

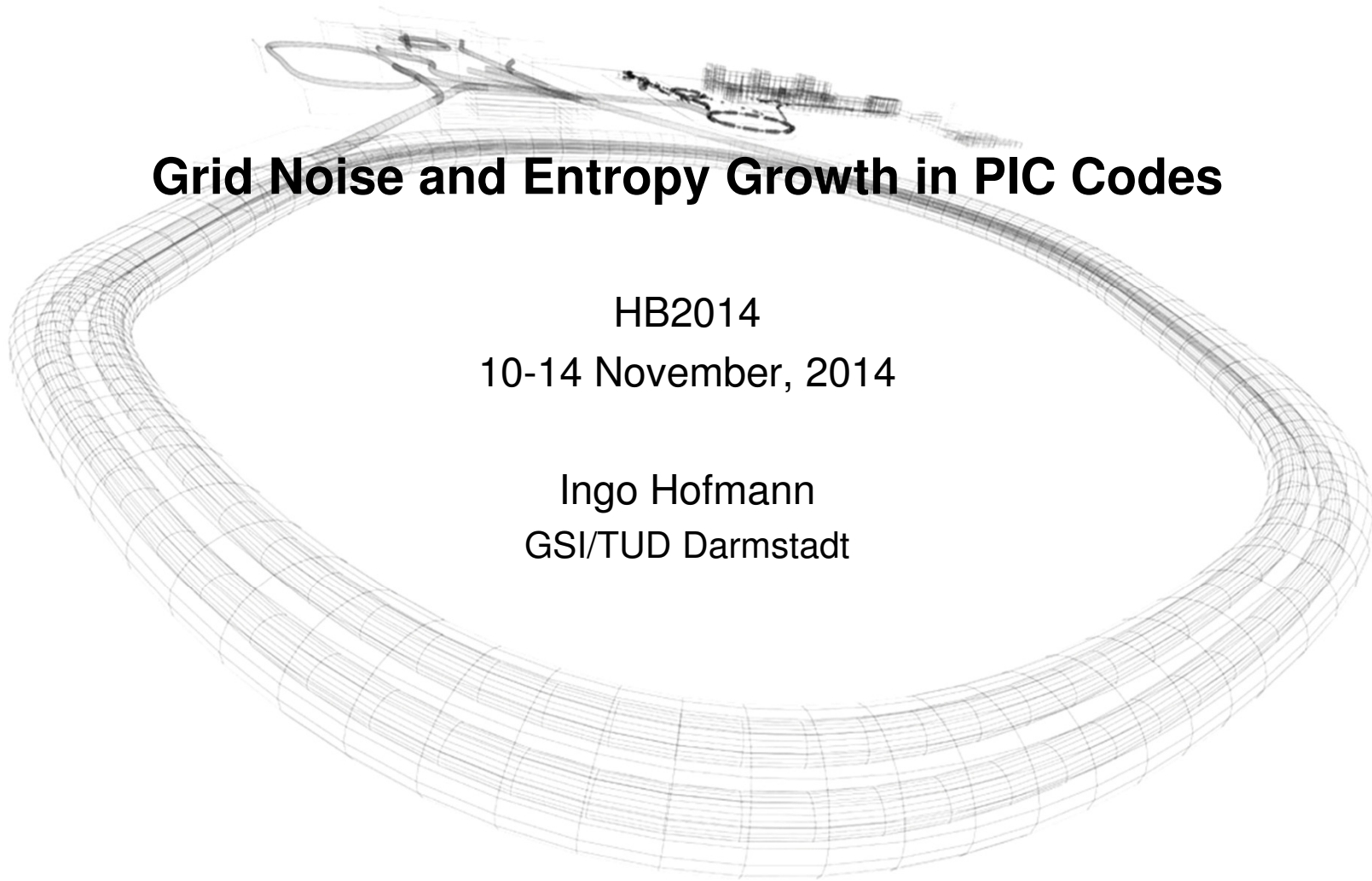


# Grid Noise and Entropy Growth in PIC Codes

HB2014

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- Theory overview
- Grid and particle collisions in 3d and 2d (bunches + FODO)
- Noise and rms entropy in equilibrium
- Application to resonant/unstable processes
- Conclusions

*Acknowledgments:*

*O. Boine-Frankenheim (co-worker)*

*J. Struckmeier (discussions)*

# Motivation

*Why do we study collisions and noise in high intensity simulation?*

- „usual“ approach: let us increase N, grid resolution, time steps → convergence
- in linacs so far we have lived well with this – questions arise for error studies (1000's of linacs) with „low“ N ( $\sim 10^5$ )
- in circular machines 100.000's of turns – cross-check resonance trapping with space charge (noise-free simulations by G. Franchetti)
- CERN space charge workshop April 2013 triggered off new interest ..

*3 papers give basis for „equilibrium“ beams:*

1. *Theory: J. Struckmeier, Part. Acc. (1994) & Phys. Rev. E (1996) & PRSTAB (2000)*
2. *3d-simulation: I. Hofmann and O. Boine-Frankenheim, to be published in PRSTAB (2014)*
3. *2d-simulation: O. Boine-Frankenheim et al. NIM (2014)*

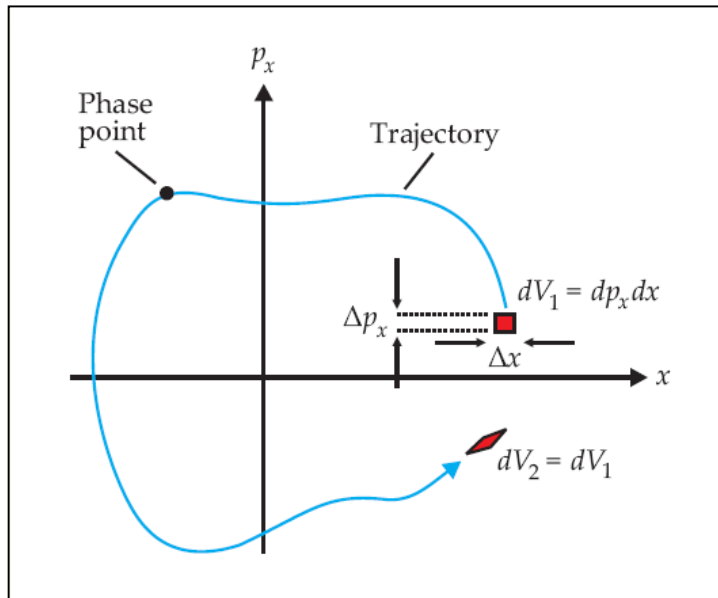
Current study extends equilibrium discussion to (typical) dynamical phenomena  
(mismatch, resonances ...)

- are above noise concepts useful?

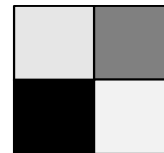
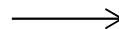
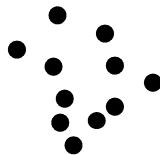
# Entropy

## Liouville - infinite resolution – coarse graining

### Liouville



- Liouville: exact areas in 2d, 4d, 6d phase space invariant for Hamiltonian flow – no growth of „infinite resolution“ entropy
- How break this?
  - for „finite resolution“ self-consistent space charge potential on a grid (as in PIC) exactly Hamiltonian flow broken → growth of entropy
  - for exactly resolved motion of directly interacting particles (collisions) in 6d:
    - collisions break Hamiltonian flow in 6d → growth of entropy
    - „infinite resolution“ entropy in 6Nd phase space still invariant → Gibbs introduced **coarse graining of phase space to obtain entropy growth**



$$S_{CG}(\rho(X)) = -k \int \bar{\rho}(X) [\ln \bar{\rho}(X)] d\Gamma$$

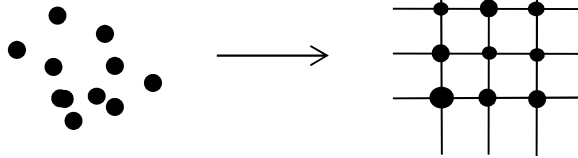
$\bar{\rho}$  is averaged over a cell

# Poisson solver grid & noise

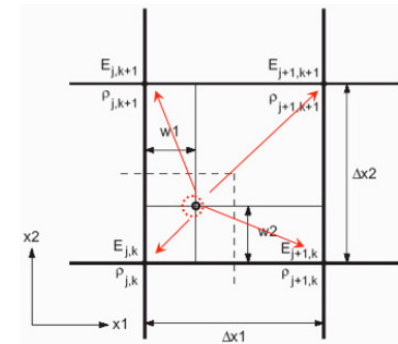
## Grid induced noise:

- not included in original „collisional“ approach of Struckmeier
- assume several sources:
  - **non-Liouvillian effect by fluctuating charges on grid**
  - focussing modulation „fast“
  - coherent flow vs. incoherent „temperature“

**PIC:** replaces infinite resolution charge density by “grid distribution”



area weighting



Phase space remains infinitely resolved

## Averaging charges over cells:

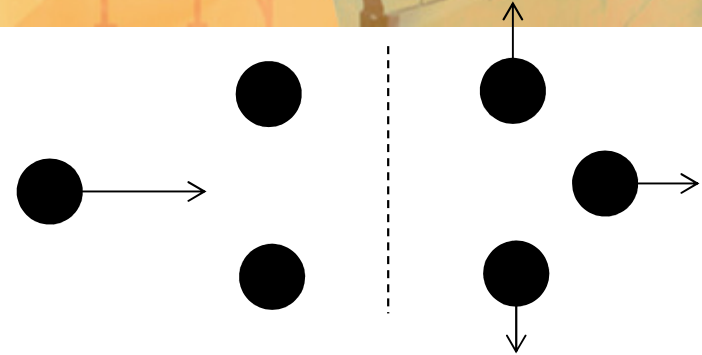
- ✓ loss of Hamiltonian character of phase space flow
- ✓ → growth of entropy

Is there a practically useful „reduced“ definition of entropy?



# “Reduced” approach for beams: J. Struckmeier suggested (1994 ff) to start from Fokker-Planck Equation in 6d phase space

- Assuming Markov process for collisions of charged particles: dynamical friction
  - no correlation between several particles
  - memory extinction
- friction and diffusion terms



collisional interaction  
"beyond" Hamiltonian  
flow in 6d

$$\frac{df}{dt} = \frac{\partial f}{\partial t} + \sum_{i=1}^3 \left( \dot{x}_i \frac{\partial f}{\partial x_i} + \dot{p}_i \frac{\partial f}{\partial p_i} \right) = \left[ \frac{\delta f}{\delta t} \right]_{Non-Liou}$$

$$\left[ \frac{\delta f}{\delta t} \right]_{Non-Liou} \equiv \sum_{i=1}^3 \frac{\partial}{\partial p_i} \{ \beta_{f;i} p_i f \} + \sum_{i,j=1}^3 \frac{m^2 \partial^2}{\partial p_i \partial p_j} \{ D_{ij}(\vec{p}, t) f \}$$

form second order  
moments

# Second order moments of Vlasov – FP equation

→ rms based (emittance, "temperature", ...)

$$\frac{d}{dt} \langle x_i^2 \rangle = \int x_i^2 \frac{\partial f}{\partial t} d\tau$$

$$\varepsilon_i^2(s) \equiv \langle x_i^2 \rangle \langle x_i'^2 \rangle - \langle x_i x_i' \rangle^2 \quad s = \beta ct$$

negligible transient effects associated with initially non-matched self-fields  
(I. Hofmann and J. Struckmeier, Part. Acc. 1987)

$$\frac{1}{\langle x_i^2 \rangle} \frac{d}{ds} \varepsilon_i^2(s) = \frac{2q}{mc^2 \beta^2 \gamma^3} \left( \langle x_i' F_i \rangle - \frac{\langle x_i x_i' \rangle}{\langle x_i^2 \rangle} \langle x_i E_i \rangle \right) - 2 \left( \frac{\beta_{f;i} \varepsilon_i^2(s)}{c \beta \gamma \langle x_i^2 \rangle} - \frac{\langle D_{ii} \rangle}{c^2 \beta^3 \gamma^3} \right)$$

$$\frac{kT_i}{mc^2 \beta^2 \gamma} = \frac{\varepsilon_i^2(s)}{\langle x_i^2 \rangle}; \quad T_{eq} = \frac{1}{3} \sum_{i=1}^3 T_i$$

$$D \equiv \langle D_{ii} \rangle; \quad \beta_f \equiv \beta_{f;i}; \quad k_f \equiv \beta_f / c \beta \gamma \quad D = \beta_f \gamma kT / m \text{ (Einstein relation)}$$

$$\frac{d}{ds} \ln \varepsilon_x^2 = k_f \left[ \frac{T_y(s)}{T_x(s)} - 1 \right] \quad \frac{d}{ds} \ln \varepsilon_y^2 = k_f \left[ \frac{T_x(s)}{T_y(s)} - 1 \right] \quad (\text{one decreasing one increasing})$$

$$\frac{d}{ds} \ln \varepsilon_x^2 \varepsilon_y^2 = k_f \left[ \frac{(T_x - T_y)^2}{T_x T_y} \right] \geq 0 \quad \text{or in 6d} \quad \frac{d}{ds} \ln \varepsilon_x^2 \varepsilon_y^2 \varepsilon_z^2 = \frac{2}{3} k_f \left[ \frac{(T_x - T_y)^2}{T_x T_y} + \frac{(T_x - T_z)^2}{T_x T_z} + \frac{(T_z - T_y)^2}{T_z T_y} \right] \geq 0$$

## Resulting in an rms **entropy law**

$$S_{6d,rms} \equiv \ln \varepsilon_x \varepsilon_y \varepsilon_z \quad \varepsilon_{6d,rms} \equiv \varepsilon_x \varepsilon_y \varepsilon_z$$
$$\frac{d}{ds} S_{6d,rms} = \frac{1}{3} k_f \left[ \frac{(T_x - T_y)^2}{T_x T_y} + \frac{(T_x - T_z)^2}{T_x T_z} + \frac{(T_y - T_z)^2}{T_y T_z} \right] \geq 0$$
$$= 0, \text{ if } T_x = T_y = T_z \text{ (equipartitioned)}$$

never decreasing – under the assumptions of its validity!

heuristic justification:

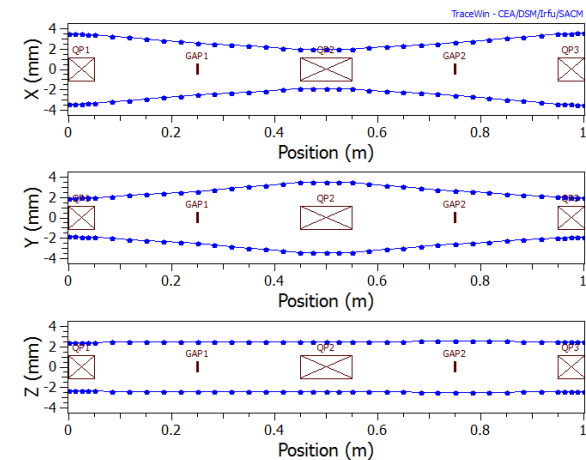
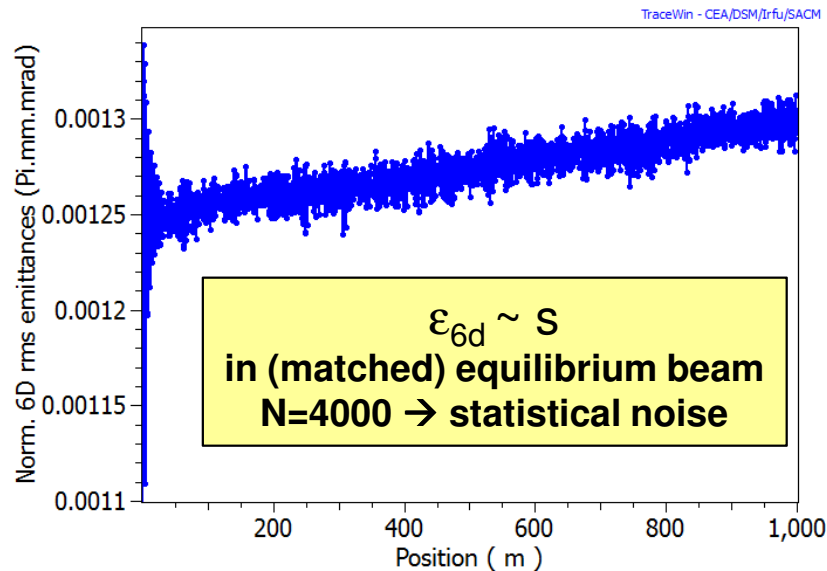
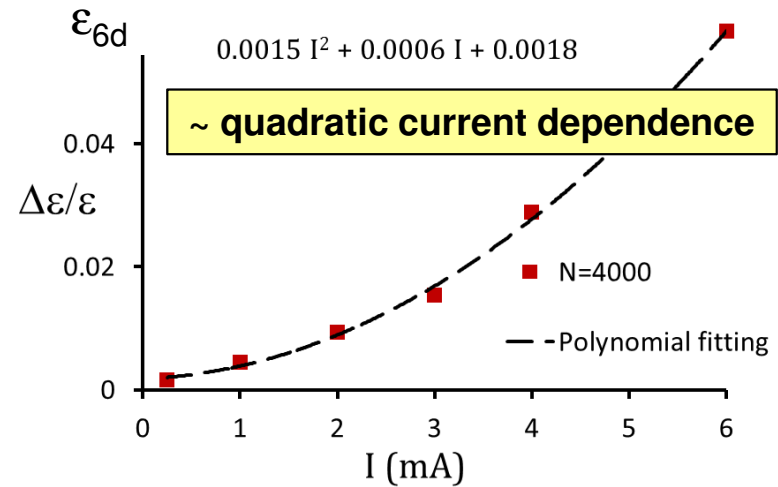
- emittances ~ "probabilities" in phase space
- → product of probabilities for "independent" events
- entropies should be additive for ~ decoupled probabilities (→ ln)
- → coupling between degrees of freedom only on microscopic level
- should be non-decreasing
- **rms approximation – no higher order correlations resolved !!!**



# TRACEWIN simulations: matched equilibrium

„Confirms“  $\epsilon_{6d}$  – as rms entropy concept

- spherical bunch:
- $k_{0x,y,z}=60^\circ$  and  $\epsilon_z=\epsilon_{xy}$
- $k_{x,y,z}=32^\circ/32^\circ/32^\circ$  (xyz Poisson solver)
- 1000 cells as reference base
- no acceleration – only RF gaps



# $\Delta\varepsilon_{6d}/\varepsilon_{6d}$ growth as function of # grid cells $n_c$

Minimizing rms entropy growth  $\rightarrow$  determine optimum grid/particle resolution

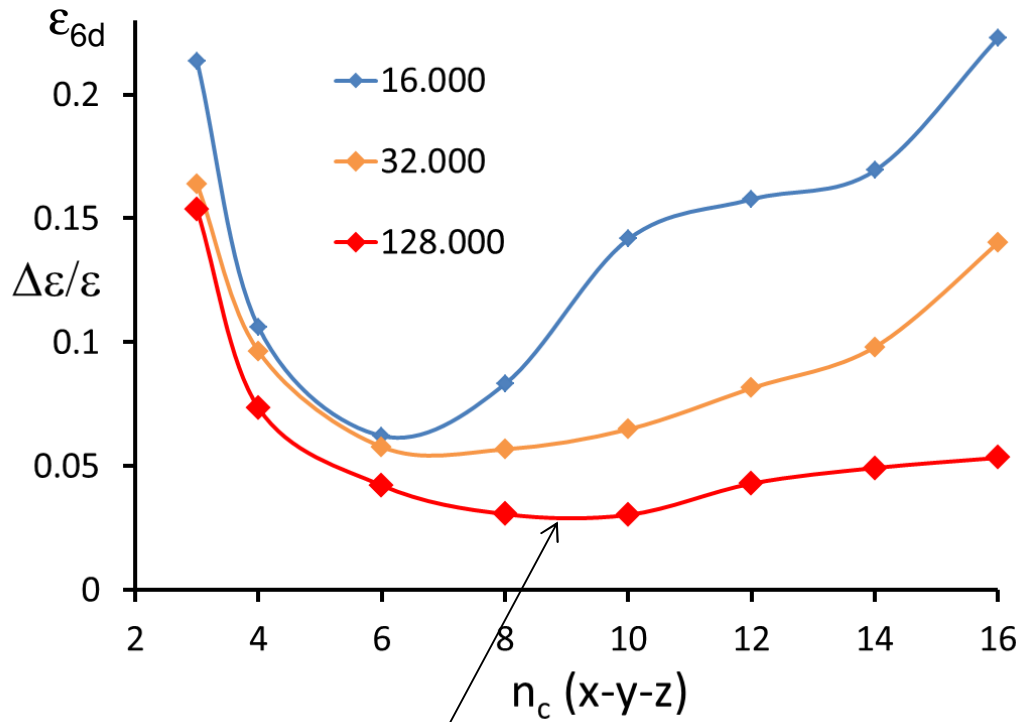
Fully isotropic (spherical) bunch

grid dominated

- not resolving Gaussian profile -

collision dominated

- charge  $\sim 1/N$  -



1 particle / cell

optimum

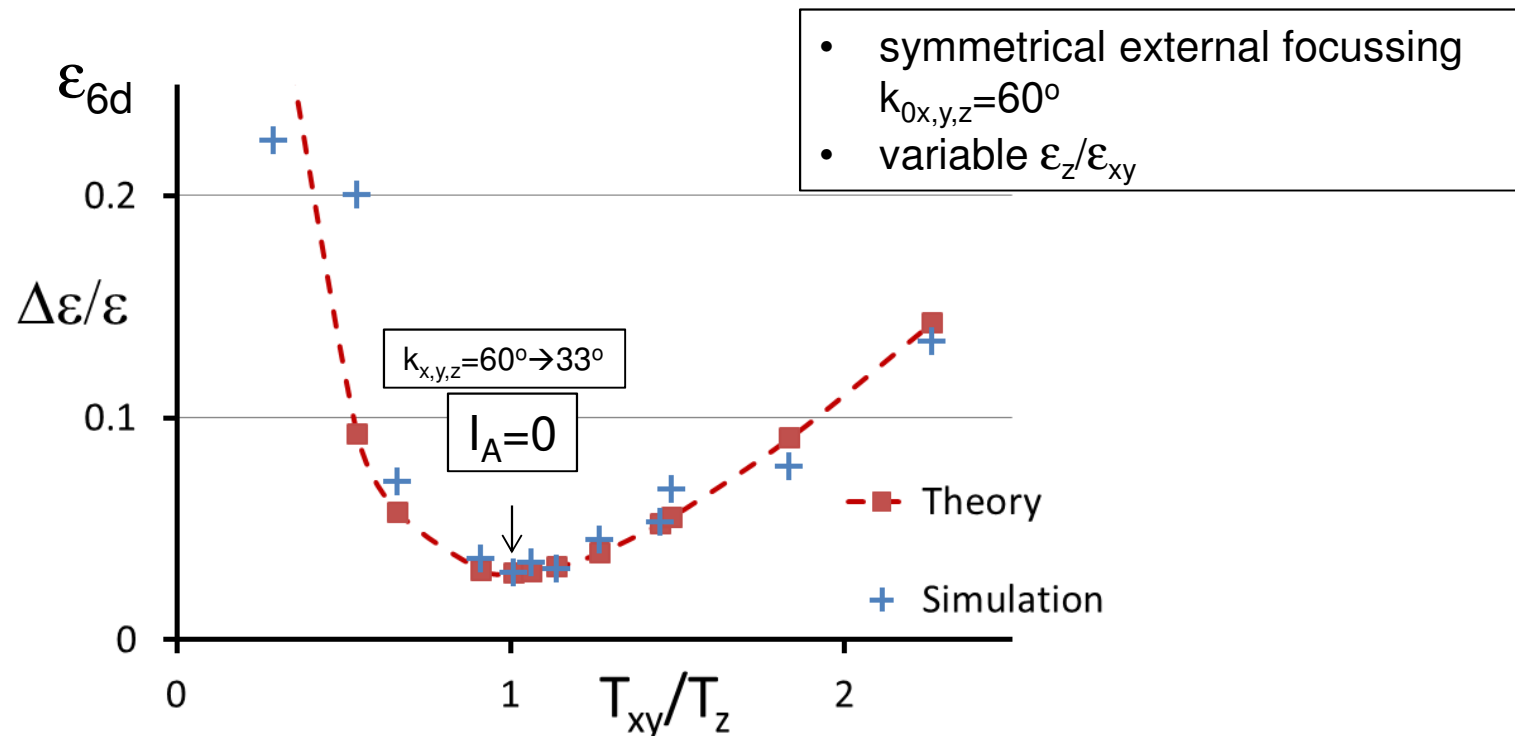
$n_c = 8$ :

