

# Matched Distributions in Cyclotrons with Higher Order Moments of the Charge Distribution

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# Outline

- 1 Introduction
- 2 Computation of Non-Linear Space Charge Map
- 3 Updating Moments
- 4 Conclusion & Outlook

# Task Description

[Fre15]

- **Goal:** Find **non-linear** mapping  $\mathcal{M}$  and  $\sigma^{(k)}$  such that

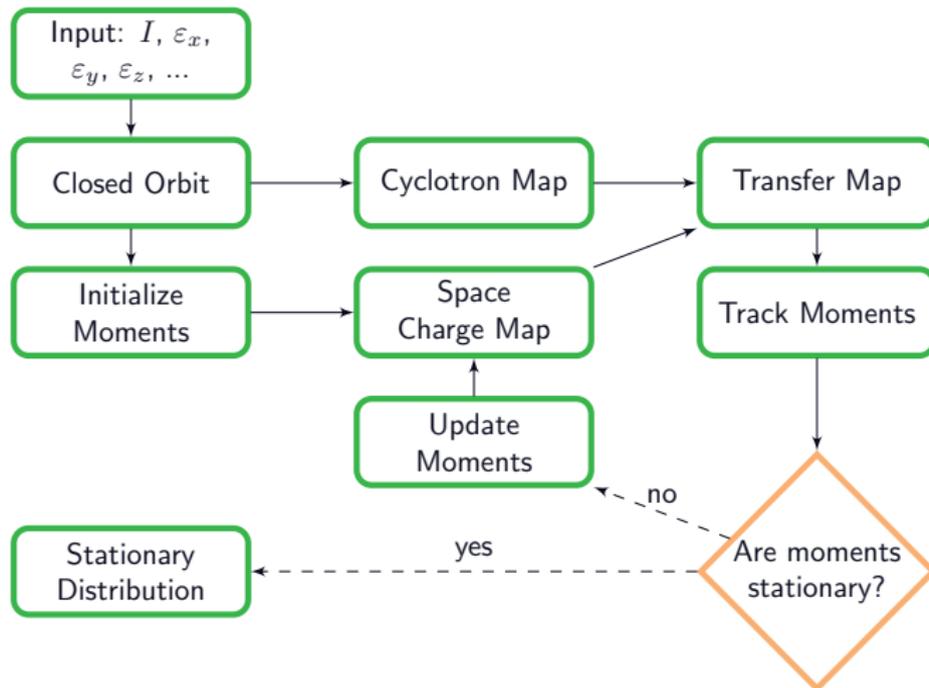
$$\begin{aligned}\sigma^{(k)}(s) &= \mathcal{M} \circ \sigma^{(k)}(0) \circ \mathcal{M}^{\mathcal{T}} \text{ and} \\ \sigma^{(k)} &= \sigma^{(k)}(s) \stackrel{!}{=} \sigma^{(k)}(0)\end{aligned}\tag{1}$$

where  $\sigma^{(k)}(t)$  are the  $k$ -th order moments of a distribution in 6-dimensional phase space and  $\circ, \mathcal{T}$  are operations.

- **Input:** Energy, emittances, intensity, field map, **higher order moments**
- **Linear Theory:**
  - C. Baumgarten. Transverse-longitudinal coupling by space charge in cyclotrons. Phys. Rev. ST Accel. Beams, 14:114201, Nov 2011.
  - C. Baumgarten. Geometrical method of decoupling. Phys. Rev. ST Accel. Beams, 15:124001, Dec 2012.
- **Tests:** PSI Injector-2, PSI Ring Cyclotron (coasting beam)

