STEADY STATE MULTIPACTING IN A MICRO-PLUSE ELECTRON GUN

Zhou Kui^{*}, Lu Xiangyang[†], Quan Shengwen, Luo Xing, Yang Ziqin, Zhao Jifei State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing 100871, China

Abstract

Multipacting is a resonant electron discharge phenomenon via secondary electron emission, while micropulse electron gun (MPG) utilizes the multipacting current in a radio-frequency (RF) cavity to produce short pulse electron beams. The concept of MPG has been proposed for many years. However, the unstable operating state of MPG vastly obstructs its practical applications. This paper presents a study on the steady state mulitpacting in a MPG. The requirements for steady state multipacting are proposed through the analysis of the interaction between the RF cavity and the beam load. Accordingly, a MPG cavity with the frequency of 2856 MHz has been designed and constructed. Various kinds of grid-anodes are tested in our primary experiments. Both the unstable and stable multipacting current have been observed. Presently, the stable output beam current has been detected at about 12.2 mA. Further experimental study is under way now.

INTRODUCTION

Multipacting is a resonant electron discharge phenomenon via secondary electron emission [1], which is frequently observed in microwave systems. When the multipacting effect occurs, it usually causes some undesirable problems, such as deterirorating the vacuum, absorbing incident power, leading to quenching of superconducting cavity, etc. So, most studies on multipacting focus on how to suppress or eliminate it. Until 1993, Fredrick M. Mako and William Peter proposed the concept of MPG [2], which utilized the multipacting current in a RF cavity to produce pulsed electron beams. Due to its self-bunching property, MPG is capable of providing high current and short pulse electron beams. In addition to that, its simple structure and high tolerance to contamination make it a potential electron source for accelerators and microwave systems.

The main problem in the development of MPG is the stability of the output beam current. To get high beam current and low emittance, previous studies on MPG always tried to feed high power into the cavity to get high surface field [3,4]. The stable multipacting current in the MPG was considered to be formed as the equilibrium result of the self-bunching effect and the space charge effect [1]. However, we believe the beam loading effect also plays an important role in the forming process of stable multipacting current.

This paper presents a study on the steady state mulitpacting in a micro-pluse electron gun with theory and experiments. In the second section, a MPG model is setup to show the basic characteristics of the MPG. The requirements for



Figure 1: The schematic diagram of the MPG model [3].

the steady state multipacting is proposed through the analysis of the interaction of the RF cavity and electron beams. In the third section, a MPG cavity with the frequency of 2856 MHz has been designed and constructed. Different kinds of grid-anodes are tested in our experiments. Both the unstable and stable multipacting current have been observed successfully. The detected stable output beam current has reached 12.2 mA. Finally, a conclusion of this paper is given.

ANALYSIS OF STEADY STATE MP

The MPG Model

The MPG model (Fig. 1) consists of three parts: an RF cavity working in the TM_{010} mode, a secondary emission surface and a grid, which is opaque to the microwave electric field but partially transparent to the electrons in order to extract the electron beams [3]. We suppose the secondary emission surface to be the cathode with secondary emission yield (SEY) δ_1 and the grid to be the grid-anode with SEY δ_2 and transmission coefficient T. Here, we adopt Vaughan's empirical formula for δ_1 and δ_2 [5].

$$\delta_1(E_i) = \delta_{\max 1} (v_1 e^{1 - v_1})^k$$
(1a)

$$\delta_2(E_i) = \delta_{\max 2} (v_2 e^{1 - v_2})^k$$
 (1b)

Where E_i is the electrons impact energy, δ_{max1} and δ_{max2} is the maximum value of δ_1 and δ_2 , $v_1 = (E_{i1} - E_0)/(E_{max1} - E_0)$, $v_2 = (E_{i2} - E_0)/(E_{max2} - E_0)$, in which E_{max1} and E_{max2} are the impact energy corresponding to δ_{max1} and δ_{max2} , E_0 being the initial energy of secondary electrons and k = 0.62 for v < 1; k = 0.25 for v > 1. For one RF period, the total effective secondary electron emission yield δ is:

^{*} Email: zhoukui@pku.edu.cn

[†] Corresponding author, email: xylu@pku.edu.cn



Figure 2: The resonant gap voltage range versus the gap d when the cavity frequency is 2856 MHz.

$$\delta = \delta_1(E_i)\delta_2(E_i)(1-T) \tag{2}$$

Now, we consider a RF cavity operating at the resonant frequency ω_0 . The gap distance between the cathode and the grid-anode is d. Under the force of the electric field, electrons in the cavity move back and forth between the cathode and the grid-anode. The transit time from one anode to the other is just one-half RF period [1]. We obtain the resonant condition:

$$V_c = \frac{m}{e} \frac{\omega_0 d(\omega_0 d - \pi v_0)}{2\sin\varphi_0 + \pi\cos\varphi_0}$$
(3)

In which, m is the mass of electron, e is the elementary electric charge, v_0 is the initial velocity of the emitted electrons and φ_0 is the initial phase of the RF field.