- 16 cells within  $\pm 3.5 \sigma$
- outside of  $3.5 \sigma$  boundary analytical approximation



# Anisotropy and grid effect

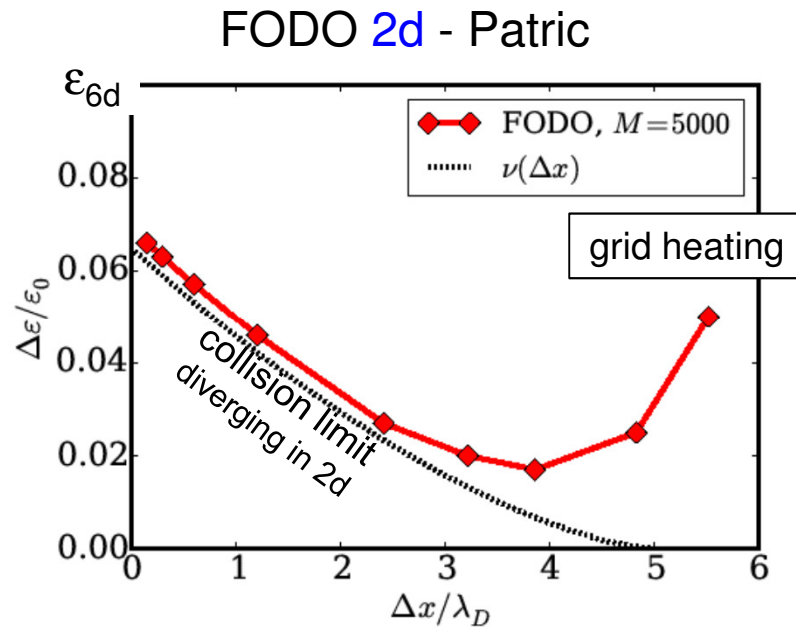
→ fits to theory with  $I_{GN}$  and  $k_f^*$  fitted to data



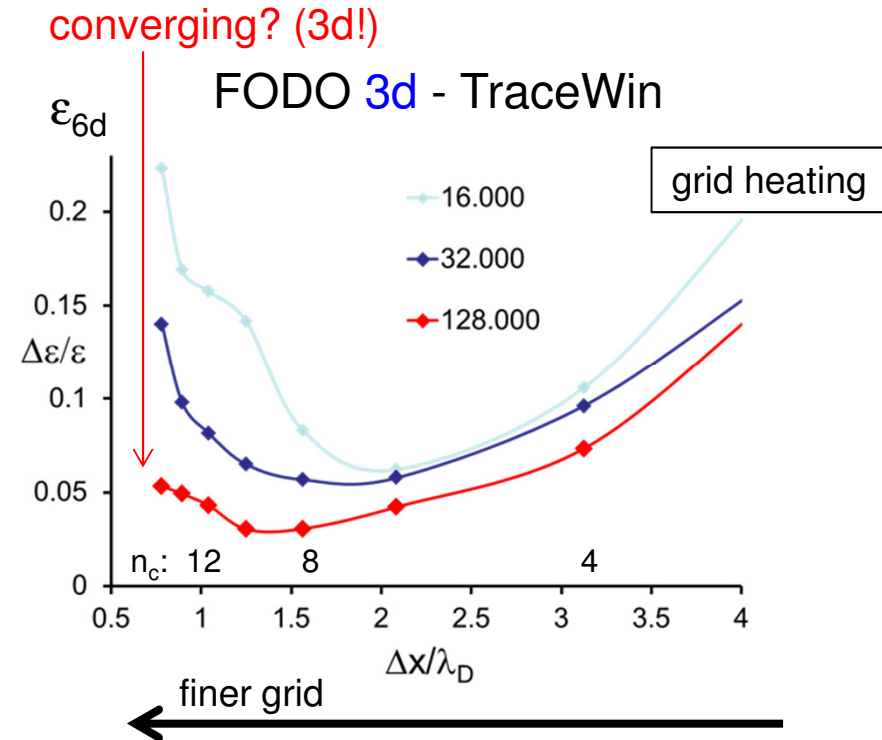
$$\frac{k_f}{3} \left( \frac{(1-r_{xy})^2}{r_{xy}} + \frac{(1-r_{xz})^2}{r_{xz}} + \frac{(1-r_{yz})^2}{r_{yz}} \right) \geq 0 \quad \longrightarrow \quad \frac{d}{ds} \ln \epsilon_x(s)\epsilon_y(s)\epsilon_z(s) = \frac{k_f^*}{3} (I_A + I_{GN})$$

$$\frac{1}{k} \Delta S = \frac{\Delta \epsilon_{6d}}{\epsilon_{6d}} = \Delta S \frac{k_f^*}{3} (I_A + I_{GN})$$

# Comparison of 3d with 2d results



Source: O. Boine-Frankenheim et al., NIM 2014



A. B. Langdon, *Effect of the spatial grid in simulation plasmas*, *J. Comput. Phys.* (1970)

→ Artificial heating if grid is too coarse:  $\Delta x \gtrsim \lambda_D$  (Debye length)

In 2d enhanced collisions if grid too fine → larger N

# From matched equilibrium beams

→ mismatched, resonances, ....

1. Can we apply our models of noise vs.  $N$  and/or  $n_c$  to dynamical situations with resonant effects? Use for optimization?
2. Are there transient/resonant/unstable phenomena, which require a more refined measure for noise/entropy than  $\epsilon_{6d,rms}$ ?
3. Do we need more theory efforts, or are phenomenological studies the only way?