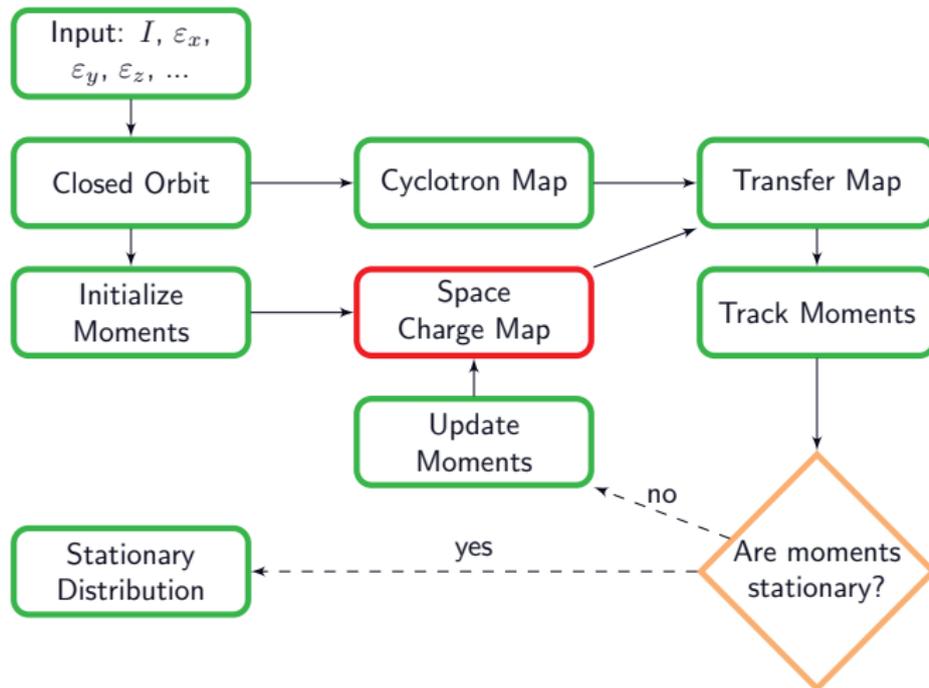
# Big Picture

## Process of finding a Matched Distribution



# Map Generation

## Context



# From Hamiltonian to Map

- Hamiltonian of particle accelerators

$$H = H_{mag} + H_{sc},$$

with magnetic part  $H_{mag}$  and space charge part

$$H_{sc} = q\phi,$$

where  $q$  is the charge and  $\phi$  the self-field potential.

- **Assumptions:**
  - No collisions, no residual gas, no walls
  - Linear magnetic force ( $\rightarrow H_{mag}$  2nd order)

# From Hamiltonian to Map

- Lie Algebra to obtain transfer maps in

$$\mathcal{M} = e^{-s :H:} \quad \text{with} \quad :H: = \frac{\partial H}{\partial \vec{q}} \frac{\partial}{\partial \vec{p}} - \frac{\partial H}{\partial \vec{p}} \frac{\partial}{\partial \vec{q}},$$

where  $(\vec{q}, \vec{p})$  are dynamical variables.

- Truncated-Power-Series-Algebra <sup>1</sup>

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<sup>1</sup>Differential Algebra package used in OPAL

## Computation of Non-Linear Space Charge Map

- **Ansatz:** Compute non-linear  $\mathcal{M}_{sc}$  starting from self-field potential of ellipsoid [Glu86]

$$\phi(x, y, z) = \frac{Q a_x a_y a_z}{4\epsilon_0} \int_0^\infty \frac{1}{\sqrt{(a_x^2 + u)(a_y^2 + u)(a_z^2 + u)}} \int_{S(u)}^\infty f(s) ds du,$$

where

$$S(u) = \frac{x^2}{a^2 + u} + \frac{y^2}{b^2 + u} + \frac{z^2}{c^2 + u},$$

with semi-principal axes  $a_x$ ,  $a_y$  and  $a_z$ .

