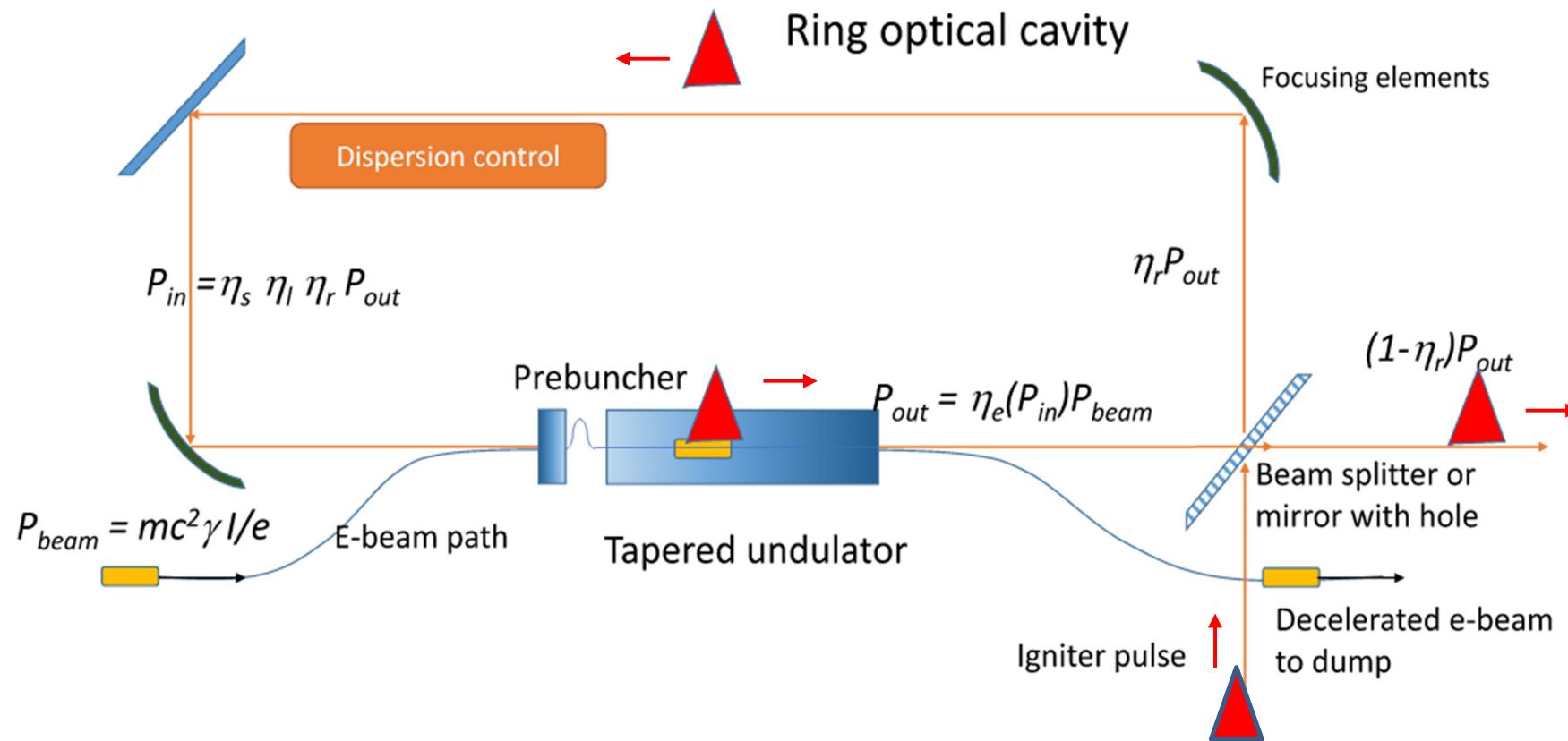


Sudar, et al *Physical review letters* 117.17 (2016): 174801.

30% RADIATIVE ENERGY EXTRACTION EFFICIENCY IN THE NUCIBUR EXPERIMENT @ $\lambda = 10.5\mu$



TAPERING ENHANCED STIMULATED SUPERRADIANCE OSCILLATOR (TESSO)

