

CHARACTERIZATION OF AN RF-DRIVEN PLASMA ION SOURCE FOR HEAVY ION FUSION*

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Abstract

We are testing a high-current-density high-brightness Argon-Ion Source for Heavy Ion Fusion applications. The 100-kV 20- μ s source has produced up to 5 mA of Ar⁺ in a single beamlet. The extraction current density is 100 mA/cm². We have measured the emittance of the beamlet, and the fraction of Ar⁺⁺ ions under several operating conditions. We present measurements of the extracted current density as a function of RF power and gas pressure (~ 2 mT), current density uniformity, and energy dispersion (due to charge exchange). We are testing a 80-kV 61-hole multi-beamlet array that will produce a total current > 200 mA. In the current experiments the beamlets are not merged into a single beam. A 500-kV experiment where the beamlets will be merged to a produce 0.5-A beam is being planned.

BACKGROUND

Following a proposal that the usual limits on brightness for compact ion-beam sources used in Heavy Ion Fusion can be circumvented by using a multi-beamlet injector [1] we have started an experimental program to examine practical issues. The final source envisioned will start with ~100 5-mA beamlets across a 100-kV gap. The beamlets will be focused by Einzel Lens while their energy is increased to about 1.6 MeV. The beamlets are then merged to produce a 0.5-A beam with a normalized emittance below 1 π -mm-mrad.

Beyond providing a low-temperature source that can provide ion emission densities of ~100 mA/cm²[2], the main physics issues involved in the multi-beamlet approach are emittance growth and envelope matching in the merging process. In computer simulations, if the initial emittance is in the range of measured values, the final emittance increases only weakly with the initial emittance of the individual beamlets.

We have available three test stands for developing the concept. The first has a small chamber with voltages below 50 kV. The second has increased diagnostics and 100 kV. The third larger system is capable of 500 kV with a 20 μ sec pulse [3].

We will first describe the characteristics of a single beamlet. We will then give early results from a "61-beamlet" array. During the next year we have a series of planned experiments that move the design closer to a source that could be used in the driver.

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STUDYING INDIVIDUAL BEAMLETS

We are using an rf plasma source to produce the Argon ions. The plasma chamber is 33 cm in diameter with multicusp permanent magnets to confine plasma. RF power (~11 MHz, > 10 kW) is applied to the source via a 2-turn, 11-cm diameter antenna inside the chamber for producing beam pulses of 20 μ s at up to 10 Hz. We have shown that we can extract 100 mA/cm² from the chamber. Optimum performance has been with ~2 mT gas in the plasma chamber.

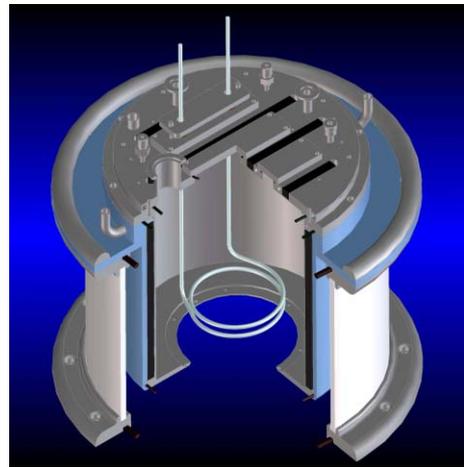


Fig. 1: RF plasma source used in the experiments.

Beam Current

Our first step was to characterize the current in a single beamlet from the rf plasma source (see Fig. 2).

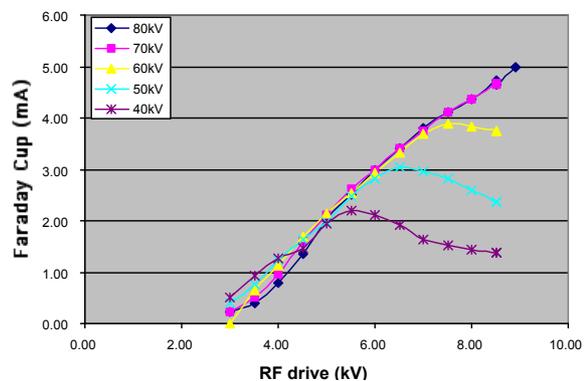


Fig. 2: Extracted current in a single beamlet at 2 mT source pressure. The aperture was 2.5 mm diameter. The curves are different gap potentials. The "10 kV RF drive" is about 12 kW of power from the rf source.

