

Echo scheme: Alternative to Slicing in Storage Rings

Frequency mixing scheme: for tunability of seeded FELs

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Outline

- ① **Echo scheme : Alternative to Slicing in Storage Rings**
 - Introduction : femtosecond radiation in storage rings
 - EEHG on storage rings (SOLEIL case)

- ② **Frequency mixing : increase of tunability in seeded FELs**

Outline

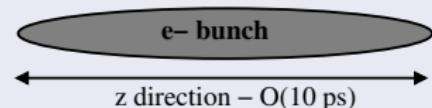
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Femtosecond radiation in storage rings

Standard pulse duration

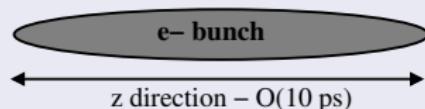
- Pulse light duration \sim electron-bunch duration



Femtosecond radiation in storage rings

Standard pulse duration

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Slicing [A. Zholents et. al., PRL 1996] [R. W. Schoenlein et al., Science 2000][S. Khan et. al, PRL 2008] ...

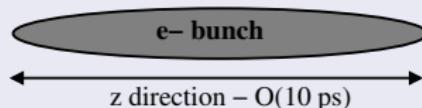
- Small part of the (incoherent) radiation kept
- **Advantages :** full tunability & hard X-rays
- **Drawback :** low peak power

- ALS (USA)
- Bessy II (Germany)
- SLS (Switzerland)
- SOLEIL (France)
- (etc.)

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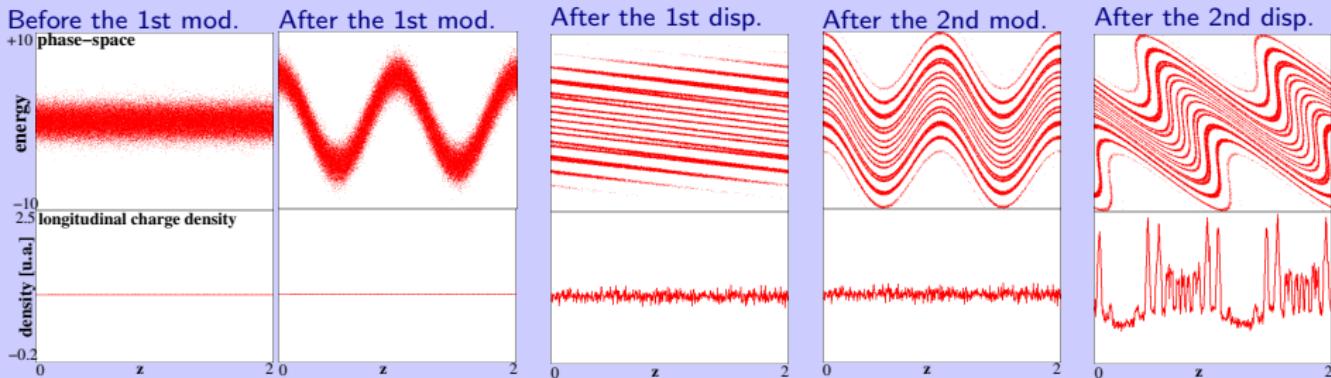
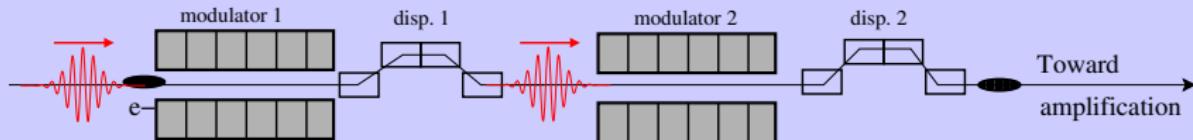
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Coherent Harmonic Generation [R. Prazeres et. al., NIMA 1991][V.N. Litvinenko, NIMA 2003] [M. Labat et. al., NIMA 2008] [G. De Ninno et. al., PRL 2008]...

- Coherent radiation at laser harmonics
- **Advantage :** higher photon flux than slicing
- **Drawbacks :** tunability & "long" wavelength

- UVSOR (Japan)
- DELTA (Germany)
- ELETTRA (Italy)
- (etc.)

EEHG scheme [G. Stupakov, PRL 2009][D. Xiang et. al, PRL 2012]...



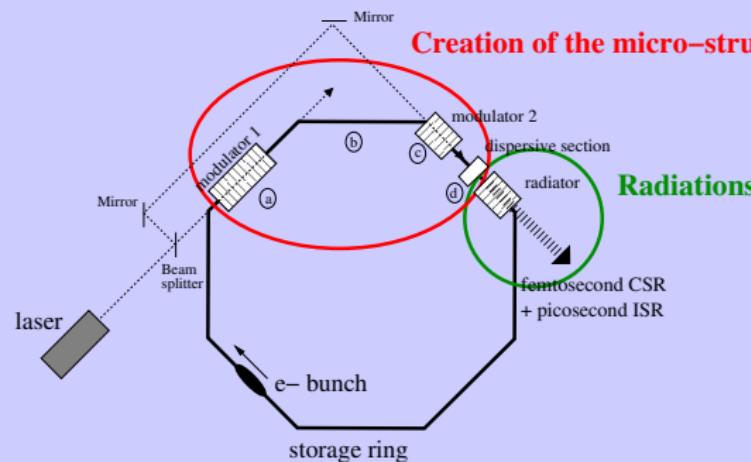
- Application to storage rings ?
- What type of radiations ?