The MPG also needs to satisfy the requirement of the self-bunching property. The self-bunching property demands the secondary electrons emitted within an appropriate initial phase range. For electrons with emitted energy $E_0 = 0$, the appropriate initial phase range is 0 to $\arctan 2/\pi$ ($\approx 32.5^{\circ}$) [1]. Substituting the phase range into the resonant condition (Eq. 3), we get the corresponding resonant gap voltage range:

$$V_{min} = \frac{m}{e} \frac{\omega_0^2 d^2}{\sqrt{\pi^2 + 4}}$$
(4a)

$$V_{max} = \frac{m}{e} \frac{\omega_0^2 d^2}{\pi}$$
(4b)

Figure 2 shows the resonant gap voltage region (light green area) varying with the gap distance d when the cavity frequency is 2856 MHz. If the gap voltage is outside of this resonant range, the electrons cannot gather together in the longitudinal direction so that the electron beam will break up.

Requirements for Steady State Multipacting

Basically, the higher the gap voltage is, the more energy will be gained. According to the secondary emission prop-

ISBN 978-3-95450-133-5





Figure 3: The total effective SEY curve versus the gap voltage V_c . (E_1 is the first crossover point and E_2 is the second one. V_1 and V_2 are the gap voltages corresponding to E_1 and E_2 . [V_{min1} , V_{max1}], [V_{min2} , V_{max2}] and [V_{min3} , V_{max3}] are the corresponding resonant gap voltage ranges when the gap distance d takes different values. In addition, $V_1 \in [V_{min1}, V_{max1}]$ and $V_2 \in [V_{min2}, V_{max2}]$.)

erty of common meterials, we can map the total effective SEY curve versus the gap voltage qualitatively (Fig. 3).

Supposing that the resonant gap voltage range is $[V_{min1}, V_{max1}]$, if the gap voltage V_c is greater (less) than V_1 , more (fewer) secondary electrons will be released, so the number of electrons N_e will increase (decrease). As a result, more (less) microwave power will be drained from the cavity as it accelerates more (fewer) secondary electrons. This leads to a higher (lower) P_b and lower (higher) P_c as well as a lower (higher) gap voltage V_c , which makes V_c closer to V_1 in subsequent RF periods. Similarly, V_c will move far away from V_2 automatically if $[V_{min2}, V_{max2}]$ is the resonant gap voltage range. Thus E_1 is the steady state operating point for the MPG, not E_2 . This conclusion is consistent with R. Kishek and Y.Y. Lau's theoretical analysis on multipactor discharge in a closed RF cavity [6,7].

However, if the resonat gap voltage range is set to be $[V_{min3}, V_{max3}]$ as many previous studies did, δ will be always greater than 1. When the gap voltage reaches the resonant range, the multipacting current will grow to a very large value in a very short time. A great deal of microwave power will be drained from the cavity to accelerate these multiplied electrons. So the gap voltage will eventually fall outside of the resonant range and the electron beam will break up due to the longitudinal debunching. Therefore, there is no stable working point in $[V_{min3}, V_{max3}]$. This explains why previous MPGs could not work stably.

To obtain steady state multipacting, the working point should be set in the vicinity of the first crossover point E_1 and $[V_{min1}, V_{max1}]$ should be set as the resonant gap voltage range. Meanwhile, V_1 must be included in this resonant range. Noting that the value of E_1 depends on the secondary emission property of the cathode and the grid-anode as well as the transmission factor T (Eq. 2), while the resonant gap voltage range [V_{min1}, V_{max1}] is determined by ω_0 and d.



Figure 4: (a) the cross section of the MPG, (b) the magnetic field distribution, (c) the electric field distribution.

Table 1: The Designed RF Parameters of the MPG Cavity

RF parameters	Designed Value
Resonant frequency f_0 (MHz) :	2856
Gap distance d (mm)	1.75
Unloaded quality factor Q_0	653
Shunt impedance r_{shunt} (M Ω)	0.036
Unloaded coupling coefficient β_0	1.25

So, as long as the above requirements are satisfied, it is possible to get steady state multipacting in the MPG.

PRIMARY EXPERIMENTS

The MPG Cavity and the Test Stand

According to the requirements for steady state MP, a micro-pulse electron gun working in the TM₀₁₀ mode with the frequency of 2856 MHz has been designed and constructed. Table 1 lists the designed RF parameters of the MPG cavity. The micro-pulse electron gun (Fig. 4a) consists of a RF cavity made of stainless steel, a cathode made of Cu-Al-Mg alloy and a grid-anode. The impact energy corresponding to the maximum secondary emission yield δ_{max} is about 1 KeV for Cu-Al-Mg alloy [8]. The gap distance between the cathode and the grid-anode is adjustable with the precision of 0.01 mm. The cathode and the gridanode are also dismountable and can be replaced easily. The distribution of the magnetic field and the electric field are shown in Fig. 4b and Fig. 4c. Obviously, the electric field is mainly concentrated in the region between the cathode and the grid-anode.