# Computation of Non-Linear Space Charge Map

- Since we defined

$$H_{sc} = q\phi,$$

we get

$$\mathcal{M}_{sc} = e^{-s:H_{sc}:} = e^{-sq:\phi:} = 1 - sq:\phi: + \frac{(sq)^2}{2}:\phi:^2 - \frac{(sq)^3}{3!}:\phi:^3 + \dots,$$

where

$$:\phi: = \frac{\partial\phi}{\partial x} \frac{\partial}{\partial p_x} + \frac{\partial\phi}{\partial y} \frac{\partial}{\partial p_y} + \frac{\partial\phi}{\partial z} \frac{\partial}{\partial p_z} - \underbrace{\frac{\partial\phi}{\partial p_x} \frac{\partial}{\partial x}}_0 - \underbrace{\frac{\partial\phi}{\partial p_y} \frac{\partial}{\partial y}}_0 - \underbrace{\frac{\partial\phi}{\partial p_z} \frac{\partial}{\partial z}}_0$$

$$= \frac{\partial\phi}{\partial x} \frac{\partial}{\partial p_x} + \frac{\partial\phi}{\partial y} \frac{\partial}{\partial p_y} + \frac{\partial\phi}{\partial z} \frac{\partial}{\partial p_z}$$

$$:\phi: = -E_x \frac{\partial}{\partial p_x} - E_y \frac{\partial}{\partial p_y} - E_z \frac{\partial}{\partial p_z}$$

# Computation of Non-Linear Space Charge Map

## Expansion of Electric Field

- Since  $\vec{E}$  independent of momenta  $\implies \phi^{(j)} \equiv 0 \forall j > 1$ , thus

$$\mathcal{M}_{sc} = e^{-sq:\phi:} = 1 - sq:\phi: = 1 + sq \left( E_x \frac{\partial}{\partial p_x} + E_y \frac{\partial}{\partial p_y} + E_z \frac{\partial}{\partial p_z} \right)$$

**Idea:** Expanding electric field

$$E_w \approx \sum_{i=0}^n c_{w,i} w^i,$$

where  $w \in \{x, y, z\}$ .

# Computation of Non-Linear Space Charge Map

## Expansion of Electric Field

- General form of space charge map:

$$\mathcal{M}_{sc}(s) = \begin{pmatrix} 1 \\ 1 + sqE_x \\ 1 \\ 1 + sqE_y \\ 1 \\ 1 + sq\gamma^2 E_z \end{pmatrix} \approx \begin{pmatrix} 1 \\ 1 + sq \sum_{i=0}^n c_{x,i} x^i \\ 1 \\ 1 + sq \sum_{i=0}^n c_{y,i} y^i \\ 1 \\ 1 + sq\gamma^2 \sum_{i=0}^n c_{z,i} z^i \end{pmatrix},$$

with variable ordering  $(x, p_x, y, p_y, z, p_z)$ .

- **Question:** How do we get coefficients  $c_{x,i}$ ,  $c_{y,i}$ ,  $c_{z,i}$ ?  
 → **Least squares minimization (LSM)**

# Computation of Non-Linear Space Charge Map

## Expansion of Electric Field

- LSM yields to a linear system of equations:

$$\underbrace{\begin{pmatrix} \langle 1 \rangle & \langle w \rangle & \langle w^2 \rangle & \cdots & \langle w^n \rangle \\ \langle w \rangle & \langle w^2 \rangle & \langle w^3 \rangle & \cdots & \langle w^{n+1} \rangle \\ \langle w^2 \rangle & \langle w^3 \rangle & \ddots & & \vdots \\ \vdots & \vdots & & \langle w^{2n-2} \rangle & \langle w^{2n-1} \rangle \\ \langle w^n \rangle & \langle w^{n+1} \rangle & \cdots & \langle w^{2n-1} \rangle & \langle w^{2n} \rangle \end{pmatrix}}_{\text{symmetric, only monomials}} \cdot \begin{pmatrix} c_{w,0} \\ c_{w,1} \\ \vdots \\ c_{w,n-1} \\ c_{w,n} \end{pmatrix} = \underbrace{\begin{pmatrix} \langle E_w \rangle \\ \langle E_w w \rangle \\ \vdots \\ \langle E_w w^{n-1} \rangle \\ \langle E_w w^n \rangle \end{pmatrix}}_{\text{E-field moments}}$$

