

SIMULATIONS OF THE LONGITUDINAL INSTABILITY IN THE NEW SLC DAMPING RINGS*

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I. INTRODUCTION

In 1992 a longitudinal, single bunch instability was observed in the SLC damping rings. [1] Beyond a threshold current of 3×10^{10} a “saw-tooth” variation in bunch length and energy spread was observed, a phenomenon that made it practically impossible to operate the SLC collider above threshold. For the 1994 run a new, low-impedance vacuum chamber was installed in both damping rings both to alleviate this problem and to shorten the bunch length. According to recent measurements the bunch length has indeed become shorter, but the “saw-tooth” instability is still seen, now beginning at currents of $1.5 - 2.0 \times 10^{10}$. [2] Fortunately, it appears to be benign and does not seem to limit SLC performance.

In an earlier paper we investigated the single bunch behavior of the SLC damping rings with the old vacuum chamber using time domain tracking and a Vlasov equation approach. [3] When compared to measurements we found: good agreement in the average values of bunch length, energy spread, and synchronous phase shift as functions of current; a 30% discrepancy in threshold current; in agreement, a mode with frequency near 2.5 times the synchrotron frequency (the so-called “sextuple” mode) as signature of the instability and the slope of the mode frequency as function of current. ¹

In the present paper we repeat the exercise of the earlier paper but with a new wakefield. The impedance which used to be inductive has become resistive, leading to different phenomena.

In a recent paper the instability in a purely resistive ring is analyzed using a Vlasov equation approach. [5] It is demonstrated that such an instability is a weak instability, with a growth rate proportional to intensity squared, and one that can be described as the coupling of two quadrupole modes with different radial mode numbers. We will compare our results with this paper. For related papers, see also Refs. [6], [7].

II. THE WAKEFIELD

In the vacuum chamber upgrade of the SLC damping rings primarily small objects that are inductive at nominal bunch lengths (~ 5 mm) were removed or modified, [2] [8] changing the character of the rings from inductive to resistive. As before, we have attempted to find an approximate Green function wakefield $W(z)$ for the new ring using the time-domain parts of the MAFIA family of computer programs, [9] taking as driving bunch a short, gaussian bunch with rms length of 1 mm. To make it causal, the part in front of bunch center ($z < 0$) was reflected and added to the back (see Fig. 1), a transformation that

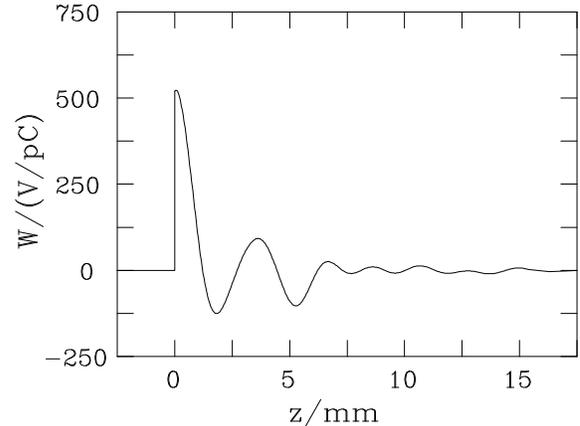


Figure 1. The wakefield used for the simulations.

preserves the real part of the impedance. We expect to be able to find beam instabilities down to wavelengths of about 1 cm.

We have had difficulty obtaining an accurate wakefield, and we are not very satisfied with what we have; it should be considered only preliminary. As we have removed the grosser, cylindrically symmetric objects (45° transitions, masks, etc.) we are left with a machine dominated by objects for which it is difficult to obtain an accurate wake function, such as, for example, the beam position monitors. Also, difficult vacuum chamber objects that could previously be ignored, such as the septum chamber, may now be important. In spite of these misgivings, and even though (as we will see) the results do not agree with the SLC measurements as well as before, we feel that this wakefield is still useful for studying the basic character of the damping ring current dependent behavior.