# Non-stationary (non-equilibrium) beams

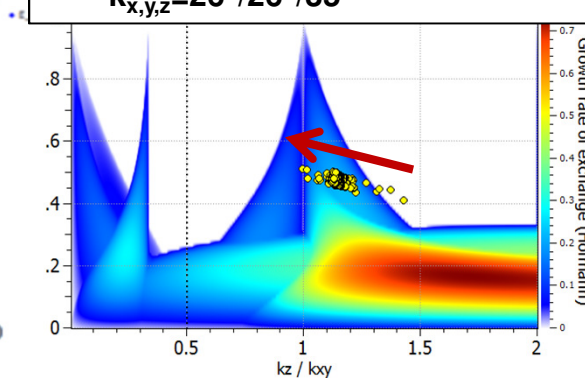
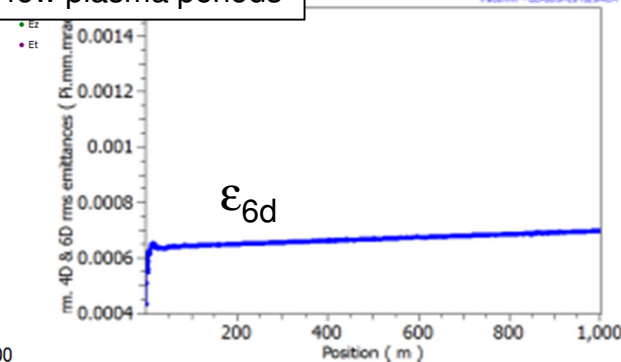
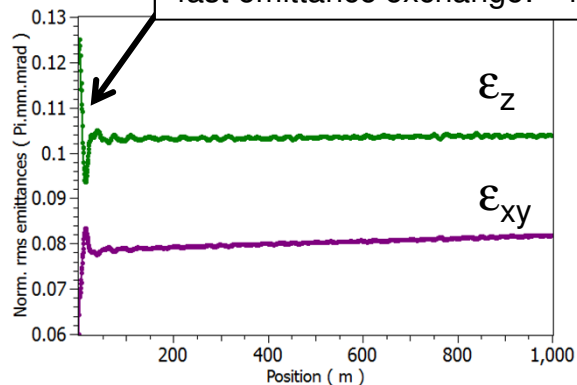
- fast emittance exchange on  $2k_z-2k_{xy}=0$  resonance (space charge octupole)

exchange practically unaffected by noise

$N=10^5$

fast emittance exchange: ~ few plasma periods

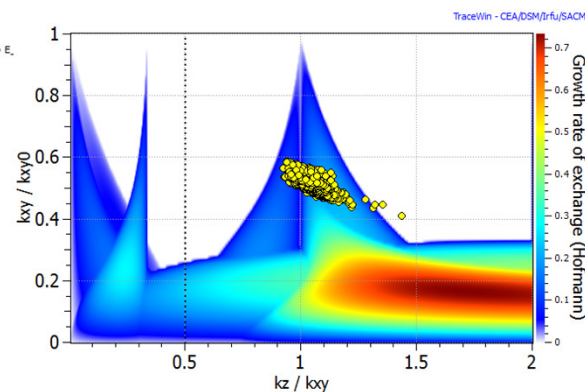
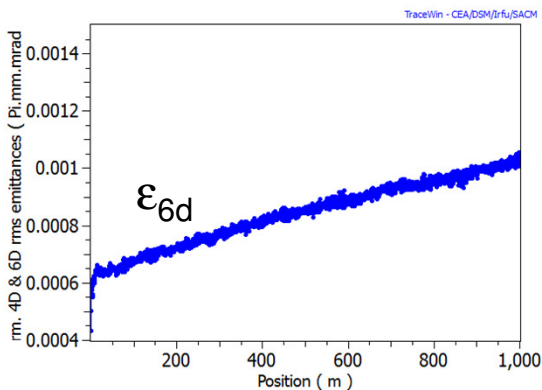
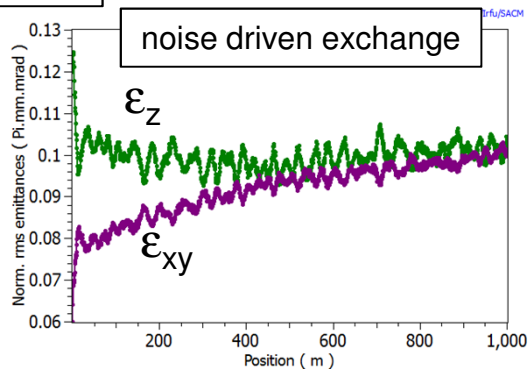
- $k_{0x,y,z}=60/60/60^\circ$  and  $\epsilon_z/\epsilon_{xy}=2$
- $k_{x,y,z}=26^\circ/26^\circ/35^\circ$



$N=10^3$

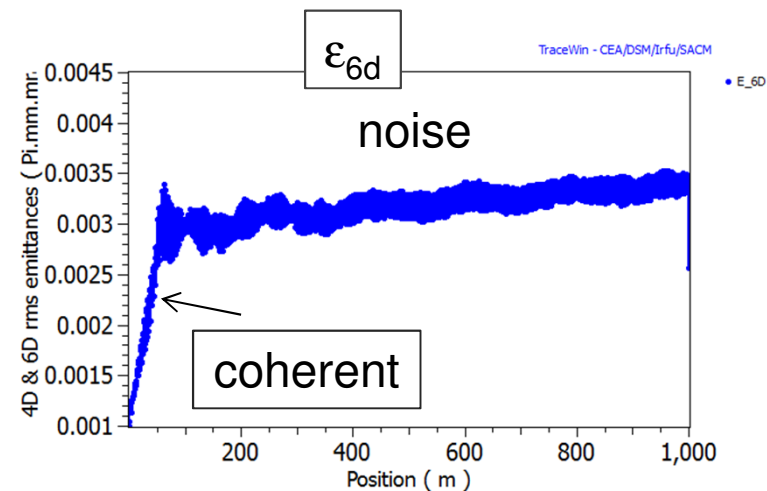
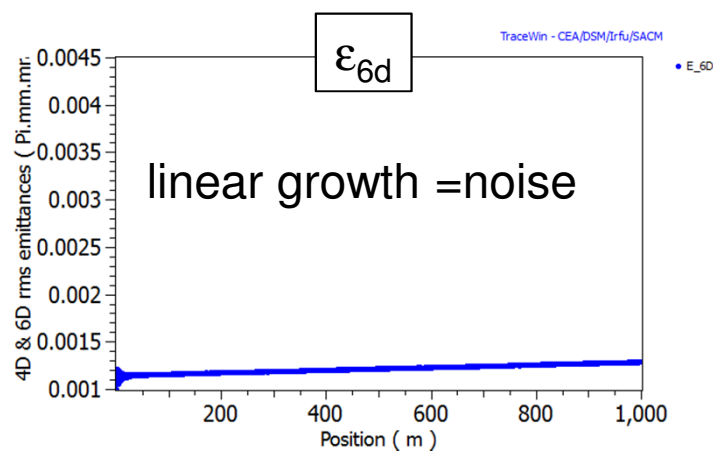
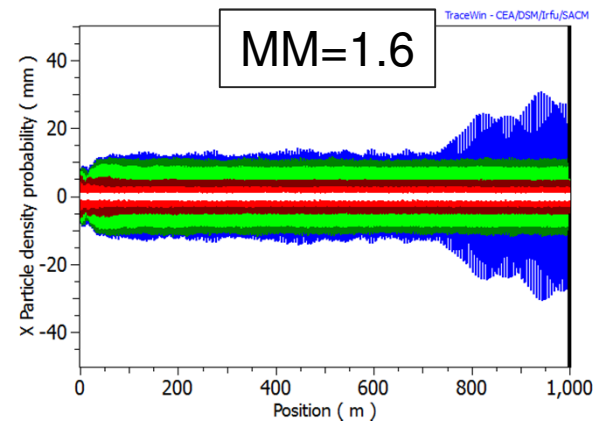
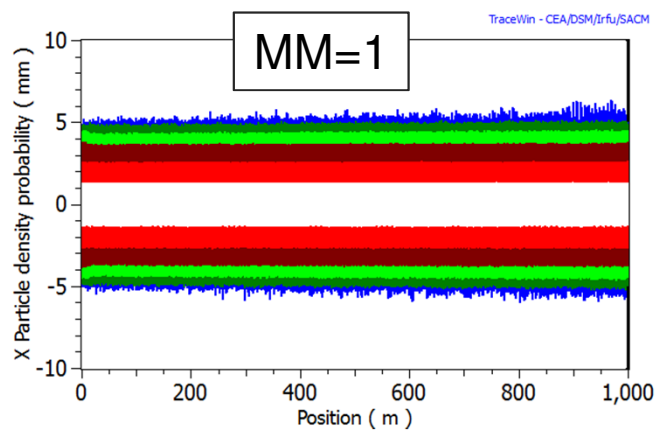
noise driven exchange