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Layout



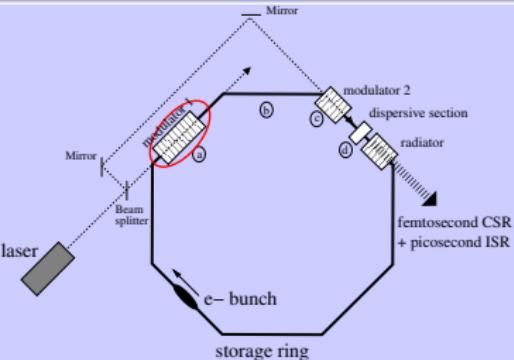
SOLEIL parameters

- $E_0 = 2.75 \text{ GeV}$
- $\sigma_E = 10^{-3}$
- Peak current $I = 134 \text{ A}$
- Bunch length $\sigma_z = 10.5 \text{ mm}$

Studies on the application of EEHG on storage rings :

- SOLEIL : [C. Evain, M.E. Couplie, A. Nadji, A. Loulergue, J.M. Filhol, A.A. Zholents, New J. Phys. 2012]
- DELTA : [R. Molo et. al., this conference - today - WEPS043]
- HEFEI : [H. Li, W. Gao, Q. Jia, L. Wang, IPAC 2013]

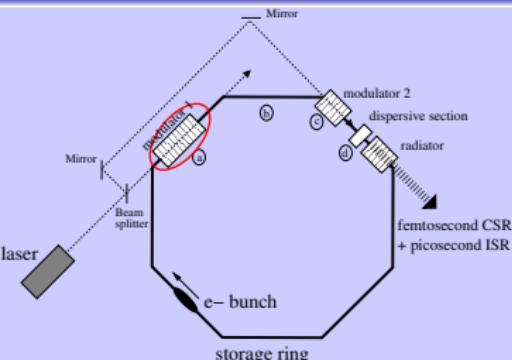
Electron bunch dynamics



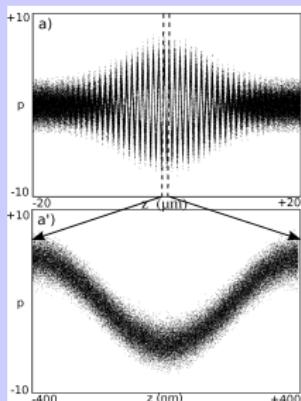
Model

- 6D
- Linear & Non-linear terms for the transport in the magnetic elements
- Energy fluctuation induced by ISR in bending magnets

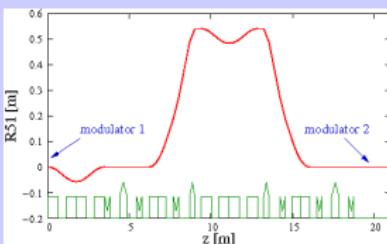
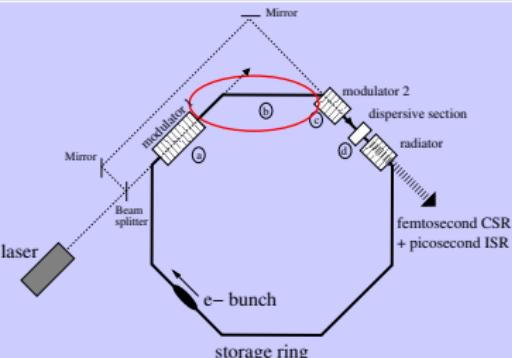
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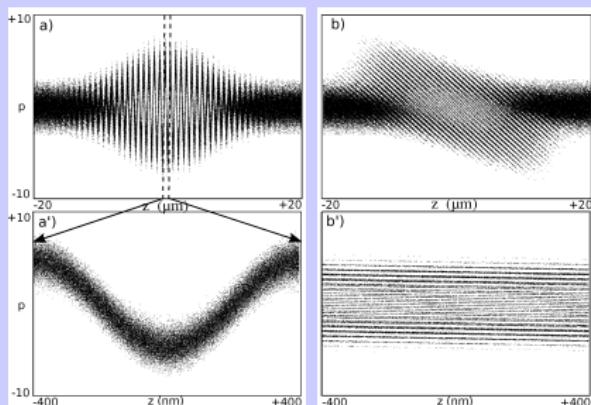
- $p_x = p_0 + A_1 \times e^{-\frac{z^2}{2(\sigma_{L1})^2}} \cos(\frac{2\pi}{\lambda_L} z) \times e^{-\frac{x^2+y^2}{w_1^2}}$
- σ_{L1} : rms laser pulse length, w_1 : laser waist
- **Parameter values taken :**
 - $W_1 = 600\mu\text{m}$ ($\sigma_x = 147\mu\text{m}$), $\lambda_L = 800$ nm
 - $\sigma_{L1} = 43$ fs (100 fs FWHM) , $A_1 = 5$
 - Modulator : 13 periods of length : 150 mm



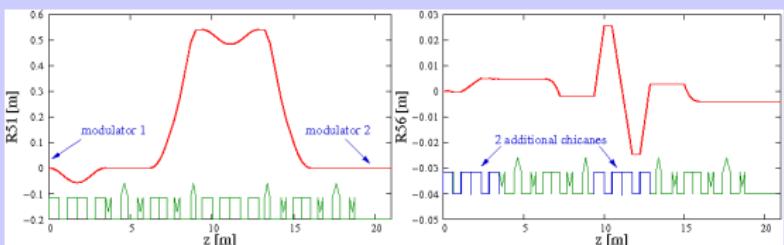
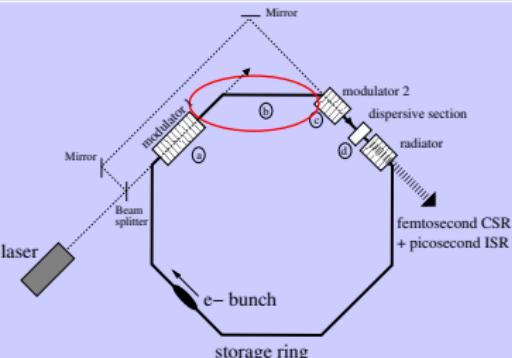
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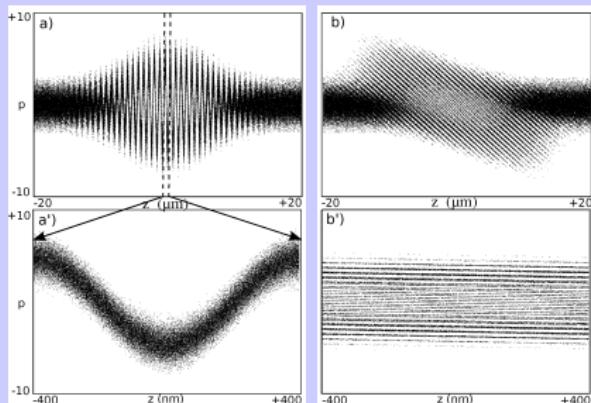
● Chasman-Green lattice \Rightarrow achromatic lattice