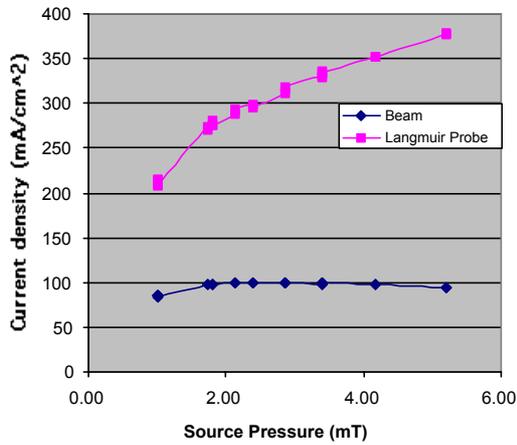


Fig. 3: Effect of source pressure on extracted current and the ion density in the chamber. Gap voltage was 80 kV and the RF drive was 8.5 kV.

As shown in Fig. 3 increasing the source pressure above 2 mT continued to increase the ion density in the chamber, but did not yield higher extracted currents from the gap. Increasing the source pressure increases charge exchange in the gap.

Charge state and charge exchange

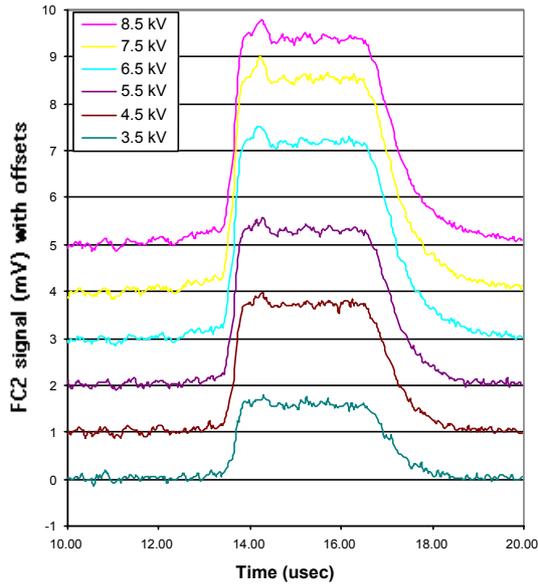


Fig. 4: Arrival time of ions at the back Faraday Cup.

A crude measurement of the ionization state was done by using time-of-flight information. We have a 1.5 meter drift section between pulsed dipole plates and a Faraday Cup. Since there was no focusing, the beam diameter at the cup was larger than it's acceptance. We switched the beam into the Faraday Cup during the center of the pulse. The Ar^{++} ions will have higher momentum, and the front edge of the pulse will arrive at the Faraday Cup before the Ar^+ ions. At drive powers higher than 5 kW, we did see a "front step" in the ion current at the cup (see Fig. 4). We

estimated that less than 5% of the extracted ions were in the Ar^{++} state.

An early concern was that collisions between the extracted ions and the background gas would yield an excessive amount of charge exchange. The newly created ions would leave the gap with lower energies and different transverse motion. Using the dipole plates as an "energy analyzer" we looked for a shift in the profile at the slit cup. We even increased the background gas pressure in the gap to enhance the effect. Our preliminary conclusion is that the energy spread is small. It is conceivable that most of the ions that undergo collisions do not reach the Faraday Cup. Further experiments are planned to improve our ability to resolve energy spread.

Emittance

We used a narrow slit and slit cup to examine the $x - x'$ phase space of a beamlet. An example is shown in Fig. 5.

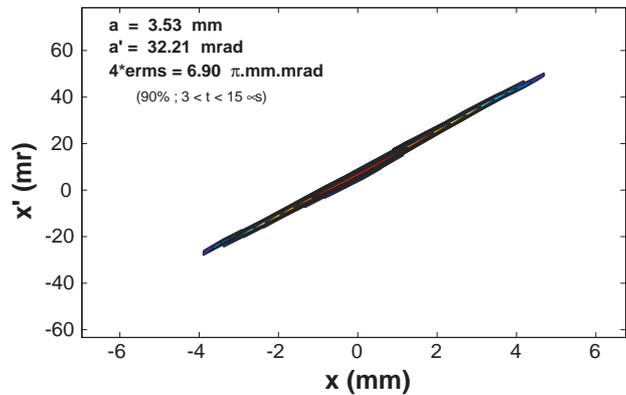


Fig. 5: Phase space of a beamlet 12 cm from the extraction plate. Operation at 80kV, 2.1 mT source pressure, 6.5 kV RF drive. Cut taken at 90% of the peak valve. Early values removed to reduce noise from the spark gaps.