$n_c=6$



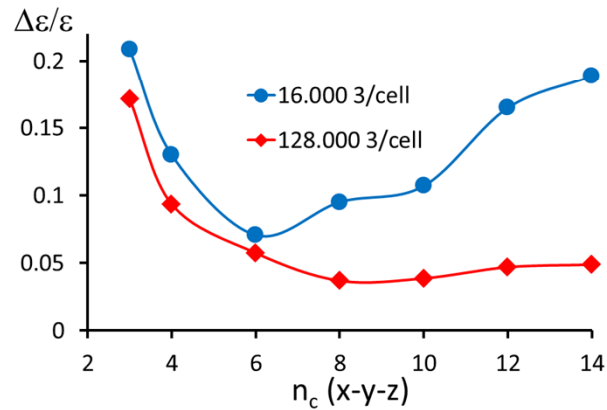
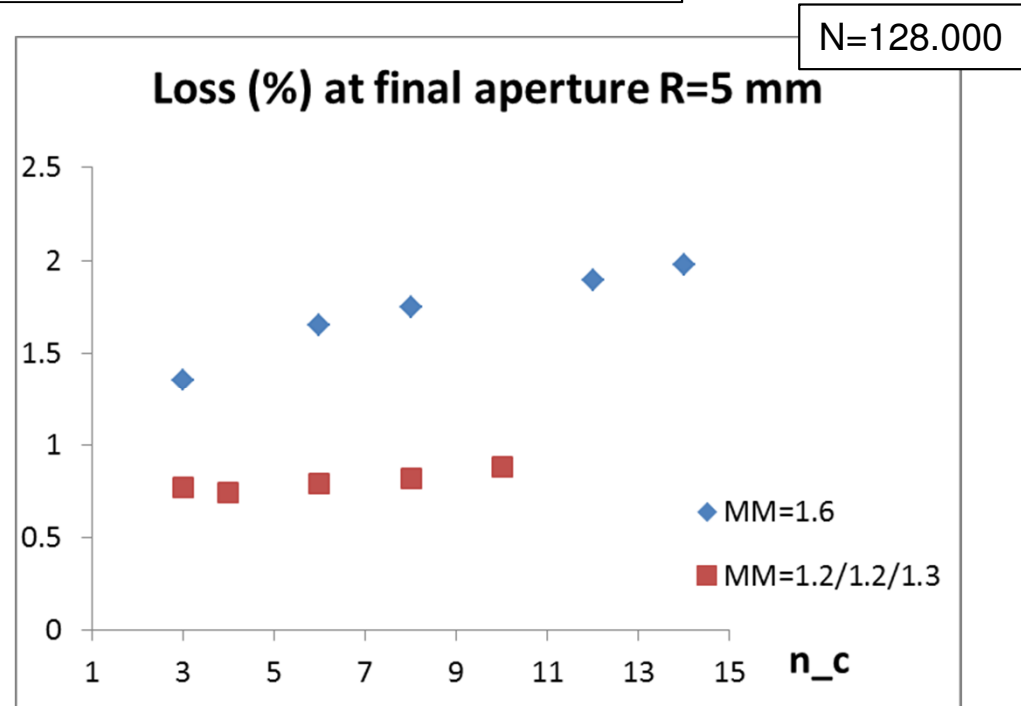
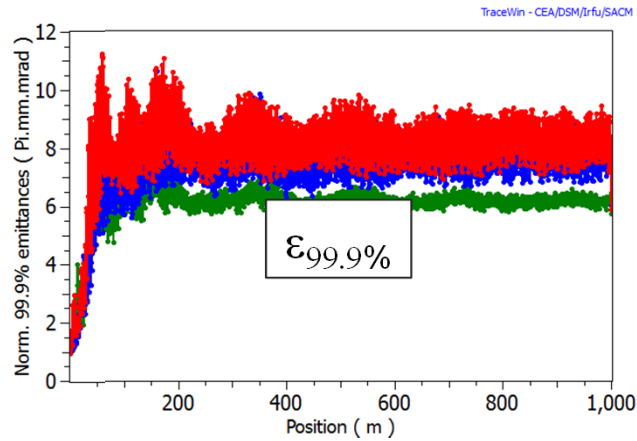
# Mismatched beams $\rightarrow$ halo formation

chose large MM factor in x,y and z &  $N=10^5$



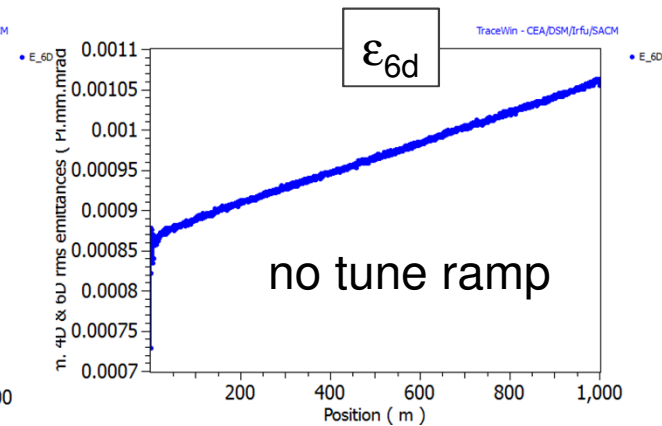
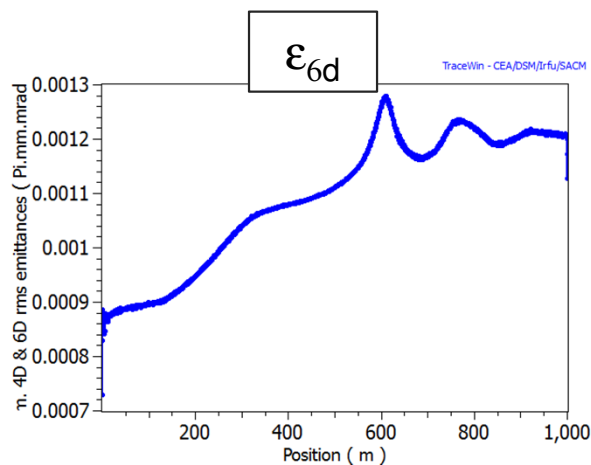
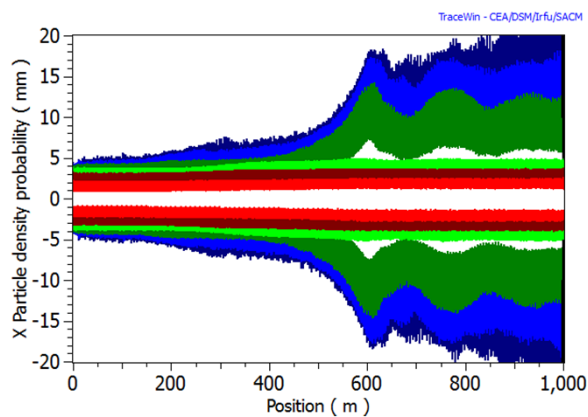
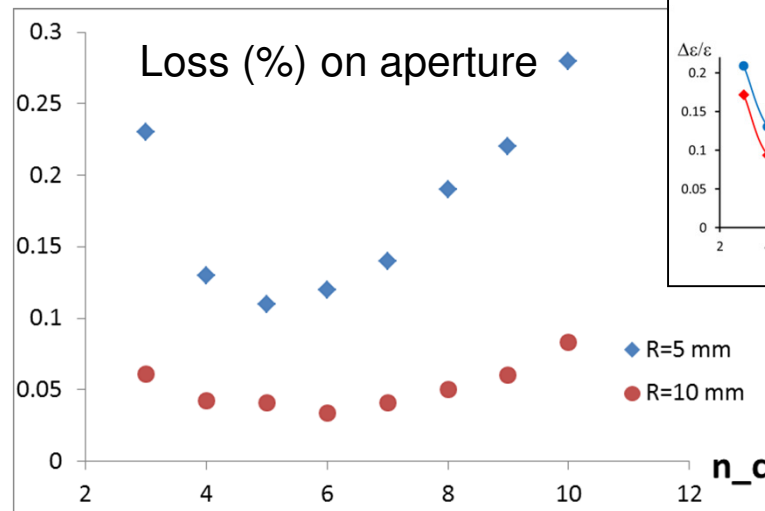
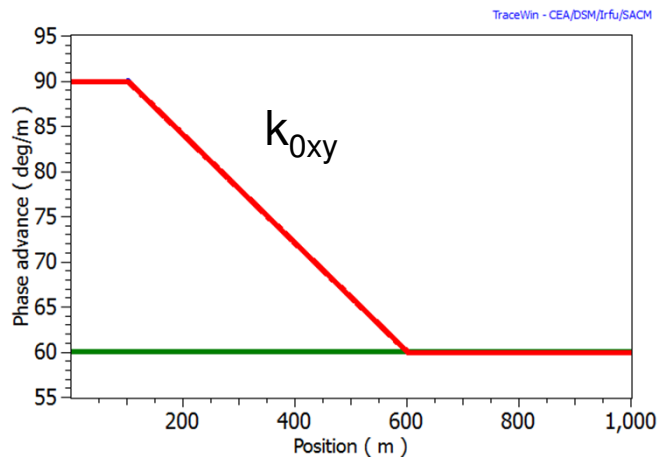
# Mismatch $\rightarrow$ halo $\sim$ 50 turns