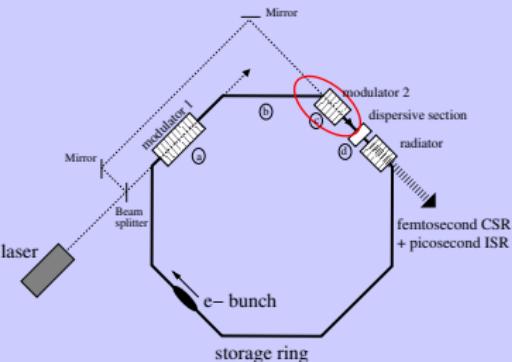
Electron bunch dynamics



- Chasman-Green lattice \Rightarrow achromatic lattice
- "Small" $R_{56}^{(1)}$ \Rightarrow additional chicane (here, $R_{56}^{(1)} = -1.5$ mm)

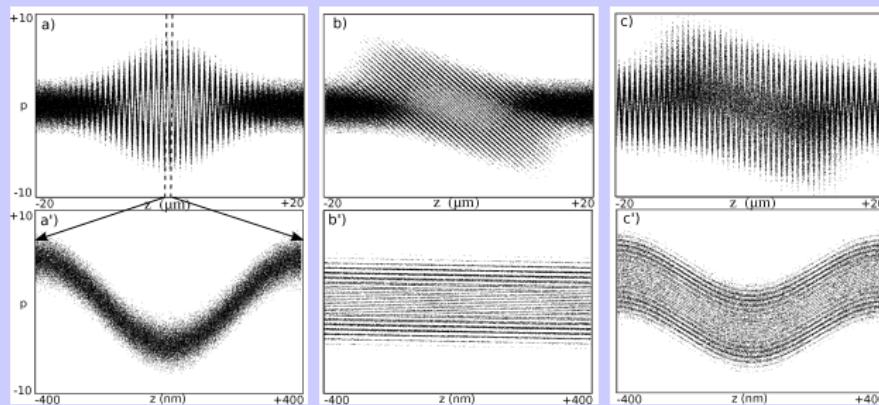


Electron bunch dynamics

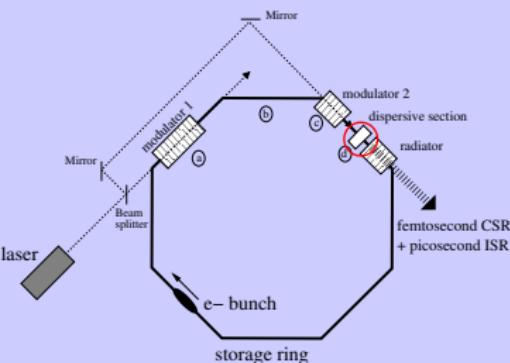


- $p = p_0 + A_2 \times e^{-\frac{z^2}{2(c\sigma_{L2})^2}} \cos(\frac{2\pi}{\lambda_L} z) \times e^{-\frac{x^2+y^2}{w_2^2}}$

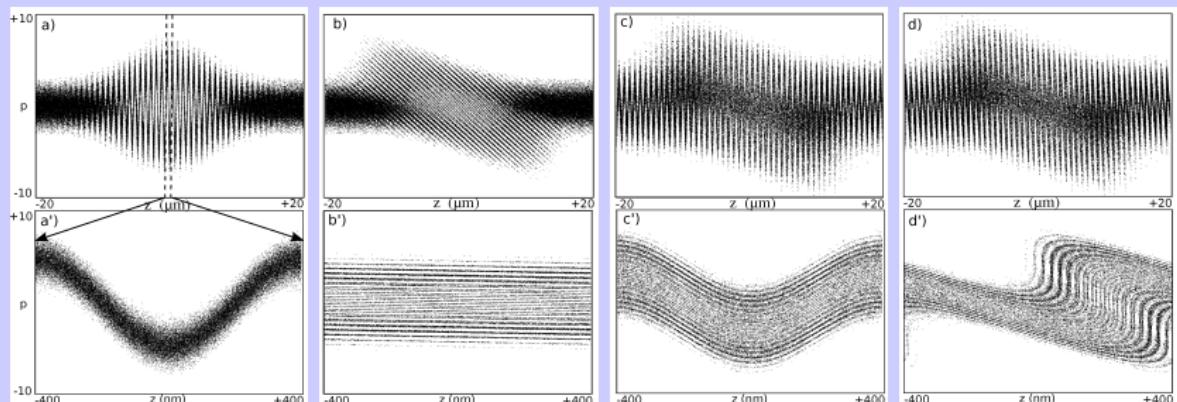
- $w_2 = 600\mu\text{m}$
- A_2 value depends of the wanted harmonic number
(here : $A_2 = 2.95$)