## Example: Linear Expansion of Electric Field

- **Assumption:**

- Centered beam, i.e.  $\langle x \rangle = \langle y \rangle = \langle z \rangle = 0$
- $E_w \approx c_{w,0} + c_{w,1}w$

- **Result:**

$$\begin{pmatrix} 1 & 0 \\ 0 & \langle w^2 \rangle \end{pmatrix} \cdot \begin{pmatrix} c_{w,0} \\ c_{w,1} \end{pmatrix} = \begin{pmatrix} 0 \\ \langle E_w w \rangle \end{pmatrix} \implies E_w \approx \frac{\langle w E_w \rangle}{\langle w^2 \rangle} w$$

where  $w \in \{x, y, z\}$  and moments  $\langle \cdot \rangle$ . Thus, we get

$$M_{sc}(s) = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ q \frac{\langle x E_x \rangle}{\langle x^2 \rangle} s & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & q \frac{\langle y E_y \rangle}{\langle y^2 \rangle} s & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & q \gamma^2 \frac{\langle z E_z \rangle}{\langle z^2 \rangle} s & 1 \end{pmatrix}.$$

## Linear Space Charge Map

- Space charge Hamiltonian

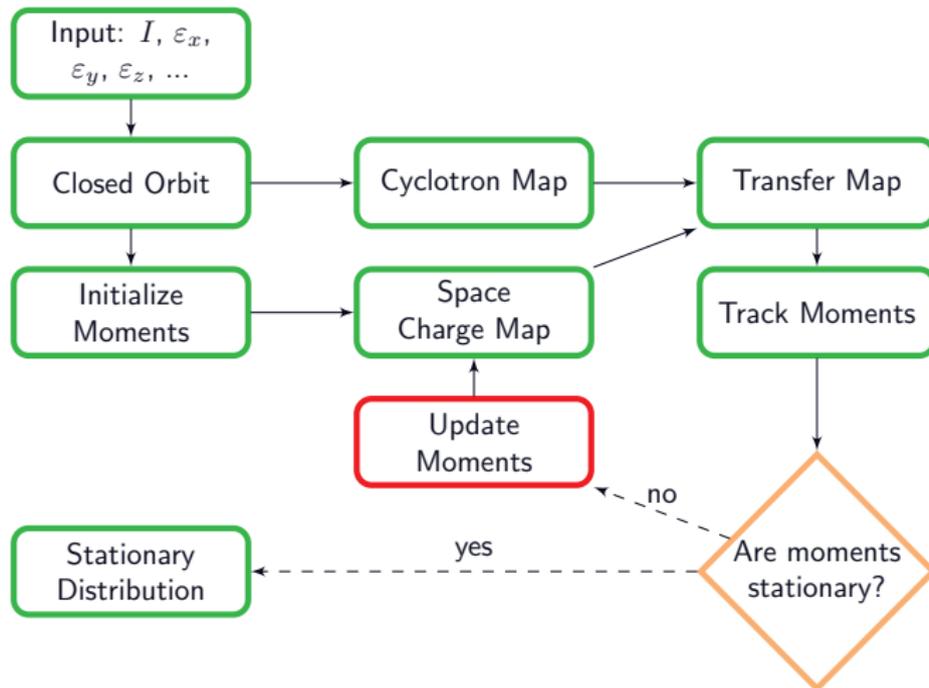
$$H_{sc} = -\frac{K_x}{2}x^2 - \frac{K_y}{2}y^2 - \frac{\gamma^2 K_z}{2}z^2,$$

with space charge strengths  $K_x$ ,  $K_y$  and  $K_z$ . This leads to

$$M_{sc}(s) = \begin{pmatrix} 1 & 0 & 0 & 0 & 0 & 0 \\ K_x s & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & K_y s & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 & K_z \gamma^2 s & 1 \end{pmatrix}.$$

# Updating Moments

## Context



## Updating Moments - Minimization Problem

### Goal:

$$\langle x^i p_x^j y^k p_y^l z^m p_z^n \rangle(s) \equiv \langle x^i p_x^j y^k p_y^l z^m p_z^n \rangle(s + L) \quad \forall s,$$

where  $L$  is the length of the orbit and  $\langle x^i p_x^j y^k p_y^l z^m p_z^n \rangle$  specifies a general moment.