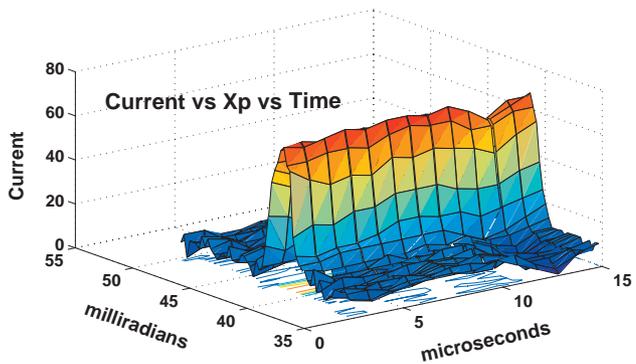


Fig. 6: The current at the slit cup verses angle from the slit as a function of time. Same parameters as in Fig. 5.

The beam parameters were stable over the 3 to 15 μs period as shown in Fig. 6. The deviation at 12 μs is believed to be from a problem with the diagnostics. The normalized emittance for the beamlet was $\sim 0.015 \pi \cdot mm \cdot mrad$.

CURRENT EXPERIMENTS

We are presently investigating a multi-beamlet extraction array. The 1.6-cm gap is held with three high-gradient insulators. The bottom beam-forming plate in Fig. 7 has Pierce Cones for each of the beamlets. The apertures in the bottom plate are 2.2 mm diameter. The holes in the top plate are 4.0 mm diameter. Figure 8 shows the change in the opacity of a kapton sheet hit with 50-kV 20- μ s beamlets from the array. The kapton sheet was 2 cm from the extraction plate. Figure 9 is a lineout of the image shown in Fig. 8.



Fig. 7: Extraction plate for the 61-beamlet array.

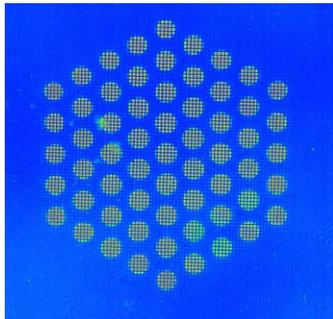


Figure 8. Image produced by ions hitting a Kapton sheet.

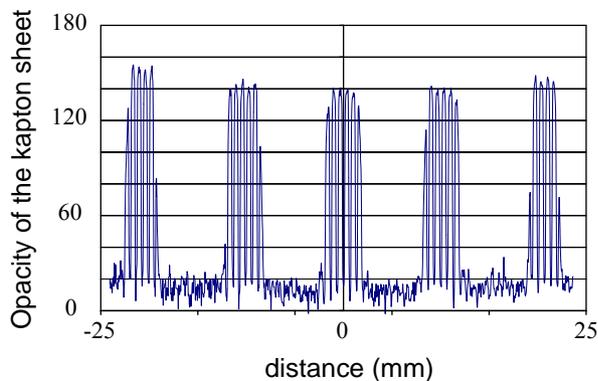


Fig. 9: Line out of the above image taken at the first row of complete beamlets. The fine structure is from a metal screen placed in front of the Kapton sheet to stop ion charge buildup.

NEXT STEPS

After further characterization of the beamlets, our next step will be to add Einzel Lens to the present configuration (shown in Fig. 10). In latter versions the lens will be used to boost the energy of the beamlets. We plan to separate the lens with high-gradient insulators. The separation of the lens is about 1 cm.

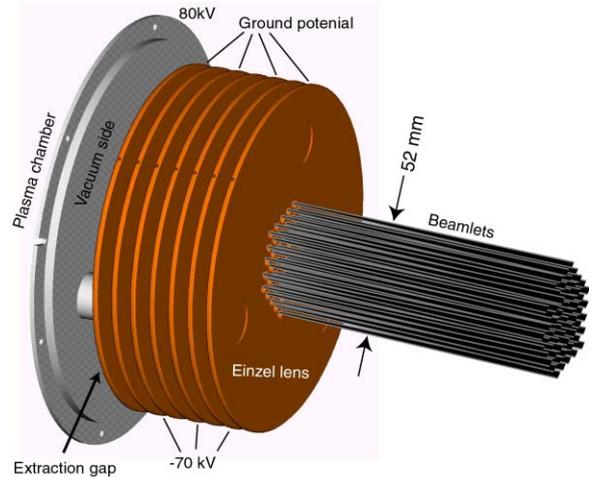


Fig. 10: Next experiment with extraction gap and Einzel Lens.

Merging of the beamlets will not be done at this stage because the perveance is too high. Also, minimizing the final transverse phase space will require curved plates for the Einzel Lens. We are designing an experiment with the extraction gap at ~ 100 kV with post acceleration up to 500 keV.

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