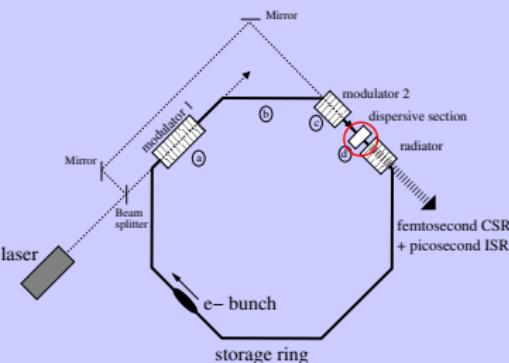
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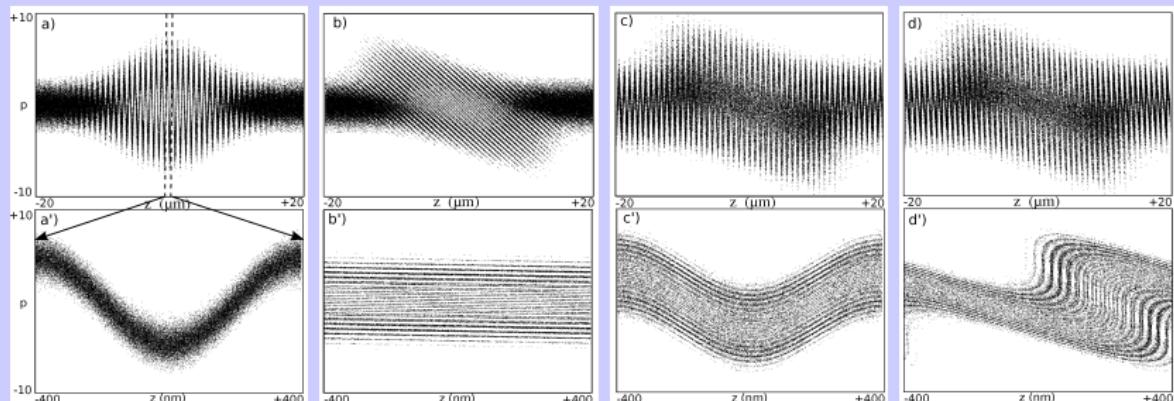
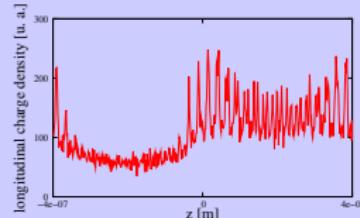
- $z \simeq z_0 + p \times R_{56}^{(2)} \frac{\sigma_E}{E_0}$
- $R_{56}^{(2)}$ value depends of the wanted harmonic number
(here : $R_{56}^{(2)} = -48\mu\text{m}$)



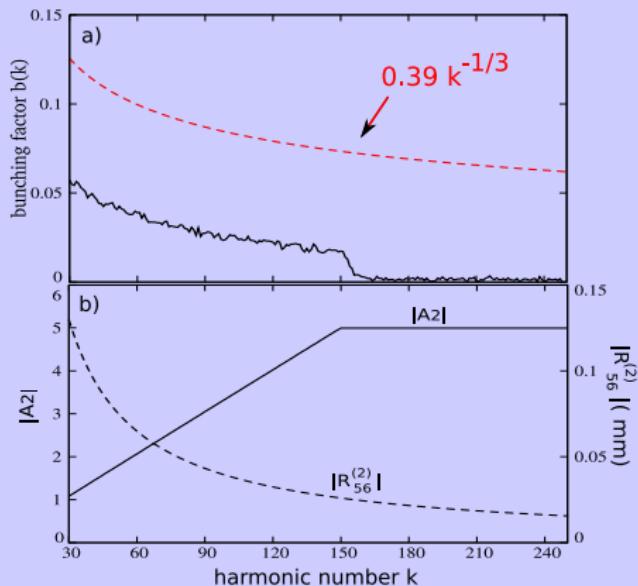
Electron bunch dynamics



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- $R_{56}^{(2)}$ value depends of the wanted harmonic number
(here : $R_{56}^{(2)} = -48\mu\text{m}$)
- Modulation of the longitudinal charge density



Bunching factor vs harmonic number



- $R_{56}^{(1)} = -4 \text{ mm}$
- $A_1 \simeq 5$
- $\sigma_{L1} = 100 \text{ fs}, \sigma_{L2} = 275 \text{ fs}$ (FWHM)

Analytical formula : [D. Xiang and G. Stupakov, PRSTAB (2010)]

Radiation peak power (TEMPO Beamline : $N_u = 19$, $\lambda_u = 80$ mm)

- Analytical formula

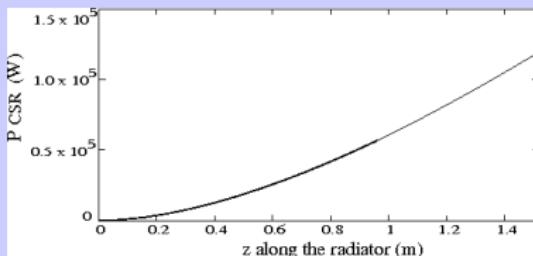
$$P_{CSR} = \pi \alpha \hbar \omega \frac{K^2}{1 + K^2/2} [JJ]^2 \frac{I_{peak}}{e} n_e b^2 \sqrt{f_2}.$$

- $n_e = \frac{I_{peak} \lambda_r N_u}{ce}$,
- $f_2 = (\sigma_r \sigma_{r'})^2 / (\sqrt{\sigma_r^2 + \sigma_x^2} \sqrt{\sigma_{r'}^2 + \sigma_{x'}^2} \sqrt{\sigma_r^2 + \sigma_y^2} \sqrt{\sigma_{r'}^2 + \sigma_{y'}^2})$,
- $\sigma_r = \sqrt{2 \lambda_r \lambda_u N_u} / 4\pi$,
- $\sigma_{r'} = \sqrt{\lambda_r / 2 \lambda_u N_u}$.

[Z. Huang and K.-J. Kim, PAC'99 p. 2495 (1999)]

- @ 26.5 nm (800/30 nm) with $b = 5\%$ and $I_{peak} = 138$ A, $P_{CSR} \simeq 187$ kW
- $P_{CSR} > 10^6 P_{slicing}$ (with $\Delta\omega/\omega = 0.05\%$)

- With GENESIS (using convharm option) :



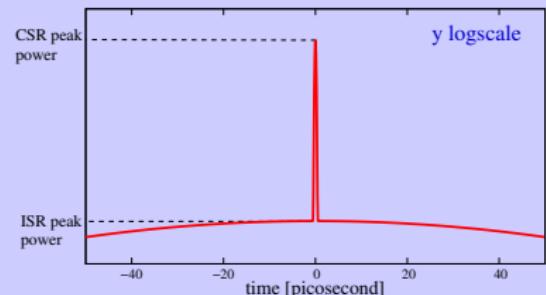
$$P'_{CSR} \simeq 120 \text{ kW}$$

Signal/noise ratio

- Signal-to-noise ratio S/N

$$\begin{aligned} S/N &= \frac{P_{CSR} \times \sigma_{L1} \eta c}{P_{ISR} \times \sigma_z} \\ &= \frac{n_e b^2 \sqrt{f_2} \sigma_{L1} c}{\sigma_z} N_u \frac{\Delta\omega}{\omega} \\ &\simeq 100 \end{aligned}$$

with $\eta = 10\%$ (η : percentage of electrons involved in the fs light pulse)