The test stand is shown in Fig. 5. The microwave power is supplied by a soild-state power amplifier with 1 kW maximum output power. The output electron beams from the grid-anode are collected by a faraday cup parallelly connected to a 50 Ω resistance and the beam current is measured by a HP 54503A oscilloscope with 1 M Ω input resistance.

Table 2: The Measured Output Beam Currents (I_b) with Various Grid-anodes

Grid-anode No.	1	2	3	4
Material	SS	SS	OFC	OFC
Т	6%	18.3%	18.3%	25%
Measured I_b (mA)	≈0.2	3.8	4.2	12.2



Figure 5: Photogragh of the test stand.

Experiment Results and Discussions

Various grid-anodes are tested in our experiments. Table 2 lists the measured output beam currents corresponding to four kinds of grid-anodes with different meterial and transmission coefficient (T). In Table 2, SS refers to stainless steel and OFC represents oxy free copper. Figure 6 shows our primary experimental results. The macro-pulse width is 15 us and the repetition rate is about 100 Hz. Figure 6a shows the output beam current when No.1 grid-anode is used. The output siganl is very small and unstable on the oscilloscope. The waveform is fluctuating with the "spikes" on the pick-up signal and reflected signal in Fig. 6b. Figure 6c performs a very stable output beam current when the grid-anode is No.2. The signal waveform can remain unchanged for more than half an hour, which reveals that we have already obtained the steady state multipacting in the MPG. The output beam current is about 1.7 mA. Moreover, it can be adjusted up to 3.8 mA, which is shown in Fig. 6d. We also obtained the stable output beam current successfully, when No.3 and No.4 grid-anode are tested (Fig. 6e and Fig.6f). The output beam current measured in the experiments are about 4.2 mA (T=18.3%) and 12.2 mA (T=25%) respectively.

The spikes appearing on the pick-up signal and the reflected signal in Fig. 6b indicate the typical unstable multipacting processes in the cavity. In this situation, the total effective SEY is above 1, but the actual working point is not in the vicinity of E_1 . So, the electrons in the cavity grows rapidly without restraint. The multiplied electrons absorb plenty of RF power, leading to the cavity voltage falling out of the resonant gap voltage range, so the electron beam collapses. After that, the cavity voltage returns to the initial state. The remaining electrons in the cavity will begin



Figure 6: Primary experimental results. (a) is the output beam current ($\approx 0.2 \text{ mA}$) when the grid-anode is No.1; (b) is the pick-up signal and reflected signal corresponding to (a); (c) is the output beam current ($\approx 1.7 \text{ mA}$) when the grid-anode is No.2; (d) is the maximum detected output beam current ($\approx 3.8 \text{ mA}$) when the grid-anode is No.2; (e) is the output beam current ($\approx 4.2 \text{ mA}$) when No.3 grid-anode is tested; (f) is the output beam current ($\approx 12.2 \text{ mA}$) when No.4 grid-anode is tested.

the multiplication again and a new cycle keeps on. So the "spikes" appears on the pick-up signal successively as well as the reflected signal due to the variation of the loaded coupling coefficient.

CONCLUSION

In summary, this paper presents a study on the steady state multipacting in a micro-pulse electron gun with theory and experiments. In theory, the requirements for steady state multipacting is proposed. The working point should be set in the vicinity of the first crossover point E_1 and $[V_{min1}, V_{max1}]$ should be set as the resonant gap voltage range. Meanwhile, V_1 must be included in this range. As for experiments, a MPG cavity with the frequency of 2856 MHz has been designed and constructed. Various kinds of gridanodes are tested in our experiments. Both the unstable and stable multipacting current have been observed. Presently, the measured stable output beam current has reached 12.2 mA. Further experimental study is under way now.

ACKNOWLEDGMENT

We gratefully acknowledge the assistance of Li Ming and Yang Xingfan in providing the solid-state power amplifier. We also thank Wang Xinping for preparing the Cu-Al-Mg alloy.

REFERENCES

- J. R. M. Vaughan, IEEE Trans. Electron Devices 35(7), 1172-1180 (1988).
- [2] F. M. Mako and W. Peter, PAC'93, Washington, 4, 2702-2704 (1993).
- [3] L. K. Len and F. M. Mako, PAC'99, New York, 1, 70-74 (1999).
- [4] J. Y. Zhai, C. X. Tang, S. X. Zheng, EPAC'06, Edinburgh, 3203-3205 (2006).
- [5] J. R. M. Vaughan, IEEE Trans. Electron Devices 36(9), 1963-1967 (1989).
- [6] R. Kishek and Y. Y. Lau, Phys. Rev. Lett. 75(6), 1218-1221 (1995).
- [7] R. Kishek, Y. Y. Lau et al., Phys. Plasmas 5(5), 2120-2126 (1998).
- [8] Q. H. Pan, The Chinese Journal of Nonferrous Metals 10(3), 374-377 (2000).