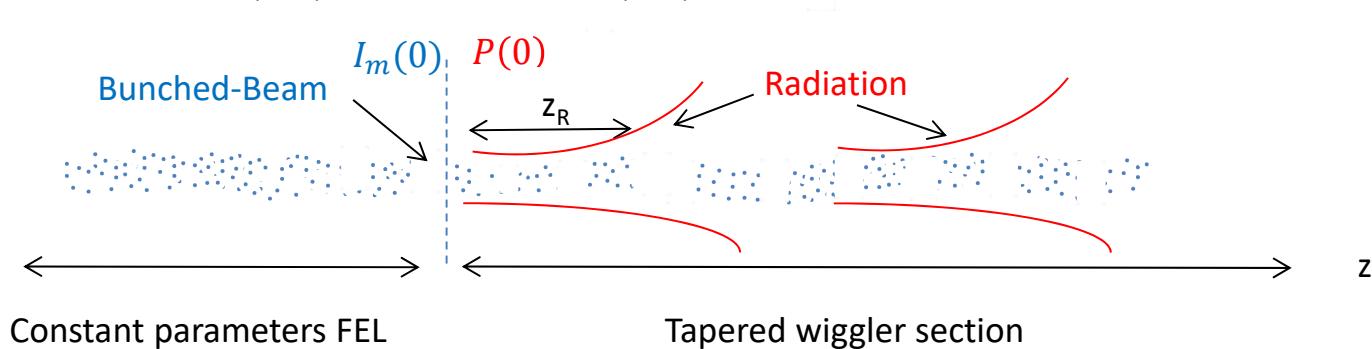


J. Duris, P. Musumeci, N. Sudar, A. Murokh, and A. Gover. Phys. Rev. Accel. Beams **21**, 080705 (2018)

P. Musumeci, "Efficiency and High Gain Amplification at 266 nm" THP073 Thursday Poster Session, 29.8.19.

Tapered wiggler optimization in seed injected FEL

D. Prosnitz et al PRA (1981); Scharlemann T. et al., PRL (1981)



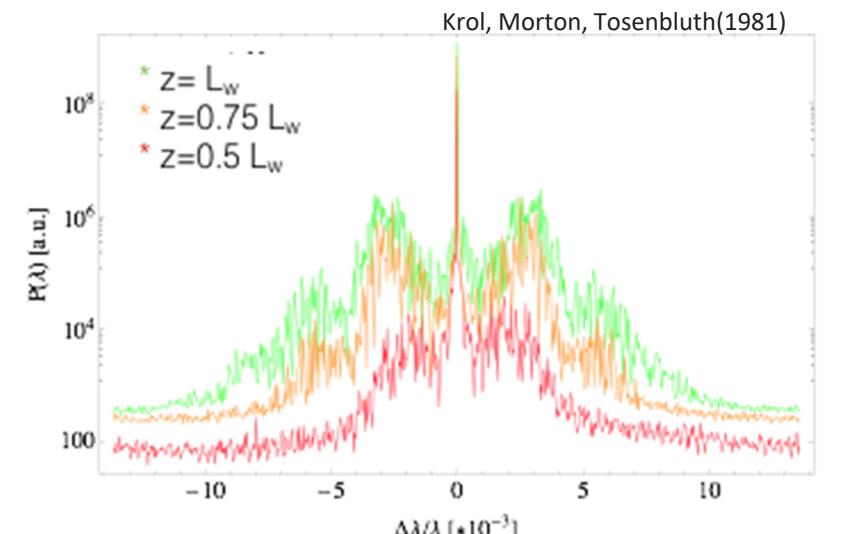
Effects degrading the fundamental processes:

- 1d theory [1]
- Diffraction[2-7]
- Sideband instability [1,3,4,7,...]
- Spectral effects; Shot-Noise [9,10]
- Phase space spread of injected beam [3,7,8]

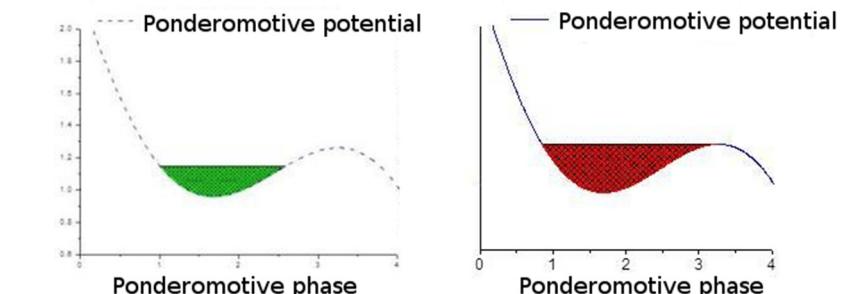
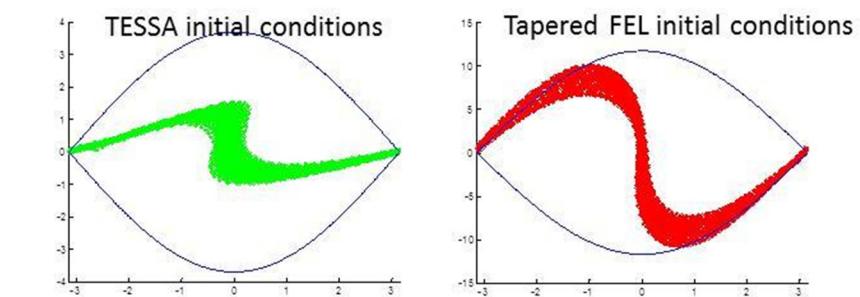
Evolving solutions:

- Fresh bunch scheme [11,12]
- Phase shifter [13,14]
- Gain modulation [8]
- Shot-noise suppression [9,10]
- Experimental tests [15]

[1] KMR (1981); Bonifacio, Casagrande (1988) [2] Fawley (1996); [3] Jiao et al (2012); [4] Emma, Pellegrini (2014);
[5] Schneidmiller, Yurkov (2015); [6] Tsai, Wu et al (2018); [7] N.Sudar (2019) [8] Emma, Sudar (2017) [9] Gover and
Dyunin (2009); [10] Ratner, Huang, Stupakov (2011) [11] Ben-Zvi et al (1992) [12] Emma C. et al (2017) [13] Ratner,
D., et al., 2010 [14] Duris, Murokh, and Musumeci (2015) [15] Wu (2017); N. Sudar et al., (2018)



Duris, Murokh, and Musumeci (2015)



Conclusions

1. Fundamental radiation emission processes of bunched beam at zero order:
-Spontaneous Superradiance (SP-SR) $\propto N^2, z^2$
-Stimulated Superradiance (ST-SR) $\propto N, z, E_0$
2. Model of periodical tightly bunched e-beam interaction with a single radiation mode in the nonlinear regime:
-SR and ST-SR in a uniform wiggler
-Tapering Enhanced Superradiance (TES), Tapering Enhance Stimulated Superradiance Amplification (TESSA) and Oscillator (TESSO).
3. Self interaction of a bunched beam in a uniform wiggler and seedless TESSA.
4. Application in THz superradiant sources based on SR emission of sub-picoSec bunches.
5. Applications of TESSA, TESSO in the THz to UV frequencies range.
6. Optimization of tapering strategy in the tapered wiggler section of X-Ray FELs.