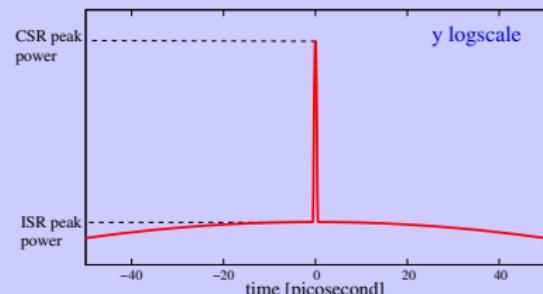


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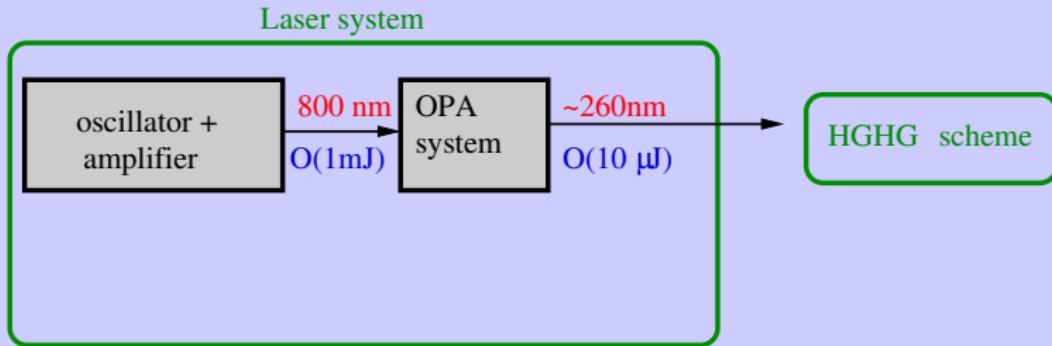
- Conclusion & Perspectives

- Application of EEHG to storage rings very promising
- Challenge : technical realization (test facility/user facility) (cf. Z. Zhao & D. Xiang talks)

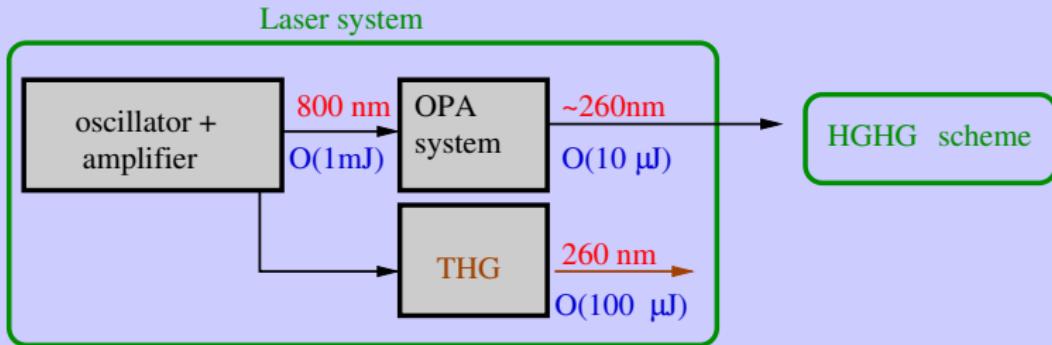
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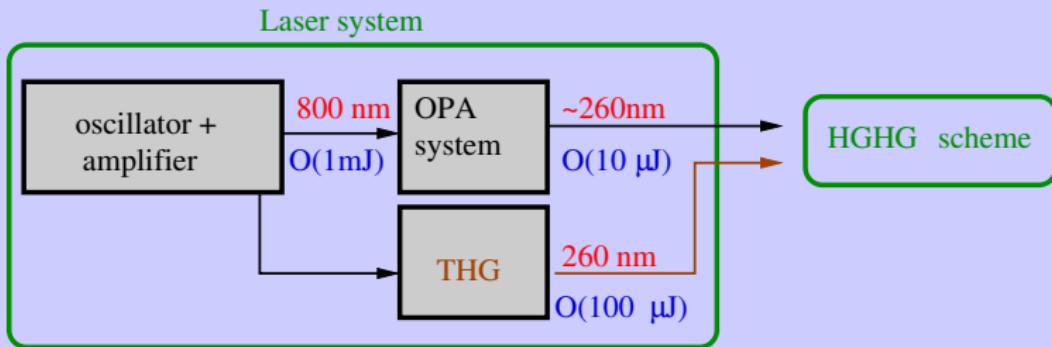
HGHG scheme + tunable seed