The induced voltage on any turn is given by

$$V_{ind}(z) = -eN \int_{-\infty}^z W(z-z') \lambda_z(z') dz', \quad (1)$$

with N the bunch population and $\lambda_z(z)$ the longitudinal charge distribution. To see that it is resistive in character at typical bunch lengths we plot in Fig. 2, λ_z and V_{ind} , for $N = 2 \times 10^{10}$ and nominal length $\sigma_{z0} = 5$ mm, the solution to the Haïssinski equation. [10] We note that roughly V_{ind} can be written as $V_{ind} \approx -eN R c \lambda_z$, with R the resistance, a constant. In fact, if we take a pure resistance with $R = 880 \Omega$, and repeat the potential well calculation, we obtain almost the same bunch shape (the dashes in Fig. 2).

III. SIMULATIONS

For tracking we let the beam be represented by N_p macroparticles; each particle i has position and energy coordinates

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¹In Ref. [3] it was stated that the measured slope was three times smaller than the calculations. However, subsequent, more accurate measurements are in agreement with the calculations. [4]

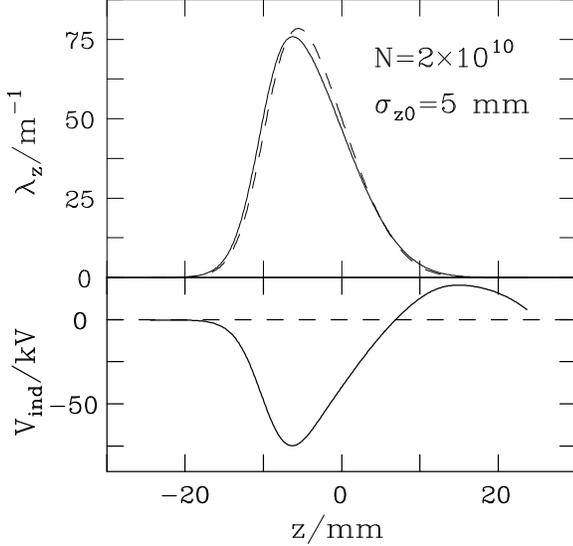


Figure 2. A potential well example.

(z_i, ϵ_i) . The properties of particle i are advanced on each turn according to the equations: [11]

$$\Delta \epsilon_i = -\frac{2T_0}{\tau_d} \epsilon_i + 2\sigma_{\epsilon 0} \sqrt{\frac{T_0}{\tau_d}} r_i + V'_{rf} z_i + V_{ind}(z_i) \quad (2)$$

$$\Delta z_i = \frac{\alpha c T_0}{E_0} (\epsilon_i + \Delta \epsilon_i) \quad , \quad (3)$$

with T_0 the revolution period, τ_d the damping time, $\sigma_{\epsilon 0}$ the nominal rms energy spread, V'_{rf} the slope of the rf voltage (a negative quantity), α the momentum compaction factor, and E_0 the machine energy; r_i is a random number from a normal set with mean 0 and rms 1. To calculate λ_z on each turn we bin the macro-particles in z .

For the simulations we take $T_0 = 118$ ns, $E_0 = 1.15$ GeV, rf frequency $\nu_{rf} = 714$ MHz, $\sigma_{\epsilon 0} = 0.07\%$, and $\tau_d = 1.7$ ms. We choose $V_{rf} = 0.8$ MV, where $\sigma_{z 0} = 4.95$ mm, and the synchrotron frequency $\nu_{s 0} = 99$ kHz. Therefore $\nu_{s 0} T_0 = 85$ turns, $\tau_d / T_0 = 14450$ turns. We take $N_p = 30,000$, and for λ_z we use 100 bins extending over $10\sigma_z$. We let the program run for 3 damping times.

As a second method of calculation we use a computer program that solves perturbatively the time independent, linearized Vlasov equation, including the effects of potential well distortion, looking for unstable modes. [12] Beyond the threshold current we assume the average energy distribution remains gaussian, with the rms width σ_ϵ increasing to keep the beam just at the threshold condition.