- **Minimization Problem:**

$$\arg \min \frac{1}{2} \left\| \langle x^i p_x^j y^k p_y^l z^m p_z^n \rangle(s + L) - \langle x^i p_x^j y^k p_y^l z^m p_z^n \rangle(s) \right\|_2^2.$$

( $\arg \min f(x) \hat{=}$  find  $x$  for which  $f(x)$  is minimal)

- **Matrix form:**

$$\arg \min \frac{1}{2} \left\| A \vec{x}_{s+L} - \vec{x}_s \right\|_2^2.$$

# Updating Moments - Constraining the System

- **Matrix form:**

$$\arg \min \frac{1}{2} \|A\vec{x}_{s+L} - \vec{x}_s\|_2^2.$$

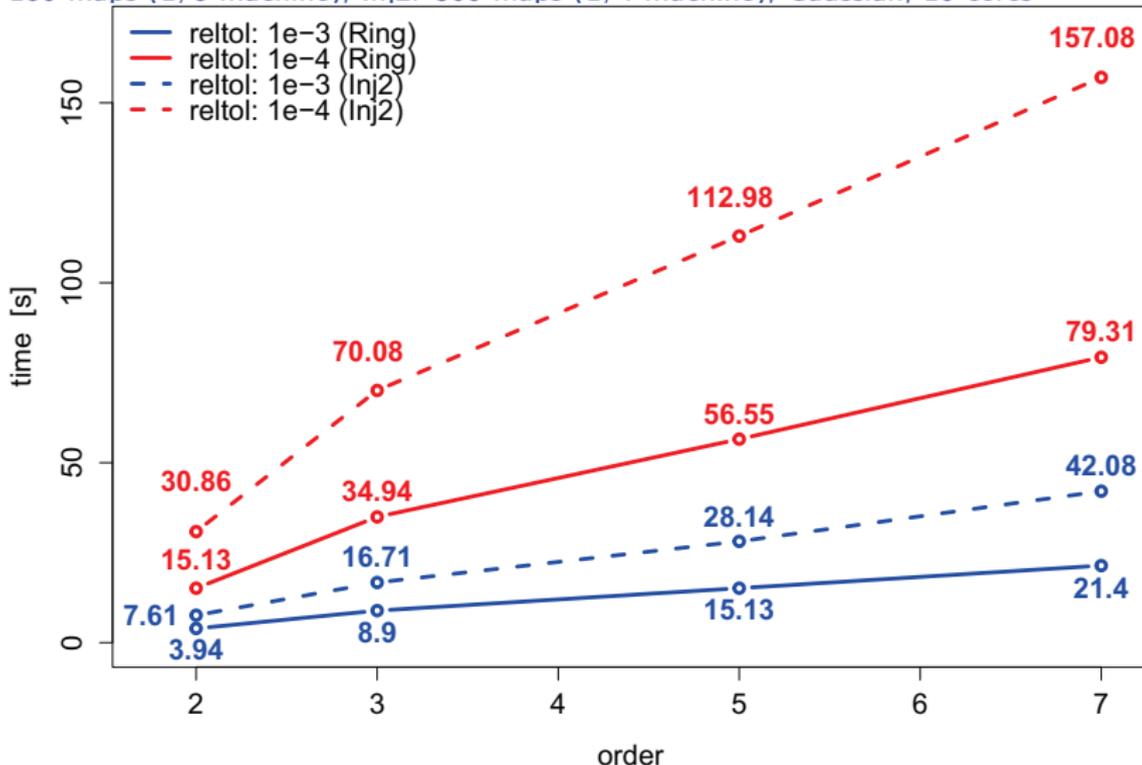
**Attention:** Matched distribution  $\rightarrow \vec{x}_{s+L} \equiv \vec{x}_s$