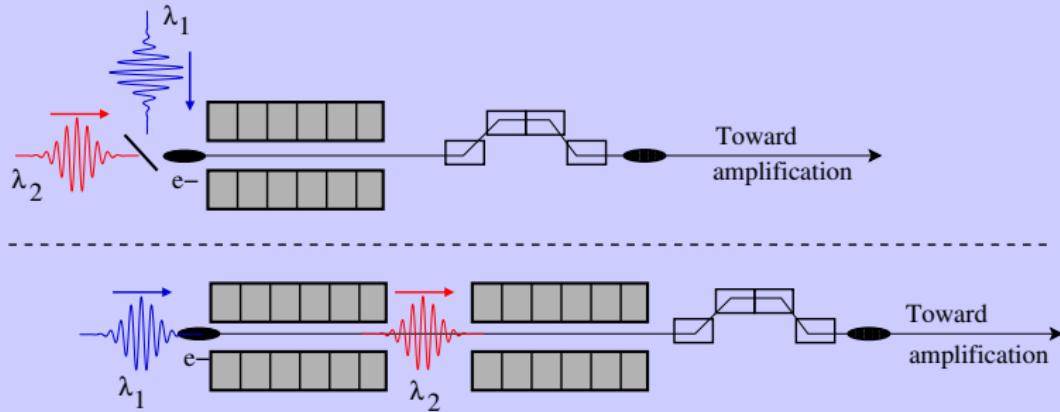
HGHG scheme + tunable seed



HGHG scheme + tunable seed + 2nd fixed seed

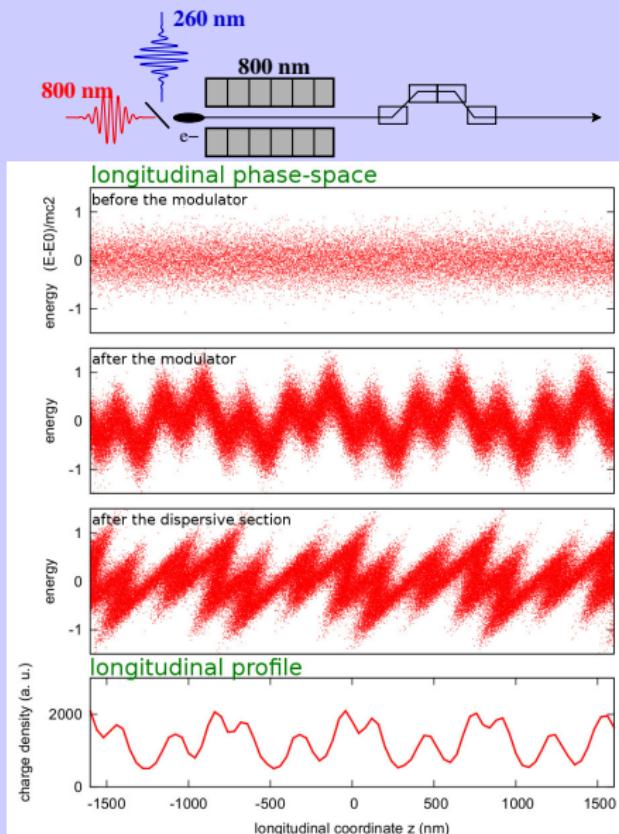


Layout - frequency mixing scheme



- Same layout that in [D. Xiang and G. Stupakov, PRSTAB 2009] for creation of micro-structures in the THz range

Electron dynamics - 1 undulator case & FERMI Parameters



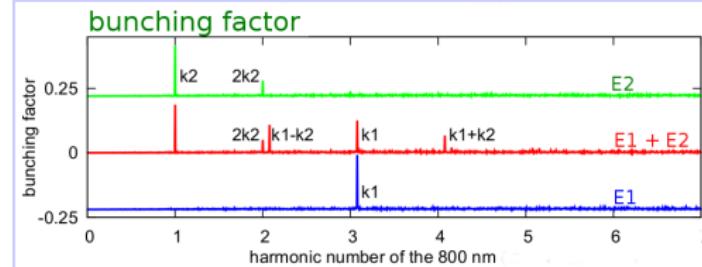
● Fermi & laser parameters

Beam : $E_0 = 1.2$ GeV, $\sigma_E = 150$ KeVModulator : 19 periods, $\lambda_{w1} = 16$ cm, $\lambda_{m1} = 800$ nmlaser : $\lambda_1 = 260$ nm, $\lambda_2 = 800$ nm, $E_1 \simeq E_2 \simeq 10 \mu\text{J}$

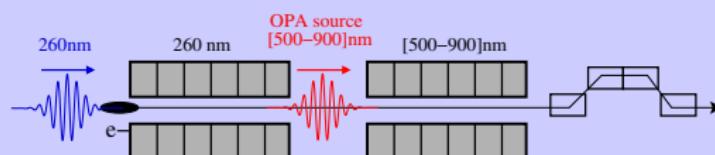
● Model for the laser/e- interaction

see e.g. [A. Zholents, FEL'09]

$$\begin{aligned}\frac{d\gamma}{dz} &= \frac{e}{m_0 c^2} E_x \beta_x, \quad \beta_x = -\frac{K}{\gamma} \sin\left(\frac{2\pi}{\lambda_{w1}} z\right) \\ E_x &= \frac{E_1}{\sqrt{1+\left(\frac{z}{z_{r1}}\right)^2}} \sin[k_1(z-ct)] + \frac{E_2}{\sqrt{1+\left(\frac{z}{z_{r2}}\right)^2}} \sin[k_2(z-ct)]\end{aligned}$$

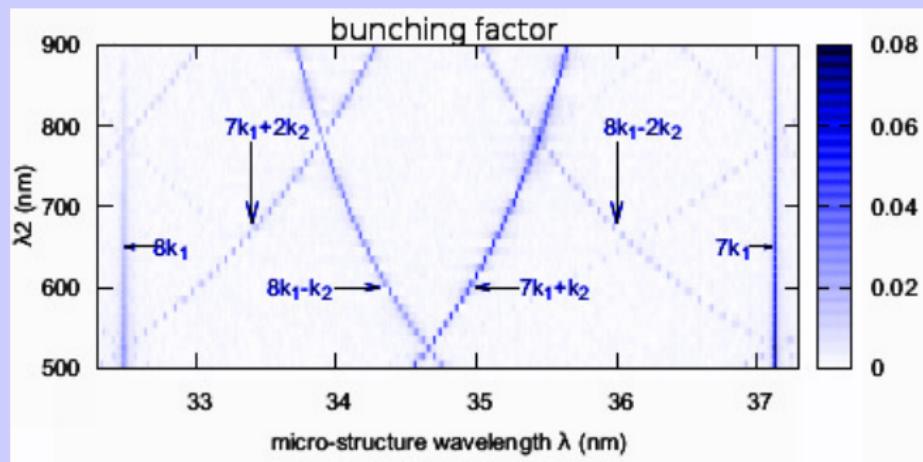


Tunability with an OPA source



Pulse energy (with 100 fs FWHM) :
(obtainable from 1 mJ pulse @800nm)

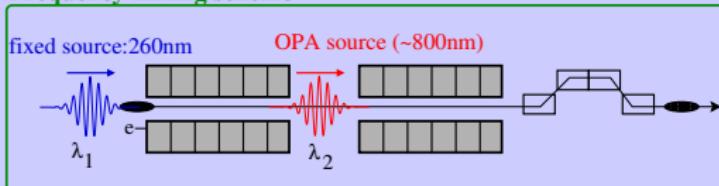
- 260 nm : 130 μ J
- [500–900] nm : 10 μ J