fast conversion into halo  $\rightarrow$  noise has no effect on this!



# Crossing of space charge driven resonances

$k_{0xy}=90^{\circ} \rightarrow 60^{\circ}$  / 500 cells

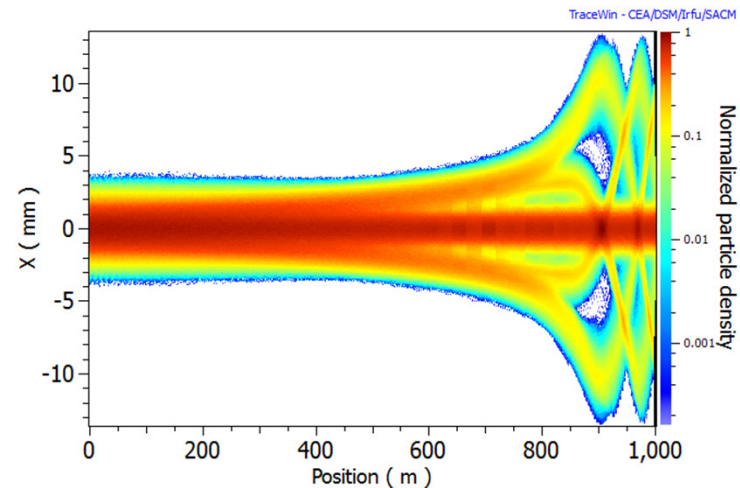
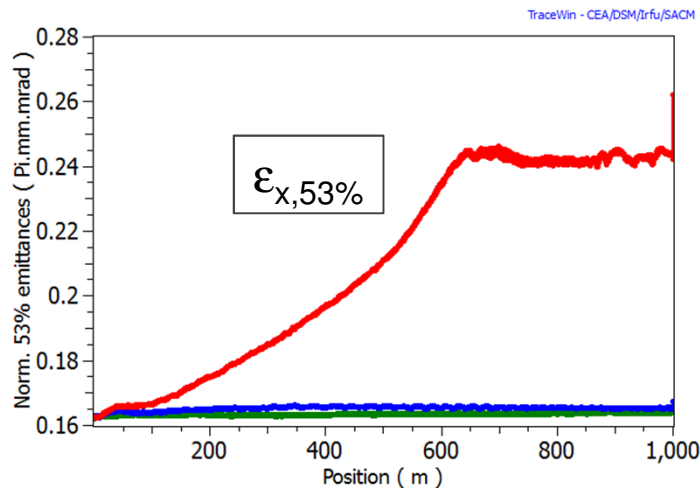
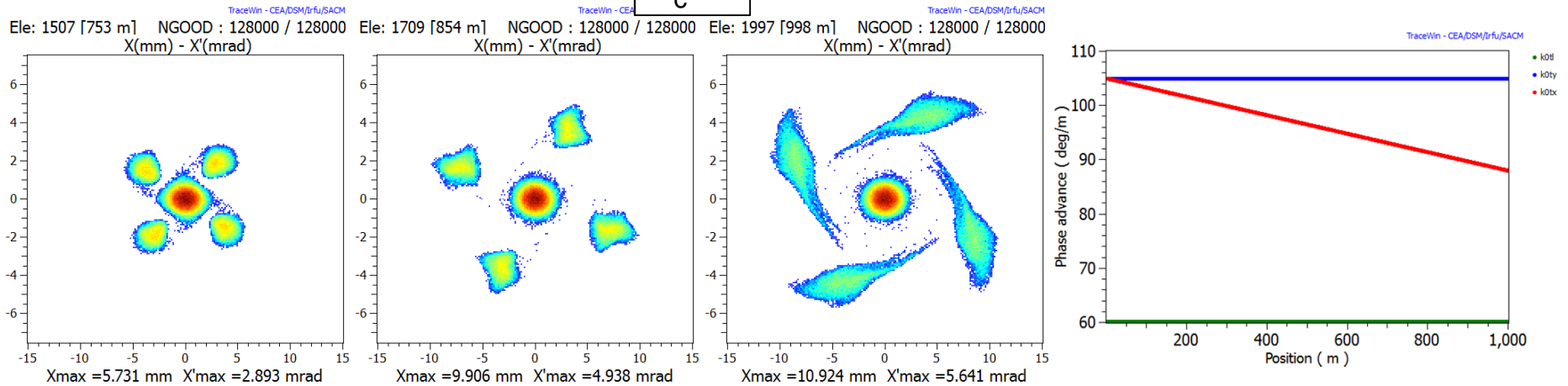




# Crossing of $90^\circ$ (4<sup>th</sup> order) space charge driven resonance $k_{0x}=105^\circ \rightarrow 88^\circ$ / 1000 cells ( $k_{0y}=105^\circ$ fixed)

128 k particles,  $n_c=4-6$ : find 4 „identical“ islands pushed outwards

$n_c=6$

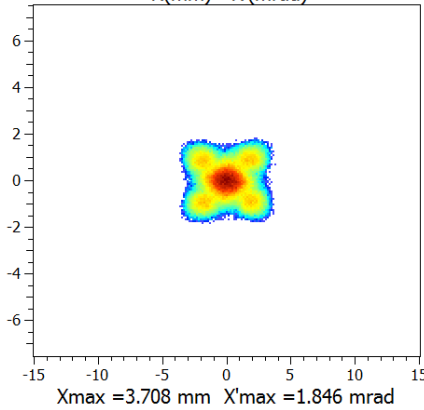




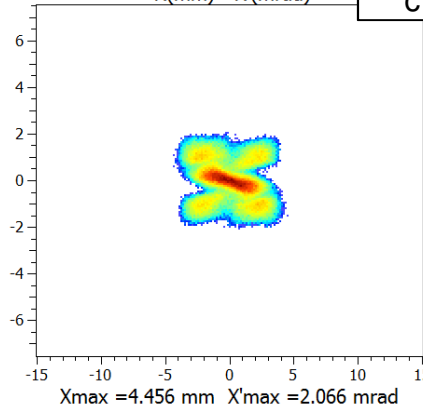
# cont'd crossing $k_{0x}=105^0 \rightarrow 88^0$ /1000 cells

128 k particles,  $n_c=7-12$ : find breaking of four-fold symmetry > 500 cells  
envelope instability ( $h=2$ ) on top of  $h=4$ ?