- **Questions:**

- How do we avoid non-physical results?  $\rightarrow$  Lagrange multipliers
- How do we incorporate emittances?
- How do we solve that problem?  
 $\rightarrow$  Fixed point computation, e.g. Newton-Raphson method

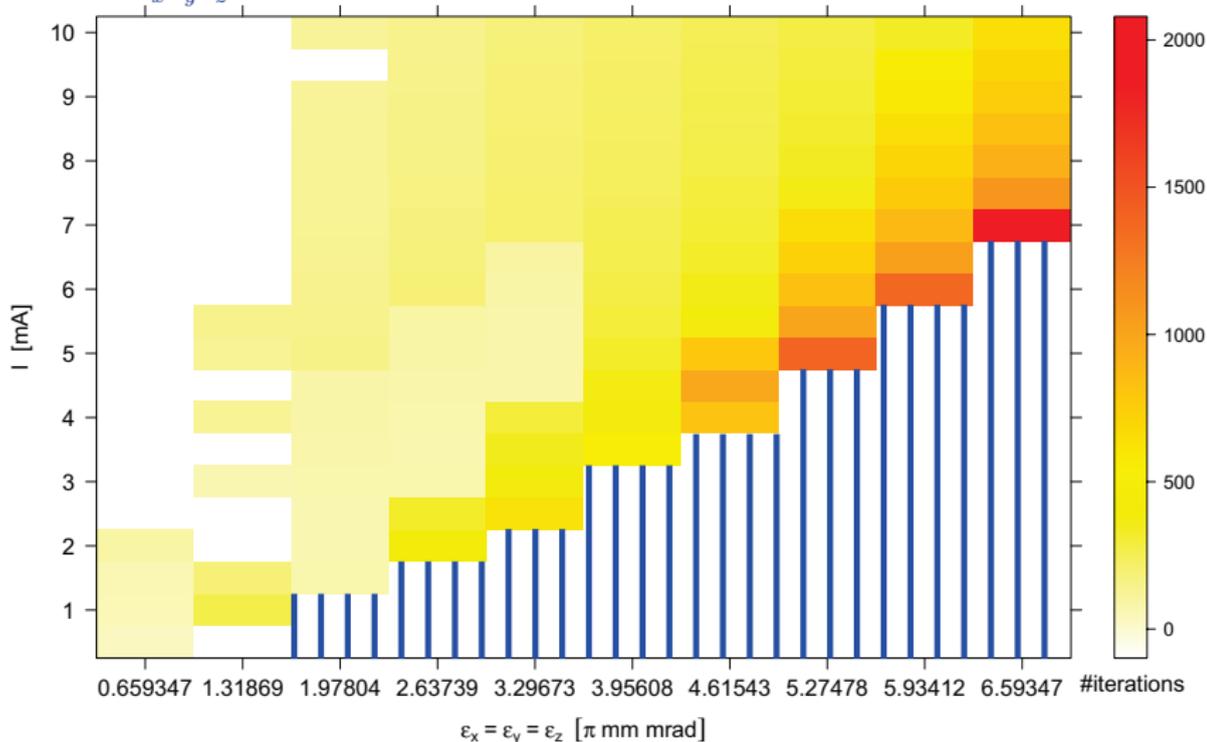
# Results: Space Charge Map Constructions $\mathcal{M}_{SC}$