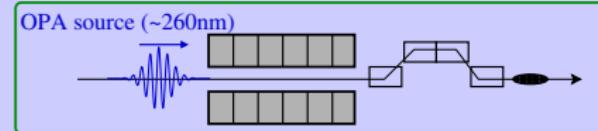
Comparison with HGHG with an OPA source

- Analytical formula from : [D. Xiang, G. Stupakov, PRSTAB 2009] (with $B_1 = 0$)

Frequency mixing scheme



HGHG scheme with an OPA source



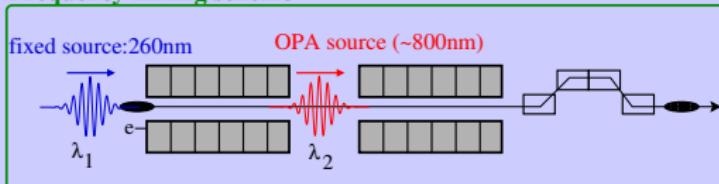
- $k = 7k_1 + k_2 \Rightarrow \lambda_r \simeq 35.5 \text{ nm}$
- $k = 7k_1 + 2k_2 \Rightarrow \lambda_r \simeq 34 \text{ nm}$
- $A_1 = 5$ (260 nm) $A_2 = 1$ (OPA source @800 nm)

- $k = 7k_1 \Rightarrow \lambda_r \simeq 37 \text{ nm}$
- $A_2 = 1$ (OPA source @260 nm)

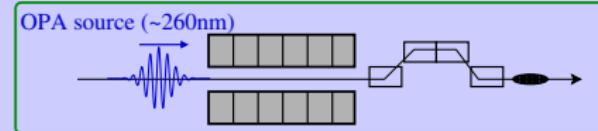
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Frequency mixing scheme



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| | fq. mixing (n=7, m=1) | fq. mixing (n=7, m=2) | HGHG (m=7) |
|-----------------|-----------------------|-----------------------|------------|
| Bunching factor | 5.1% | 2.3% | 0.0034% |
| Ratio with HGHG | ≈ 1500 | ≈ 650 | |

Analytical expression of the bunching factor

- From : [D. Xiang, G. Stupakov, PRSTAB 2009] (with $B_1 = 0$)

$$\begin{aligned} k &= nk_1 + mk_2 \quad \text{with } n,m : \text{integer} \\ bf(k) &= \left| e^{-\frac{1}{2}[(Km+n)B_2]^2} \times J_m [-(Km+n)A_2 B_2] \times J_n [-(Km+n)A_1 B_2] \right| \end{aligned}$$

- A_1, A_2 : energy modulation amplitude (in energy spread unit)
- B_2 : dispersive strength (normalized)
- if $(A_1 < 1 \& A_2 < 1) \Rightarrow \underline{bf_{n,m} \propto E_1^n \times E_2^m}$ (since $A_i \propto E_i$ [A. Zholents FEL'06])
- Same **power law** that in **non-linear optics** : $P = \chi^{n+m} E_1^n \times E_2^m$
(P : polarization, χ^{n+m} : susceptibility of the NL material)

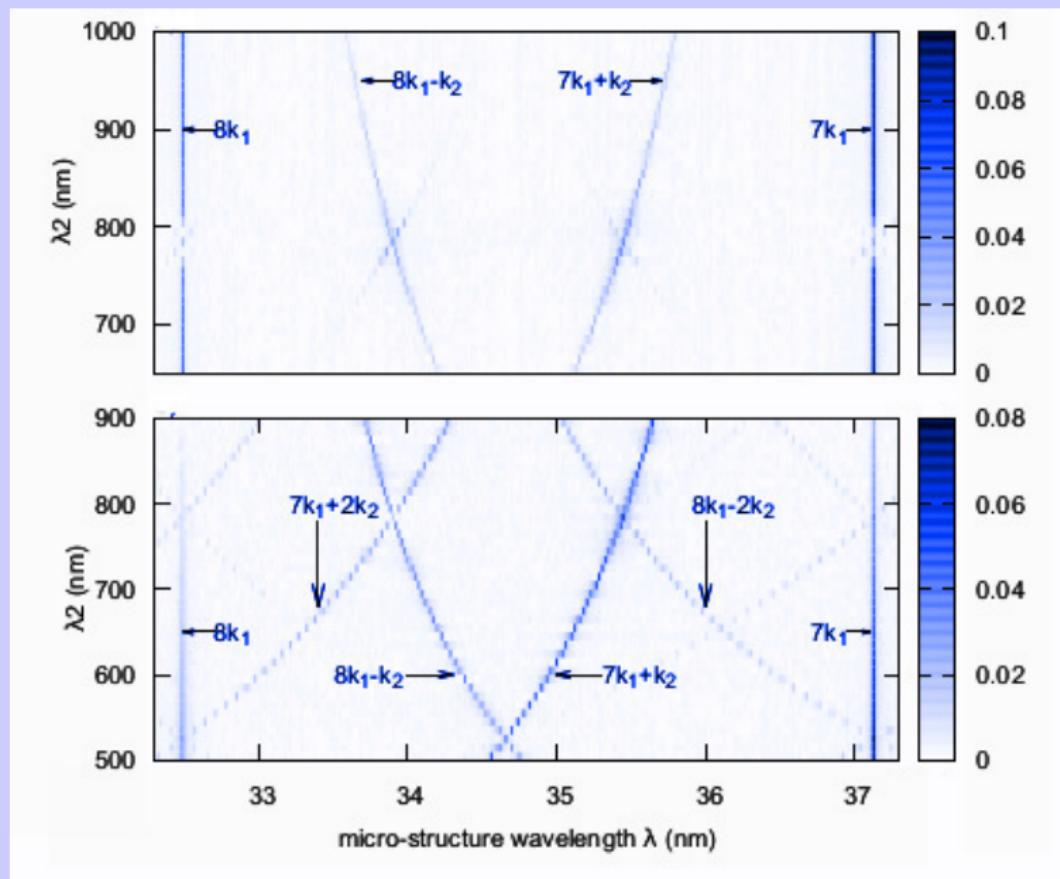
Conclusion

- **EEHG on storage rings**

- Femtosecond coherent synchrotron radiation
- Promising characteristics (wavelength, peak power, signal/noise)
- Challenge : technical realization
(see also application on DELTA storage ring [[WEPS043](#)])

- **Frequency mixing scheme**

- For tunability at "short" wavelengths, it is very efficient to use of an 2nd laser pulse (with higher energy & not tunable), coupled with a tunable source.