Ele: 901 [450 m] NGOOD : 128000 / 128000  
X(mm) - X'(mrad)

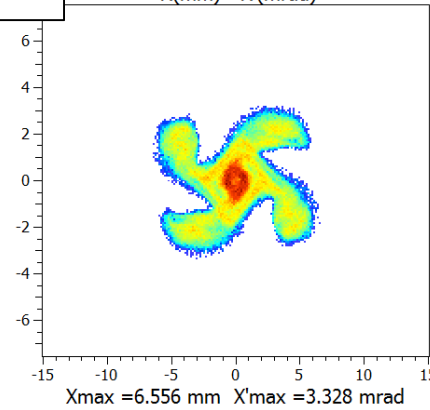


Ele: 1141 [570 m] NGOOD : 128000 / 128000  
X(mm) - X'(mrad)

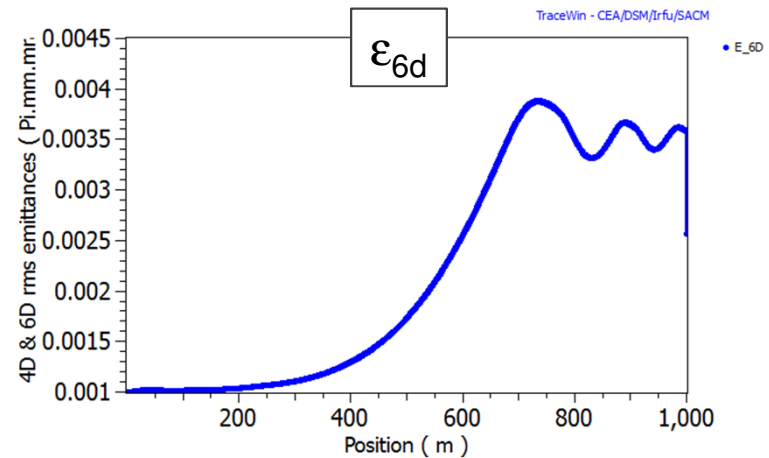
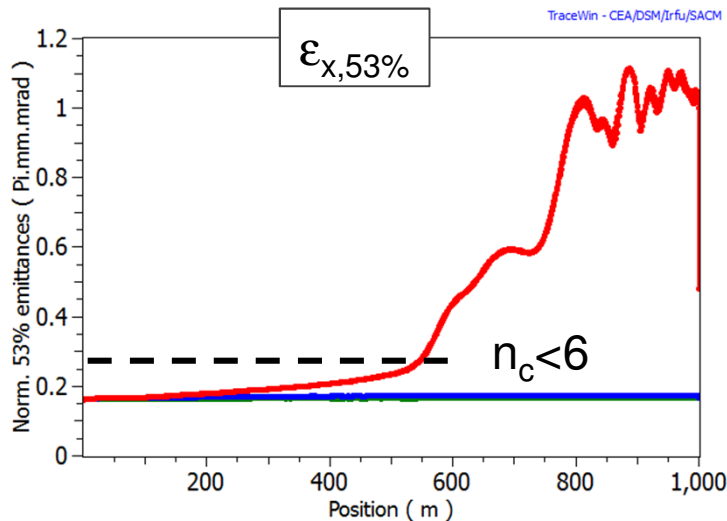
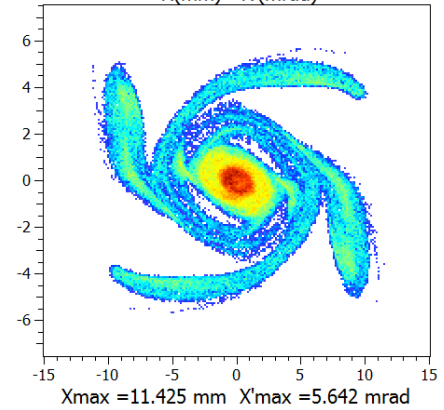


$n_c=12$

Ele: 1537 [768 m] NGOOD : 128000 / 128000  
X(mm) - X'(mrad)

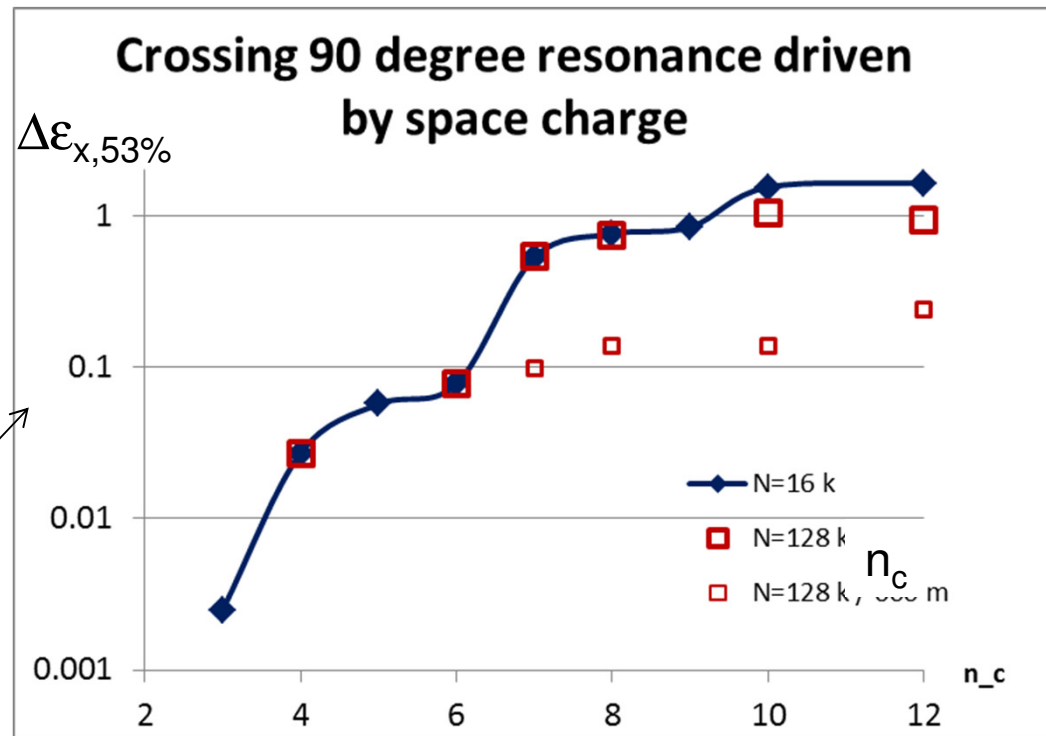
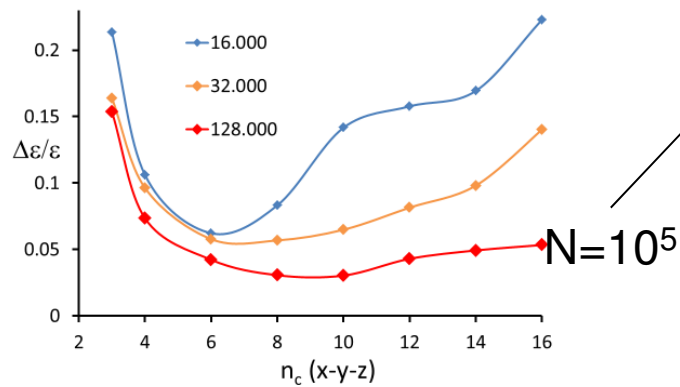


Ele: 1995 [997 m] NGOOD : 128000 / 128000  
X(mm) - X'(mrad)



Summary: crossing  $k_{0x}=105^0 \rightarrow 88^0$  /1000 cells

question: why more growth for finer grid, if additional mode is only 2nd order?



# Summary

- ✓ 3 d noise by grid well represented by rms entropy in equilibrium beam
- ✓ Grid heating ( $n_c < 5$ ) or collisional heating ( $n_c > 10$ ) for small/large number of grid cells  $n_c \rightarrow$  optimum  $n_c$  for given  $N$
- ✓ Anisotropy effect in „good“ agreement with theory
- ✓ Dynamical problems:
  - ✓ our noise modelling seems irrelevant for fast processes – „no“effect of noise
  - ✓ slow resonance crossing: retrieve noise dependence on grid, but different for island trapping  $\rightarrow$  much increased halo for larger  $n_c$  ?
- ✓  $\rightarrow$  needs more work especially towards resonance (collective) effects