IV. RESULTS

A. The Instability Threshold

For this wakefield the instability threshold is normally easy to find from the turn-by-turn tracking results. Below threshold the moments of the distributions are well behaved, above threshold they undergo macroscopic oscillations (see Fig. 3). The threshold $N_{th} \approx 1.15 \times 10^{10}$.

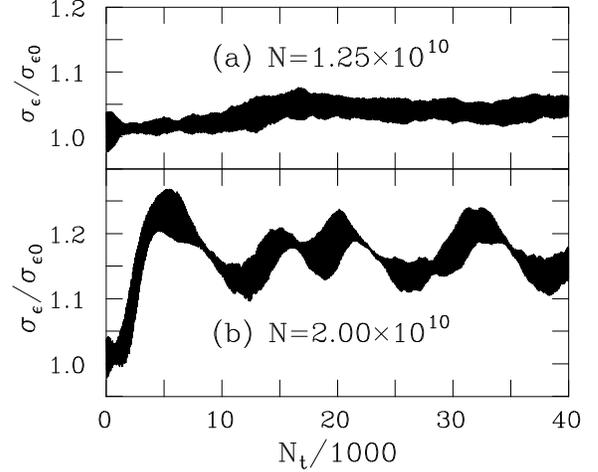


Figure 3. The turn-by-turn rms energy spread just above threshold (a) and at a higher current. (b)

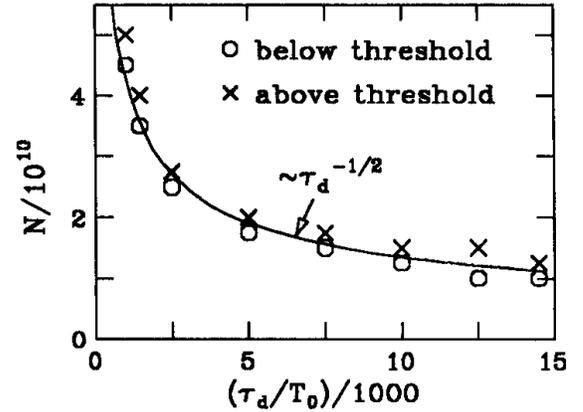


Figure 4. N_{th} vs. τ_d obtained by tracking.

When we artificially reduce the damping time in tracking we find that the threshold N_{th} increases significantly (see Fig. 4). Fitting to a power law we find that N_{th} varies approximately as $\tau_d^{-1/2}$ (the curve in Fig. 4). This agrees with the weak growth expected in a purely resistive machine, which varies as $\sim e^{aN^{2t}}$, with a a constant and t time. [5] Note that when we repeat the tracking procedure for the old, inductive vacuum chamber the threshold increases by only 30% as the damping time is decreased by a factor of 15. The two instabilities are quite different: the old was a strong instability and the new a weak one.

According to a formula in Ref. [5] (Eq. 27) we can raise the threshold by $\sim 1 \times 10^{10}$ if we add a pure inductance of $L = 2$ nH (by adding a term $-cL\lambda'_z$ to V_{ind}). This is roughly what we find; and at $N = 4 \times 10^{10}$ $\sigma_\epsilon / \sigma_{\epsilon 0}$ has decreased from 1.5 to 1.1. This suggests that even a small amount of tune spread can damp this instability.