Ring: 180 maps (1/8 machine), Inj2: 360 maps (1/4 machine); Gaussian, 16 cores



# Results: Linear Stability PSI Injector-2

Uniform ( $f = \frac{1}{8a_x a_y a_z}$ ),  $E = 2$  MeV



## Results: Non-Linear Matching

- **Goal:**

$$\sigma^{(3)}(s) = \mathcal{M} \circ \sigma^{(3)}(0) \circ \mathcal{M}^T \text{ and}$$

$$\sigma^{(3)} = \sigma^{(3)}(s) \stackrel{!}{=} \sigma^{(3)}(0)$$

- **Assumptions:**

- beam centered, i.e.  $\langle x \rangle = \langle y \rangle = \langle z \rangle = 0$
- symmetric,  $\langle E_x \rangle = \langle E_y \rangle = \langle E_z \rangle = 0$
- moments of order  $> 3$  equal zero

- **Arbitrary distribution perturbation:**

$$f(x, y, z) = f_{Gaussian}(x, y, z) \cdot \exp \left[ -k \langle x^3 \rangle \frac{x^3}{\langle x^2 \rangle^{3/2}} \right]$$

with constant  $k = 0.1$ .

- Initially 3rd order moments uniformly in  $[-0.01, 0.01]$

# Non-Linear Matching

- $L_2$ -error:  $9.8451 \cdot 10^{-4}$

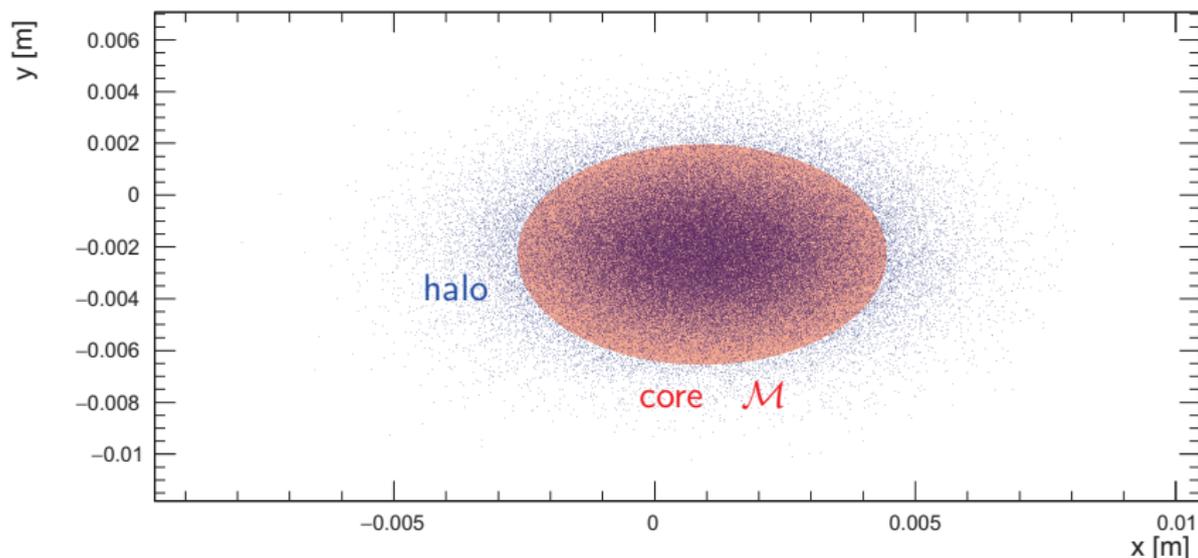
$\langle x^3 \rangle = -0.1946,$	$\langle x^2 p_x \rangle = -7.5797 \cdot 10^{-5},$	$\langle x^2 y \rangle = 0.0954,$	$\langle x^2 p_y \rangle = -0.0136,$
$\langle x^2 z \rangle = -0.0786,$	$\langle x^2 p_z \rangle = -0.0393,$	$\langle x p_x^2 \rangle = -0.0131,$	$\langle x p_x y \rangle = 0.0042,$
$\langle x p_x p_y \rangle = 0.0095,$	$\langle x p_x z \rangle = -0.0231,$	$\langle x p_x p_z \rangle = 5.1948 \cdot 10^{-5},$	$\langle x y^2 \rangle = 0.0070,$
$\langle x y p_y \rangle = 0.0012,$	$\langle x y z \rangle = 0.0399,$	$\langle x y p_z \rangle = 0.0263,$	$\langle x p_y^2 \rangle = 0.0006,$
$\langle x p_y z \rangle = 0.0108,$	$\langle x p_y p_z \rangle = -0.0035,$	$\langle x z^2 \rangle = -0.015,$	$\langle x z p_z \rangle = -0.0285,$
$\langle x p_x^2 \rangle = -0.0029,$	$\langle p_x^3 \rangle = 0.0004,$	$\langle p_x^2 y \rangle = 0.0004,$	$\langle p_x^2 p_y \rangle = -0.0007,$
$\langle p_x^2 z \rangle = -0.0051,$	$\langle p_x^2 p_z \rangle = -0.0072,$	$\langle p_x y^2 \rangle = 0.0007,$	$\langle p_x y p_y \rangle = 0.0004,$
$\langle p_x y z \rangle = 0.0061,$	$\langle p_x y p_z \rangle = 0.0023,$	$\langle p_x p_y^2 \rangle = 0.0004,$	$\langle p_x p_y z \rangle = 0.0049,$
$\langle p_x p_y p_z \rangle = 0.0052,$	$\langle p_x z^2 \rangle = -0.020,$	$\langle p_x z p_z \rangle = -0.0124,$	$\langle p_x p_z^2 \rangle = 0.0003,$
$\langle y^3 \rangle = -0.0017,$	$\langle y^2 p_x \rangle = -0.0015,$	$\langle y^2 z \rangle = -0.0310,$	$\langle y^2 p_z \rangle = 0.0033,$
$\langle y p_y^2 \rangle = -0.0001,$	$\langle y p_y z \rangle = 0.0009,$	$\langle y p_y p_z \rangle = -1.1216 \cdot 10^{-5},$	$\langle y z^2 \rangle = 0.0185,$
$\langle y z p_z \rangle = 0.0055,$	$\langle y p_z \rangle = 0.0009,$	$\langle p_y^3 \rangle = -0.0006,$	$\langle p_y^2 z \rangle = -0.0054,$
$\langle p_y^2 p_z \rangle = 0.0007,$	$\langle p_y z^2 \rangle = 0.0185,$	$\langle p_y z p_z \rangle = 0.0086,$	$\langle p_y p_z^2 \rangle = -0.0009,$
$\langle z^3 \rangle = -0.1298,$	$\langle z^2 p_x \rangle = -0.0073,$	$\langle z p_x^2 \rangle = -0.0154,$	$\langle p_z^3 \rangle = -0.0011$