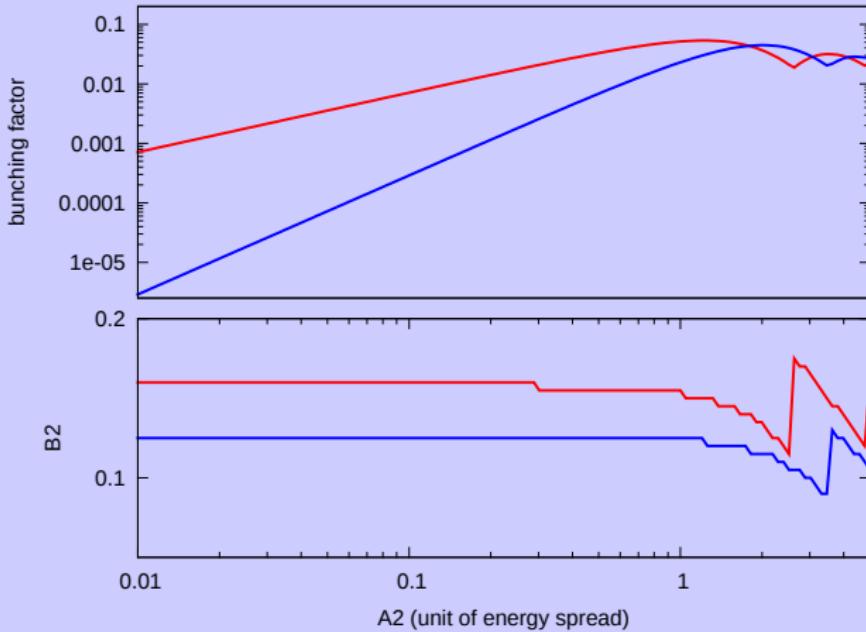
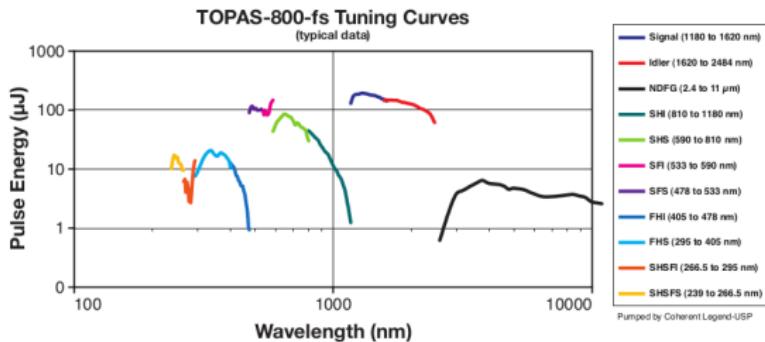


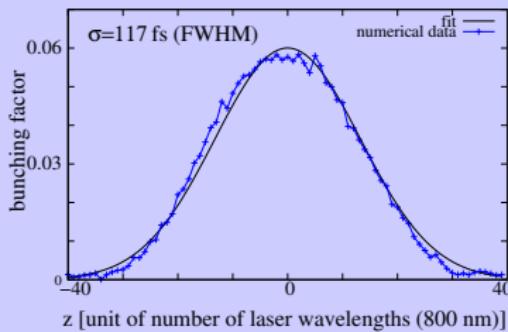
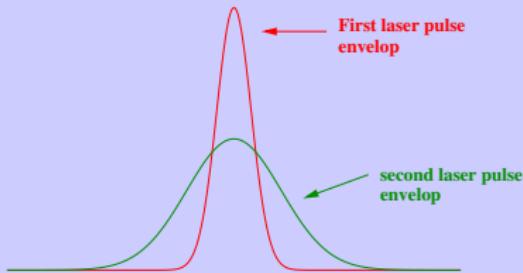
FIGURE: $K=3.07$, $n=7$, blue : $m=1$ and red : $m=2$. $A_1 = 5$

Typical Tuning Curve
(pumped by 1 mJ, <40 fs, 1 kHz, 800 nm Legend)



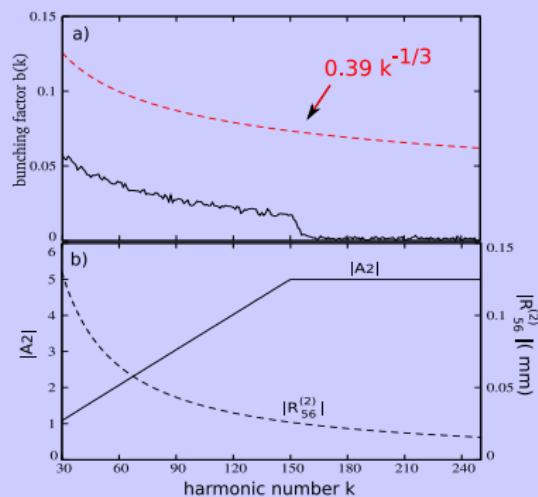
Duration of the CSR pulse

- value taken : $\sigma_{L1} = 100$ fs (FWHM), $\sigma_{L2} = 275$ fs FWHM



- No transverse dispersion and weak longitudinal one compared to the one in slicing \Rightarrow shorter pulse than with slicing scheme ($\sigma_{slicing} \simeq 200$ fs FWHM)

bunching factor with non-linear terms in the transport between the two modulators

 $R_{56}^{(1)} = 4 \text{ mm - linear transport}$  $R_{56}^{(1)} = 4 \text{ mm - non-linear transport}$ 