B. Average Bunch Properties

Above threshold the oscillations in the moments of the distributions obtained by tracking can be large (at $N = 3.5 \times 10^{10}$ the variation in σ_z is $\pm 28\%$) and the pattern can vary greatly. It depends sensitively on, for example, N_p , which (for practical rea-

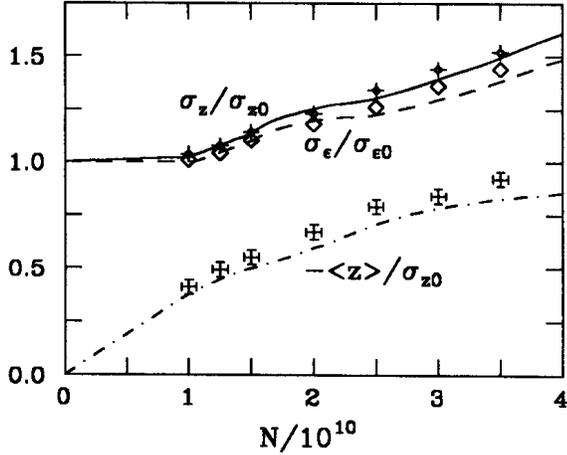


Figure 5. Average bunch properties vs N . Shown are tracking results (plotting symbols) and the Vlasov method results (curves).

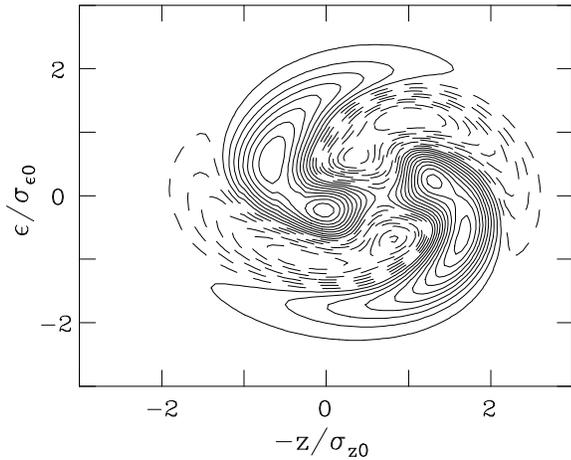


Figure 6. Mode shape at $N = 2 \times 10^{10}$.

sons) it is difficult for us to increase significantly, while keeping τ_d fixed. However, the average amplitudes appear to be rather insensitive to changes in N_p . Fig. 5 gives the average values of σ_z , σ_ϵ , and $\langle z \rangle$ as functions of N , and also shows the Vlasov equation solution for comparison. The agreement is quite good. Note that for such a resistive impedance the potential well distortion is small; the bunch lengthening is largely due to the increase in energy spread above threshold. Thus, by adding 2 nH of pure inductance to reduce the energy spread at 4×10^{10} the 60% increase in σ_z is reduced to 30%.

Fig. 6 gives contours of phase space of the unstable mode at $N = 2 \times 10^{10}$ as calculated by the Vlasov method. Tracking gives a similar result. We see a quadrupole mode, with a slight asymmetry, that has been shifted forward.

C. The Spectrum

The mode frequencies as function of N , as obtained by the Vlasov method are shown in Fig. 7. A dot represents a stable mode, an 'X' an unstable mode, with its size proportional to the growth rate. The strongest unstable mode is a quadrupole mode beginning at $N = 1 \times 10^{10}$ with $\nu = 1.95\nu_{s0}$, and then

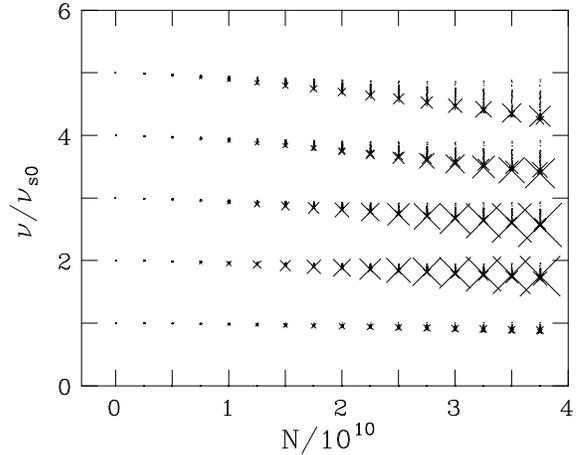


Figure 7. Modes obtained by the Vlasov method.