Converged!

# Conclusion

- **Conclusions:**
  - TPSA based
  - Good agreement of linear theory
  - Lagrange method allows extension of constraints
  - Non-linear matching is working
- **Open Questions:**
  - What is the physical relevance of non-linear matching?
  - Linear Stability behaviour? Why do we observe gaps?
  - System constraints? Higher order invariants?
  - Relative error at peak of max 15 % in the field expansion for 9th order (spherical Gaussian).
- **Improvements of the presented method:**
  - Convergence rate
  - Computation time (parallelisation?)
  - Initialisation of moments
- **Take home points:**
  - Non-linear matching
  - Non-linear maps

## Particle Non-Linear Particle Core Model



- Particles outside feel non-linear forces in existing particle core models
- Higher order maps for the core  $\rightarrow$  particles passing the core feel non-linear forces

## References

- [Fre15] M. Frey. Master's thesis, PSI and ETH Zurich, 2015.  
[http://amas.web.psi.ch/people/aadelmann/  
ETH-Accel-Lecture-1/projectscompleted/cse/  
MScThesis\\_Frey.pdf](http://amas.web.psi.ch/people/aadelmann/ETH-Accel-Lecture-1/projectscompleted/cse/MScThesis_Frey.pdf).
- [Glu86] R. L. Gluckstern. Scalar potential for charge distributions with ellipsoidal symmetry. Fermilab Internal Report TM-1402, May 1986.