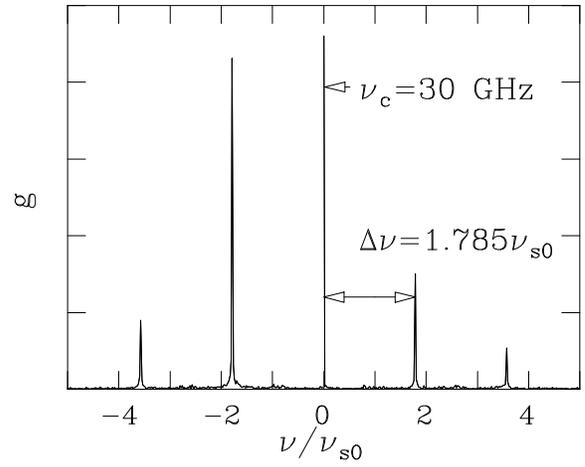


Figure 8. Spectrum for $N = 4 \times 10^{10}$, $\tau_d/T_0 = 1450$.

continuing with a slope of $-0.07\nu_{s0}/10^{10}$. As was the case for a purely resistive machine the instability can be described by coupling of two radial modes with the same azimuthal mode number.

For the tracking results, by Fourier transforming any of the turn-by-turn moments of the distribution we can obtain the spectrum. Alternatively, we can simulate what a spectrum analyzer does by calculating [13]

$$g(\omega) = \left| \sum_k e^{i\omega k T_0} \tilde{\lambda}_{zk}(\omega) \right| \quad (4)$$

with k the turn number and $\tilde{\lambda}_{zk}(\omega)$ the Fourier transform of the distribution on the k^{th} turn. We find sidebands at $1.9\nu_{s0}$, but because of numerical noise we have poor resolution. To give an example with good resolution let us set $N_p = 150,000$, artificially reduce τ_d by a factor of 10, and consider $N = 4 \times 10^{10}$. In Fig. 8 we display the result near a central frequency $\nu_c = 30$ GHz. We see sidebands of the revolution frequency separated by $1.785\nu_{s0}$. Note that the sideband amplitudes are not of equal height. In general, potential well distortion tends to result in an asymmetric mode shape; therefore, for frequencies $\nu_c \gtrsim c/2\pi\sigma_z$ we expect the sidebands to be of unequal height.

V. COMPARISON WITH MEASUREMENTS [2]

In the measurements the bunch length is smaller and the beam profile more asymmetric than before, and above threshold a frequency just below $2\nu_{s0}$ is observed, which are consistent with a resistive wakefield and our simulations. In detail, the agreement is not good unless we assume our wake is missing about 2 nH ($|Z/n| = 0.1 \Omega$) of pure inductance. The measurements give: $N_{th} = 1.5 - 2.0 \times 10^{10}$, and at $N = 4 \times 10^{10}$ $\sigma_e/\sigma_{e0} = 1.15$, and $\sigma_z/\sigma_{z0} = 1.25$; the calculations including the inductance give: $N_{th} = 2.0 \times 10^{10}$, and at $N = 4 \times 10^{10}$ $\sigma_e/\sigma_{e0} = 1.10$, and $\sigma_z/\sigma_{z0} = 1.30$. The unstable mode frequency at threshold and the slope, given by measurements (calculations, including 2 nH): 1.77 (1.87) ν_{s0} and -0.06 (-0.07) $\nu_{s0}/10^{10}$.

Finally, how can we understand the reduction of the measured threshold when the damping ring impedance was reduced? In the old, inductive machine there was a strong instability observed at 3×10^{10} . In an inductive machine there is a large incoherent synchrotron tune spread which will Landau damp weaker instabilities which otherwise might appear at lower currents. By removing mostly inductive elements, and thereby changing the character of the ring to a resistive one, we have removed this tune spread, and presumably are now able to observe one of these weaker instabilities.

VI. ACKNOWLEDGEMENTS

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